

19.4: The Rigid Rotor

The rigid rotor is a simple model of a rotating stick in three dimensions (or, if you prefer, of a molecule). We consider the stick to consist of two point-masses at a fixed distance. We then reduce the model to a one-dimensional system by considering the rigid rotor to have one mass fixed at the origin, which is orbited by the reduced mass μ , at a distance r . The cartesian coordinates, x, y, z , are then replaced by three spherical polar coordinates: the co-latitude (zenith) angle θ , the longitudinal (azimuth) angle ϕ , and the distance r . The TISEq of the system in spherical coordinates is:

$$-\frac{\hbar^2}{2I} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right] \psi(r) = E_\ell \psi(r), \quad (19.4.1)$$

where $I = \mu r^2$ is the moment of inertia. After a little effort, the eigenfunctions can be shown to be the spherical harmonics $Y_\ell^{m_\ell}(\theta, \phi)$.¹ The eigenvalues are simply:

$$E_\ell = \frac{\hbar^2}{2I} \ell(\ell + 1), \quad (19.4.2)$$

where $\ell = 0, 1, 2, \dots$ is the *azimuthal quantum number*, and $m_\ell = -\ell, -\ell + 1, \dots, \ell - 1, \ell$ is the *magnetic quantum number*. Each energy level E_ℓ is $(2\ell + 1)$ -fold degenerate in m_ℓ . Notice that the energy does not depend on the second index m_ℓ , and the functions with fixed $m_\ell = -\ell, -\ell + 1, \dots, \ell - 1, \ell$ have the same energy. Since this problem was, in fact, a one-dimensional problem, it results in just one quantum number ℓ , similarly to the previous two cases. The index m_ℓ that appears in the spherical harmonics will assume some importance in future chapters.

1. For a description of the spherical harmonics see [here](#)

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