

22.3.6: vi. Problem Solutions

Q1

a. All the Slater determinants have in common the $|1s\alpha 1s\beta 2s\alpha 2s\beta|$ "core" and hence this component will not be written out explicitly for each case.

$${}^3P(M_L = 1, M_S = 1) = |p_1\alpha p_0\alpha| \quad (22.3.6.1)$$

$$= \left| \frac{1}{\sqrt{2}}(p_x + ip_y)\alpha(p_z)\alpha \right| \quad (22.3.6.2)$$

$$= \frac{1}{\sqrt{2}}(|p_x\alpha p_z\alpha| + i|p_y\alpha p_z\alpha|) \quad (22.3.6.3)$$

$${}^3P(M_L = 0, M_S = 1) = |p_1\alpha p_{-1}\alpha| \quad (22.3.6.4)$$

$$= \left| \frac{1}{\sqrt{2}}(p_x + ip_y)\alpha \right| \quad (22.3.6.5)$$

$$\frac{1}{\sqrt{2}}(p_x - ip_y)\alpha = \frac{1}{2}(|p_x\alpha p_x\alpha| - i|p_x\alpha p_y\alpha| + i|p_y\alpha p_x\alpha| + |p_y\alpha p_y\alpha|) \quad (22.3.6.6)$$

$$= \frac{1}{2}(0 - i|p_x\alpha p_y\alpha|) \quad (22.3.6.7)$$

$$= -i|p_x\alpha p_y\alpha| \quad (22.3.6.8)$$

$${}^3P(M_L = -1, M_S = 1) = |p_{-1}\alpha p_0\alpha| \quad (22.3.6.9)$$

$$= \left| \frac{1}{\sqrt{2}}(p_x - ip_y)\alpha(p_z)\alpha \right| \quad (22.3.6.10)$$

$$= \frac{1}{\sqrt{2}}(|p_x\alpha p_z\alpha| - i|p_y\alpha p_z\alpha|) \quad (22.3.6.11)$$

As you can see, the symmetry of each of these states cannot be labeled with a single irreducible representation of the C_{2v} point group. For example, $|p_x\alpha p_z\alpha|$ is xz (B_1) and $|p_y\alpha p_z\alpha|$ is yz (B_2) and hence the ${}^3P(M_L, M_S = 1)$ functions are degenerate for the C atom and any combination of these functions would also be degenerate. Therefore we can choose new combinations which can be labeled with "pure" C_{2v} point group labels.

$${}^3P(xz, M_S = 1) = |p_x\alpha p_z\alpha| \quad (22.3.6.12)$$

$$= \frac{1}{\sqrt{2}}({}^3P(M_L = 1, M_S = 1) + {}^3P(M_L = -1, M_S = 1)) = {}^3B_1 \quad (22.3.6.13)$$

$${}^3P(yx, M_S = 1) = |p_y\alpha p_x\alpha| \quad (22.3.6.14)$$

$$= \frac{1}{i}({}^3P(M_L = 0, M_S = 1)) = {}^3A_2 \quad (22.3.6.15)$$

$${}^3P(yz, M_S = 1) = |p_y\alpha p_z\alpha| \quad (22.3.6.16)$$

$$= \frac{1}{i\sqrt{2}}({}^3P(M_L = 1, M_S = 1) - {}^3P(M_L = -1, M_S = 1)) = {}^3B_2 \quad (22.3.6.17)$$

