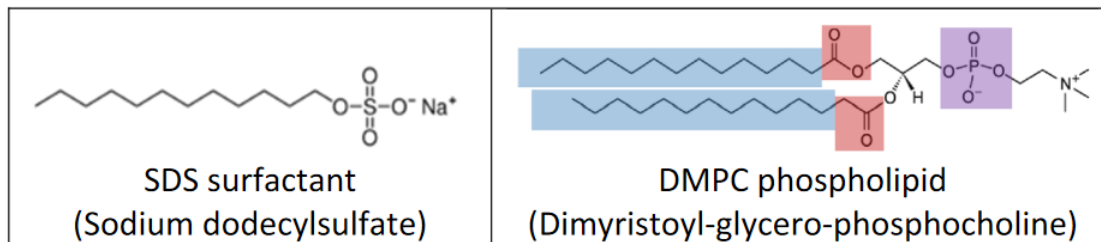
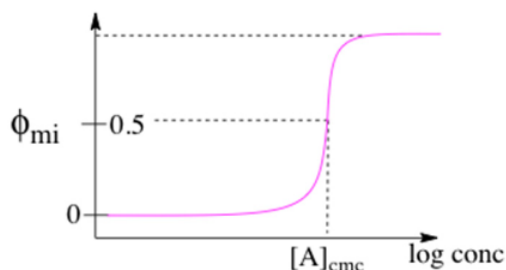


## 19.1: Micelle Formation

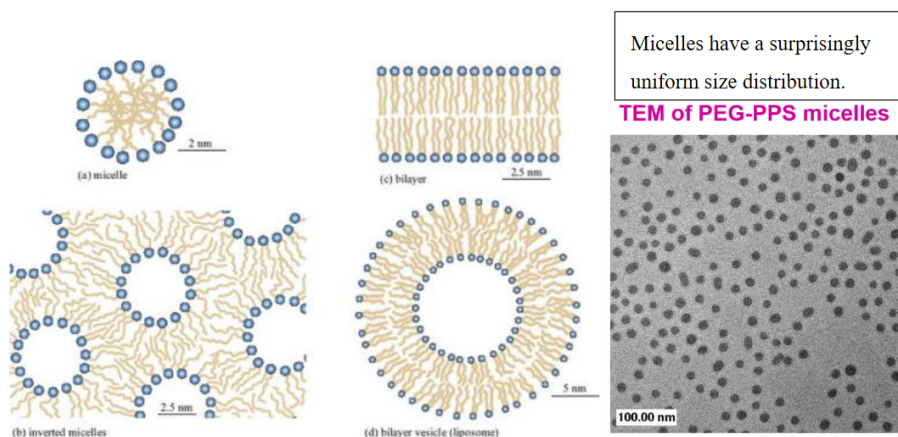
In particular, we will focus on micellar structures formed from a single species of amphiphilic molecule in aqueous solution. These are typically lipids or surfactants that have a charged or polar head group linked to one or more long hydrocarbon chains.



Such amphiphiles assemble into a variety of structures, the result of which depends critically on the concentration, composition, and temperature of the system. For SDS surfactant, micelles are favored. These condense hydrophobic chains into a fluid like core and present the charged head groups to the water. The formation of micelles is observed above a **critical micelle concentration** (CMC).

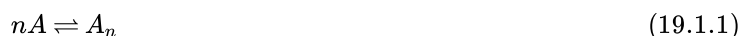


As the surfactant is dissolved, the solution is primarily monomeric at low concentration, but micelles involving 30–100 molecules suddenly appear for concentrations greater than the CMC.



Reprinted from <http://swartz-lab.epfl.ch/page-20594-en.html>.

