

## 14.1: Time-Dependent Vector Potentials

The full N-electron non-relativistic Hamiltonian  $H$  discussed earlier in this text involves the kinetic energies of the electrons and of the nuclei and the mutual Coulombic interactions among these particles

$$H = \sum_{a=1,M} -\left(\frac{\hbar^2}{2m_a}\right) \nabla_a^2 + \sum_j \left[ \left(-\frac{\hbar^2}{2m_e}\right) \nabla_j^2 - \sum_a Z_a \frac{e^2}{r_{j,a}} \right] + \sum_{j < k} \frac{e^2}{r_{j,k}} + \sum_{a < b} Z_a Z_b \frac{e^2}{R_{a,b}}.$$

When an electromagnetic field is present, this is not the correct Hamiltonian, but it can be modified straightforwardly to obtain the proper  $H$ .

### The Time-Dependent Vector $\mathbf{A}(\mathbf{r}, t)$ Potential

The only changes required to achieve the Hamiltonian that describes the same system in the presence of an electromagnetic field are to replace the momentum operators  $\mathbf{P}_a$  and  $\mathbf{p}_j$  for the nuclei and electrons, respectively, by  $(\mathbf{P}_a - Z_a e/c \mathbf{A}(\mathbf{R}_a, t))$  and  $(\mathbf{p}_j - e/c \mathbf{A}(\mathbf{r}_j, t))$ . Here  $Z_a e$  is the charge on the  $a$ th nucleus,  $-e$  is the charge of the electron, and  $c$  is the speed of light.

The vector potential  $\mathbf{A}$  depends on time  $t$  and on the spatial location  $\mathbf{r}$  of the particle in the following manner:

$$\mathbf{A}(\mathbf{r}, t) = 2\mathbf{A}_0 \cos(\omega t - \mathbf{k} \cdot \mathbf{r}).$$

The circular frequency of the radiation  $\omega$  (radians per second) and the wave vector  $\mathbf{k}$  (the magnitude of  $\mathbf{k}$  is  $|\mathbf{k}| = \frac{2\pi}{\lambda}$ , where  $\lambda$  is the wavelength of the light) control the temporal and spatial oscillations of the photons. The vector  $\mathbf{A}_0$  characterizes the strength (through the magnitude of  $\mathbf{A}_0$ ) of the field as well as the direction of the  $\mathbf{A}$  potential; the direction of propagation of the photons is given by the unit vector  $\mathbf{k}/|\mathbf{k}|$ . The factor of 2 in the definition of  $\mathbf{A}$  allows one to think of  $\mathbf{A}_0$  as measuring the strength of both  $e^{i(\omega t - \mathbf{k} \cdot \mathbf{r})}$  and  $e^{-i(\omega t - \mathbf{k} \cdot \mathbf{r})}$  components of the  $\cos(\omega t - \mathbf{k} \cdot \mathbf{r})$  function.

### The Electric $\mathbf{E}(\mathbf{r}, t)$ and Magnetic $\mathbf{H}(\mathbf{r}, t)$ Fields

The electric  $\mathbf{E}(\mathbf{r}, t)$  and magnetic  $\mathbf{H}(\mathbf{r}, t)$  fields of the photons are expressed in terms of the vector potential  $\mathbf{A}$  as

$$\mathbf{E}(\mathbf{r}, t) = -\frac{1}{c} \frac{\partial \mathbf{A}}{\partial t} = \frac{\omega}{c} \mathbf{A}_0 \sin(\omega t - \mathbf{k} \cdot \mathbf{r})$$

$$\mathbf{H}(\mathbf{r}, t) = \nabla \times \mathbf{A} = \mathbf{k} \times \mathbf{A}_0 \sin(\omega t - \mathbf{k} \cdot \mathbf{r}).$$

The  $\mathbf{E}$  field lies parallel to the  $\mathbf{A}_0$  vector, and the  $\mathbf{H}$  field is perpendicular to  $\mathbf{A}_0$ ; both are perpendicular to the direction of propagation of the light  $\mathbf{k}/|\mathbf{k}|$ .  $\mathbf{E}$  and  $\mathbf{H}$  have the same phase because they both vary with time and spatial location as  $\sin(\omega t - \mathbf{k} \cdot \mathbf{r})$ . The relative orientations of these vectors are shown below.



Figure 14.1.1: Insert caption here!

### The Resulting Hamiltonian

Replacing the nuclear and electronic momenta by the modifications shown above in the kinetic energy terms of the full electronic and nuclear-motion hamiltonian results in the following **additional** factors appearing in  $H$ :

$$H_{int} = \sum_j \left[ \frac{ie\hbar}{m_e c} \mathbf{A}(\mathbf{r}_j, t) \cdot \nabla_j + \left( \frac{e^2}{2m_e c^2} \right) |\mathbf{A}(\mathbf{r}_j, t)|^2 \right] + \sum_a \left[ \left( iZ_a \frac{e\hbar}{m_a c} \right) \mathbf{A}(\mathbf{R}_a, t) \cdot \nabla_a + \left( \frac{Z_a^2 e^2}{2m_a c^2} \right) |\mathbf{A}(\mathbf{R}_a, t)|^2 \right].$$

These so-called interaction perturbations  $H_{int}$  are what induces transitions among the various electronic/vibrational/rotational states of a molecule. The one-electron additive nature of  $H_{int}$  plays an important role in determining the kind of transitions that  $H_{int}$  can induce. For example, it causes the most intense electronic transitions to involve excitation of a single electron from one orbital to another (e.g., the [Slater-Condon rules](#)).

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