Now we can do likewise for the five degenerate 1D states:

$$\begin{aligned}
 {}^1D(M_L = 2, M_S = 0) &= |p_1\alpha p_1\beta| & (22.3.6.18) \\
 &= \left| \frac{1}{\sqrt{2}}(p_x + ip_y)\alpha \frac{1}{\sqrt{2}}(p_x + ip_y)\beta \right| & (22.3.6.19) \\
 &= \frac{1}{2}(|p_x\alpha p_x\beta| + i|p_x\alpha p_y\beta| + i|p_y\alpha p_x\beta| - |p_y\alpha p_y\beta|) & (22.3.6.20) \\
 {}^1D(M_L = -2, M_S = 0) &= |p_{-1}\alpha p_{-1}\beta| & (22.3.6.21) \\
 &= \left| \frac{1}{\sqrt{2}}(p_x - ip_y)\alpha \frac{1}{\sqrt{2}}(p_x - ip_y)\beta \right| & (22.3.6.22) \\
 &= \frac{1}{2}(|p_x\alpha p_x\beta| - i|p_x\alpha p_y\beta| - i|p_y\alpha p_x\beta| - |p_y\alpha p_y\beta|) & (22.3.6.23) \\
 {}^1D(M_L = 1, M_S = 0) &= \frac{1}{\sqrt{2}}(|p_0\alpha p_1\beta| - |p_0\beta p_1\alpha|) & (22.3.6.24) \\
 &= \frac{1}{\sqrt{2}} \left(|(p_z)\alpha \frac{1}{\sqrt{2}}(p_x + ip_y)\beta| - |(p_z)\beta \frac{1}{\sqrt{2}}(p_x + ip_y)\alpha| \right) & (22.3.6.25) \\
 &= \frac{1}{2}(|p_z\alpha p_x\beta| + i|p_z\alpha p_y\beta| - |p_z\beta p_x\alpha| - i|p_z\beta p_y\alpha|) & (22.3.6.26) \\
 {}^1D(M_L = -1, M_S = 0) &= \frac{1}{\sqrt{2}}(|p_0\alpha p_{-1}\beta| - |p_0\beta p_{-1}\alpha|) & (22.3.6.27) \\
 &= \frac{1}{\sqrt{2}} \left(|(p_z)\alpha \frac{1}{\sqrt{2}}(p_x - ip_y)\beta| - |(p_z)\beta \frac{1}{\sqrt{2}}(p_x - ip_y)\alpha| \right) & (22.3.6.28) \\
 &= \frac{1}{2}(|p_z\alpha p_x\beta| - i|p_z\alpha p_y\beta| - |p_z\beta p_x\alpha| + i|p_z\beta p_y\alpha|) & (22.3.6.29) \\
 {}^1D(M_L = 0, M_S = 0) &= \frac{1}{\sqrt{6}}(2|p_0\alpha p_0\beta| + |p_1\alpha p_{-1}\beta| + |p_{-1}\alpha p_1\beta|) & (22.3.6.30) \\
 &= \frac{1}{\sqrt{6}} \left(2|p_z\alpha p_z\beta| + \left| \frac{1}{\sqrt{2}}(p_x + ip_y)\alpha \frac{1}{\sqrt{2}}(p_x - ip_y)\beta \right| + \left| \frac{1}{\sqrt{2}}(p_x - ip_y)\alpha \frac{1}{\sqrt{2}}(p_x + ip_y)\beta \right| \right) & (22.3.6.31) \\
 &= \frac{1}{\sqrt{6}} \left(2|p_z\alpha p_z\beta| + \frac{1}{2}(|p_x\alpha p_x\beta| - i|p_x\alpha p_y\beta| + i|p_y\alpha p_x\beta| + |p_y\alpha p_y\beta|) \right. & (22.3.6.32) \\
 &\quad \left. + \frac{1}{2}(|p_x\alpha p_x\beta| + i|p_x\alpha p_y\beta| - i|p_y\alpha p_x\beta| + |p_y\alpha p_y\beta|) \right) \\
 &= \frac{1}{\sqrt{6}}(2|p_z\alpha p_z\beta| + |p_x\alpha p_x\beta| + |p_y\alpha p_y\beta|) & (22.3.6.33)
 \end{aligned}$$