To begin investigating this phenomenon, we can start by simplifying the equilibrium to a two-state form:



$K_n$  is the equilibrium constant for assembling a micelle with  $n$  amphiphiles from solution.  $n$  is called the **aggregation number**.

$$K_n = \frac{[A_n]}{[A]^n} = e^{-\Delta G_{micelle}^0/k_B T} \quad (19.1.2)$$

The total number of  $A$  molecules present is the sum of the free monomers and those monomers present in micelles:

$$CTOT = [A] + n[A_n]. \quad (19.1.3)$$

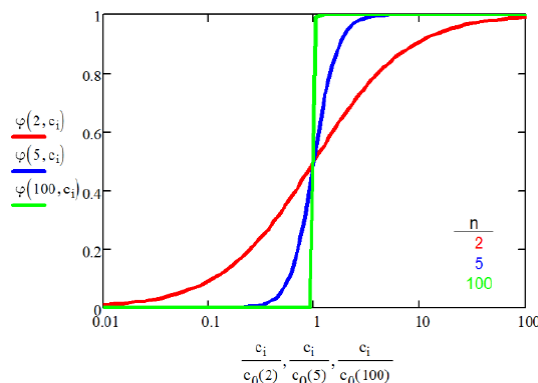
The fraction of monomers present in micelles:

$$\phi_m = \frac{n[A_n]}{CTOT} = \frac{n[A_n]}{[A] + n[A_n]} = \frac{nK_n[A]^{n-1}}{1 + nK_n[A]^{n-1}} \quad (19.1.4)$$

This function has an inflection point at the CMC, for which the steepness of the transition increases with  $n$ . Setting  $\phi_m = 0.5$ , we obtain the CMC ( $c_0$ ) as

$$c_0 = [A]_{cmc} = (nK_n)^{-\frac{1}{n-1}} \quad (19.1.5)$$

Function steepens with aggregation number  $n$ :



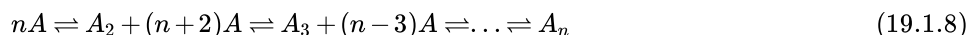
Thus for large  $n$ , and cooperative micelle formation:

$$\Delta G_{micelle}^0 = -RT \ln c_0 \quad (19.1.6)$$

Note the similarity of Equation 19.1.2 to the results for fractional helicity in the helix-coil transition:

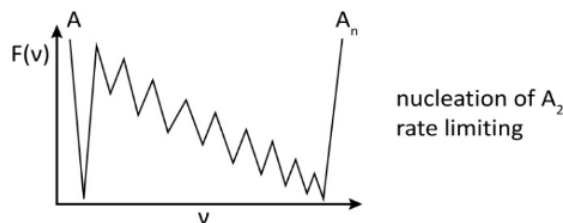
$$\frac{s^n}{1 + s^n} \quad (19.1.7)$$

This similarity indicates that a cooperative model exists for micelle formation in which the aggregation number reflects the number of cooperative units in the process. Cooperativity can be obtained from models that require surmounting a high nucleation barrier before rapidly adding many more molecules to reach the micelle composition. The simplest description of such a process would proceed in a step-wise growth form (a zipper model) for  $n$  copies of monomer  $A$  assembling into a single micelle  $A_n$ .



$$K_n = \prod_{i=1}^{n-1} K_i \quad K_i = \frac{k_f(i \rightarrow i+1)}{k_r(i+1 \rightarrow i)} \quad (19.1.9)$$

Examples of how the energy landscape looks as a function of oligomerization number  $v$  are shown below. However, if you remove the short-range correlation, overall we expect the shape of the energy landscape to still be two-state depending on the nucleation mechanism.



This picture is overly simple though, since it is not a one-dimensional chain problem. Rather, we expect that there are equilibria connecting all possible aggregation number clusters to form larger aggregates. A more appropriate description of the free energy barrier for nucleating a micelle is similar to classical nucleation theory for forming a liquid droplet from vapor.

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D. H. Boal, Mechanics of the Cell, 2nd ed. (Cambridge University Press, Cambridge, UK, 2012), p. 250.

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