

24.5: The Partition Function for A Gas of Indistinguishable, Non-interacting, Separable-modes Molecules

We represent the successive molecular energy levels as ϵ_i and the successive translational, rotational, vibrational, and electronic energy levels as $\epsilon_{t,a}$, $\epsilon_{r,b}$, $\epsilon_{v,c}$, and $\epsilon_{e,d}$. Now the first subscript specifies the energy mode; the second specifies the energy level. We approximate the successive energy levels of a diatomic molecule as

$$\epsilon_1 = \epsilon_{t,1} + \epsilon_{r,1} + \epsilon_{v,1} + \epsilon_{e,1}$$

$$\epsilon_2 = \epsilon_{t,2} + \epsilon_{r,1} + \epsilon_{v,1} + \epsilon_{e,1}$$

...

$$\epsilon_i = \epsilon_{t,a} + \epsilon_{r,b} + \epsilon_{v,c} + \epsilon_{e,d}$$

...

In [Section 22.1](#), we find that the partition function for the molecule becomes

$$z = \sum_{a=1}^{\infty} \sum_{b=1}^{\infty} \sum_{c=1}^{\infty} \sum_{d=1}^{\infty} g_{t,a} g_{r,b} g_{v,c} g_{e,d} \times \exp \left[\frac{-(\epsilon_{t,a} + \epsilon_{r,b} + \epsilon_{v,c} + \epsilon_{e,d})}{kT} \right]$$

$$= z_t z_r z_v z_e$$

where z_t , z_r , z_v , and z_e are the partition functions for the individual kinds of motion that the molecule undergoes; they are sums over the corresponding energy levels for the molecule. This is essentially the same argument that we use in [Section 22.1](#) to show that the partition function for an N -molecule system is a product of N molecular partition functions:

$$Z = z^N.$$

We are now able to write the partition function for a gas containing N molecules of the same substance. Since the molecules of a gas are indistinguishable, we use the relationship

$$Z_{\text{indistinguishable}} = \frac{1}{N!} z^N = \frac{1}{N!} (z_t z_r z_v z_e)^N$$

To make the notation more compact and to emphasize that we have specialized the discussion to the case of an ideal gas, let us replace “ $Z_{\text{indistinguishable}}$ ” with “ Z_{IG} ”. Also, recognizing that $N!$ enters the relationship because of molecular indistinguishability, and molecular indistinguishability arises because of translational motion, we regroup the terms, writing

$$Z_{\text{IG}} = \left[\frac{(z_t)^N}{N!} \right] (z_r)^N (z_v)^N (z_e)^N$$

Our goal is to calculate the thermodynamic properties of the ideal gas. These properties depend on the natural logarithm of the ideal-gas partition function. This is a sum of terms:

$$\ln Z_{\text{IG}} = \ln \left[\frac{(z_t)^N}{N!} \right] + N \ln z_r + N \ln z_v + N \ln z_e$$

In our development of classical thermodynamics, we find it convenient to express the properties of substance on a per-mole basis. For the same reasons, we focus on evaluating $\ln Z_{\text{IG}}$ for one mole of gas; that is, for the case that N is Avogadro's number, \bar{N} . We now examine the relationships that enable us to evaluate each of these contributions to $\ln Z_{\text{IG}}$.

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