Analogous to the three 3p states we can also choose combinations of the five degenerate 1D states which can be labeled with "pure" C_{2v} point group labels:

$$\begin{aligned}
 {}^1D(xx - yy, M_S = 0) &= |p_x\alpha p_x\beta| - |p_y\alpha p_y\beta| & (22.3.6.34) \\
 &= ({}^1D(M_L = 2, M_S = 0) + {}^1D(M_L = -2, M_S = 0)) = {}^1A_1 & (22.3.6.35) \\
 {}^1D(yx, M_S = 0) &= |p_x\alpha p_y\beta| + |p_y\alpha p_x\beta| & (22.3.6.36) \\
 &= \frac{1}{i}({}^1D(M_L = 2, M_S = 0) - {}^1D(M_L = -2, M_S = 0)) = {}^1A_2 & (22.3.6.37) \\
 {}^1D(zx, M_S = 0) &= |p_z\alpha p_y\beta| - |p_z\beta p_y\alpha| & (22.3.6.38) \\
 &= ({}^1D(M_L = 1, M_S = 0) + {}^1D(M_L = -1, M_S = 0)) = {}^1B_1 & (22.3.6.39) \\
 {}^1D(zy, M_S = 0) &= |p_z\alpha p_y\beta| - |p_z\beta p_y\alpha| & (22.3.6.40) \\
 &= \frac{1}{i}({}^1D(M_L = 1, M_S = 0) - {}^1D(M_L = -1, M_S = 0)) = {}^1B_2 & (22.3.6.41) \\
 {}^1D(2zz + xx + yy, M_S = 0) &= \frac{1}{\sqrt{6}}(2|p_z\alpha p_z\beta| + |p_x\alpha p_x\beta| + |p_y\alpha p_y\beta|) & (22.3.6.42) \\
 &= {}^1D(M_L = 0, M_S = 0) = {}^1A_1 & (22.3.6.43)
 \end{aligned}$$

The only state left is the 1S :

$${}^1S(M_L = 0, M_S = 0) = \frac{1}{\sqrt{3}}(|p_0\alpha p_0\beta| - |p_1\alpha p_{-1}\beta| - |p_{-1}\alpha p_1\beta|) \quad (22.3.6.44)$$

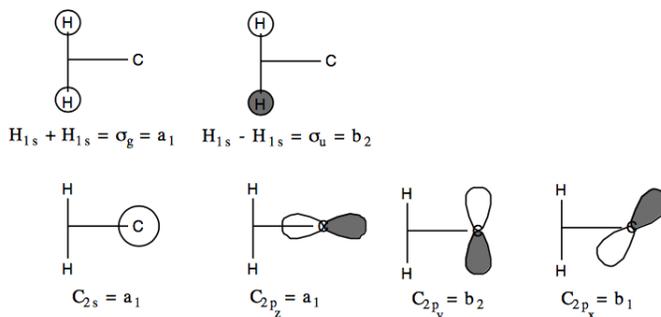
$$= \frac{1}{\sqrt{3}} \left(|p_z\alpha p_z\beta| - \left| \frac{1}{\sqrt{2}}(p_x + ip_y)\alpha \frac{1}{\sqrt{2}}(p_x - ip_y)\beta \right| - \left| \frac{1}{\sqrt{2}}(p_x - ip_y)\alpha \frac{1}{\sqrt{2}}(p_x + ip_y)\beta \right| \right) \quad (22.3.6.45)$$

$$= \frac{1}{\sqrt{3}} \left(|p_z\alpha p_z\beta| - \frac{1}{2}(|p_x\alpha p_x\beta| - i|p_x\alpha p_y\beta| + i|p_y\alpha p_x\beta| + |p_y\alpha p_y\beta|) - \frac{1}{2}(|p_x\alpha p_x\beta| + i|p_x\alpha p_y\beta| - i|p_y\alpha p_x\beta| + |p_y\alpha p_y\beta|) \right) \quad (22.3.6.46)$$

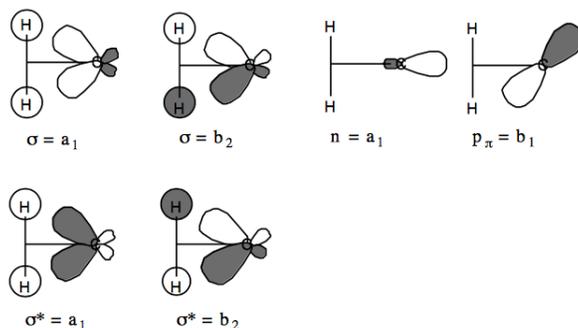
$$= \frac{1}{\sqrt{3}}(|p_z\alpha p_z\beta| - |p_x\alpha p_x\beta| - |p_y\alpha p_y\beta|) \quad (22.3.6.47)$$

Each of the components of this state are A_1 and hence this state has A_1 symmetry.

b. Forming SALC-AOs from the C and H atomic orbitals would generate the following:

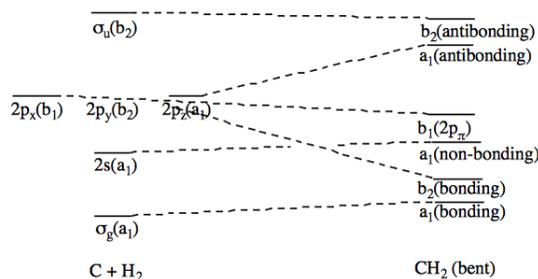


The bonding, nonbonding, and antibonding orbitals of CH₂ can be illustrated in the following manner:



c.

Orbital-correlation diagram for the reaction $C + H_2 \rightarrow CH_2$ (bent)



d. - e. It is necessary to determine how the wave functions found in part a. correlate with states of the CH₂ molecule:

$${}^3P(xz, M_S = 1); {}^3B_1 = \sigma_g^2 s^2 p_x p_z \rightarrow \sigma^2 n^2 p_x p_z \quad (22.3.6.48)$$

$${}^3P(yx, M_S = 1); {}^3A_2 = \sigma_g^2 s^2 p_x p_y \rightarrow \sigma^2 n^2 p_x p_y \quad (22.3.6.49)$$

$${}^3P(yz, M_S = 1); {}^3B_2 = \sigma_g^2 s^2 p_y p_z \rightarrow \sigma^2 n^2 p_y p_z \quad (22.3.6.50)$$

$${}^1D(xz - yy, M_S = 0); {}^1A_1 \rightarrow \sigma^2 n^2 p_x^2 - \sigma^2 n^2 p_y^2 \quad (22.3.6.51)$$

$${}^1D(yx, M_S = 0); {}^1A_2 \rightarrow \sigma^2 n^2 p_x p_y \quad (22.3.6.52)$$

$${}^1D(zx, M_S = 0); {}^1B_1 \rightarrow \sigma^2 n^2 p_x p_z \quad (22.3.6.53)$$

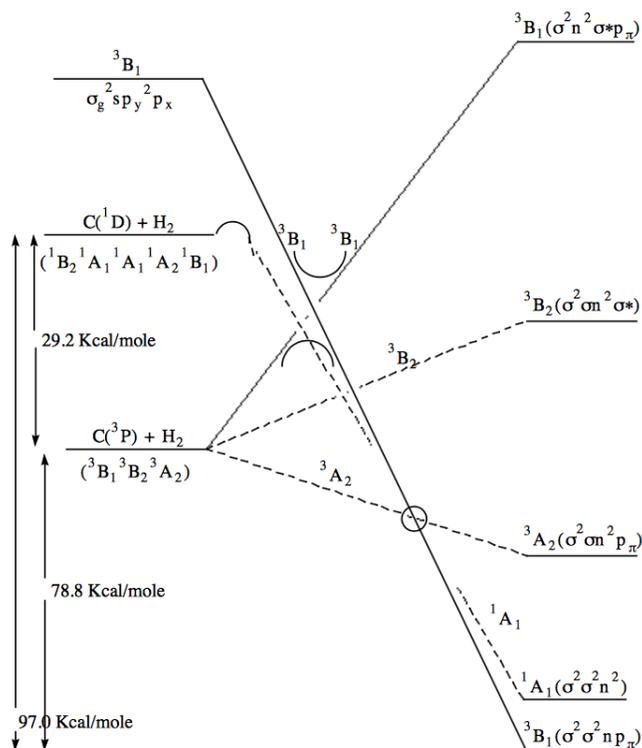
$${}^1D(zy, M_S = 0); {}^1B_2 \rightarrow \sigma^2 n^2 p_y p_z \quad (22.3.6.54)$$

$${}^1D(2zz + xx + yy, M_S = 0); {}^1A_1 \rightarrow 2\sigma^2 n^2 p_z^2 + \sigma^2 n^2 p_x^2 + \sigma^2 n^2 p_y^2 \quad (22.3.6.55)$$

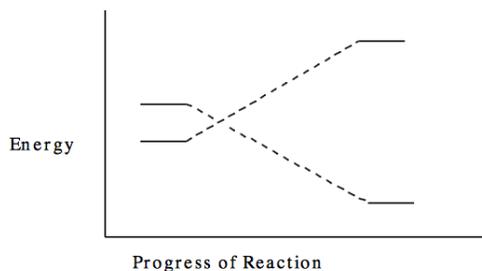
Note, the C + H₂ state to which the lowest ${}^1A_1(\sigma^2 n^2 p_x^2)CH_2$ state decomposes would be $\sigma_g^2 s^2 p_y^2$. This state ($\sigma_g^2 s^2 p_y^2$) cannot be obtained by a simple combination of the 1D states. In order to obtain pure $\sigma_g^2 s^2 p_y^2$ it is necessary to combine 1S with 1D . For example,

$$\sigma_g^2 s^2 p_y^2 = \frac{1}{6} \left(\sqrt{6} {}^1D(0, 0) - 2\sqrt{3} {}^1S(0, 0) \right) - \frac{1}{2} ({}^1D(2, 0) + {}^1D(-2, 0)).$$

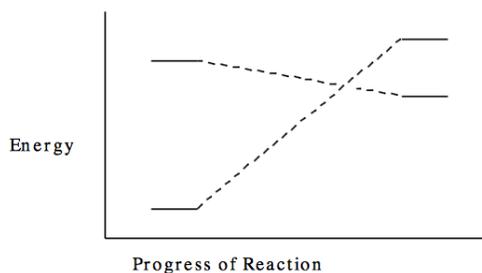
This indicates that a CCD must be drawn with a barrier near the 1D asymptote to represent the fact that 1A_1 CH₂ correlates with a mixture of 1D and 1S carbon plus hydrogen. The C + H₂ state to which the lowest ${}^3B_1(\sigma^2 n^2 p_x)CH_2$ state decomposes would be $\sigma_g^2 s^2 p_x p_z$.



f. If you follow the 3B_1 component of the $C(^3P) + H_2$ (since it leads to the ground state products) to 3B_1CH_2 you must go over an approximately 20 Kcal/mole barrier. Of course this path produces 3B_1CH_2 product. Distortions away from C_{2v} symmetry, for example to $C-s$ symmetry, would make the a_1 and b_2 orbitals identical in symmetry (a'). The b_1 orbitals would maintain their identity group going to a'' symmetry. Thus 3B_1 and 3A_2 (both $^3A''$ in C_s symmetry and odd under reflection through the molecular plane) can mix. The system could thus follow the 3A_2 component of the $C(^3P) + H_2$ surface to the place (marked with a circle on the CCD) where it crosses the 3B_1 surface upon which it then moves and continues to products. As a result, the barrier would be lowered. You can estimate when the barrier occurs (late or early) using thermodynamic information for the reaction (i.e. slopes and asymptotic energies). For example, an early barrier would be obtained for a reaction with the characteristics:



and a late barrier would be obtained for a reaction with the characteristics:



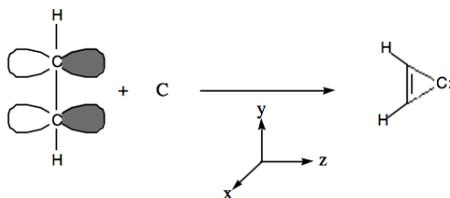
This relation between reaction endothermicity or exothermicity is known as the Hammond postulate. Note that the $C(^3P_1) + H_2 \rightarrow CH_2$ reaction of interest here (see the CCD) has an early barrier.

g. The reaction $C(^1D) + H_2 \rightarrow CH_2(^1A_1)$ should have no symmetry barrier (this can be recognized by the following the $^1A_1(C(^1D) + H_2)$ reactants down to the $^1A_1(CH_2)$ products on the CCD).

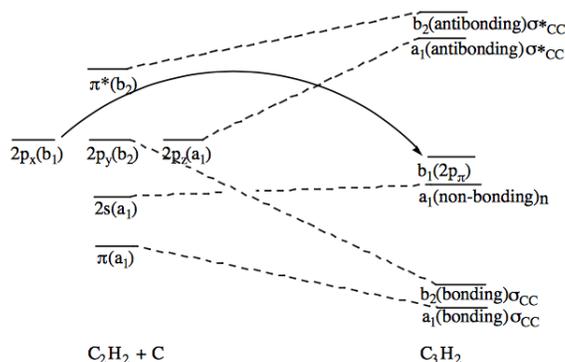
Q2

This problem in many respects is analogous to problem 1.

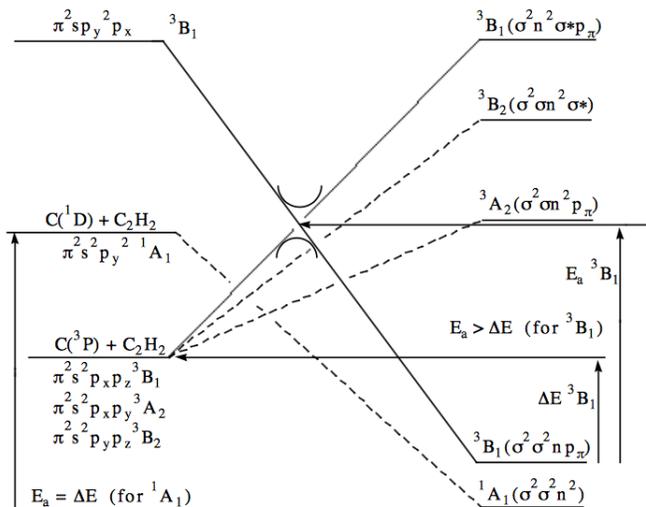
The 3B_1 surface certainly requires a two configuration CI wavefunction; the $\sigma^2\sigma^2np_x(\pi^2p_y^2sp_x)$ and the $\sigma^2n^2p_x\sigma^*(\pi^2s^2p_xp_z)$. The 1A_1 surface could use the $\sigma^2\sigma^2n^2(\pi^2s^2p_y^2)$ only but once again there is a combination of 1D determinants which gives purely this configuration ($\pi^2s^2p_y^2$). Thus mixing of both 1D and 1S determinants are necessary to yield the required $\pi^2s^2p_y^2$ configuration. Hence even the 1A_1 surface would require a multiconfigurational wavefunction for adequate description.



Orbital-correlation diagram for the reaction $C_2H_2 + C \rightarrow C_3H_2$



Configuration correlation diagram for the reaction $C_2H_2 + C \rightarrow C_3H_2$.



Q3

a.

$$\langle \sigma_g | \sigma_g \rangle = \left\langle \frac{1s_A + 1s_B}{\sqrt{2 + 2S}} \middle| \frac{1s_A + 1s_B}{\sqrt{2 + 2S}} \right\rangle \quad (22.3.6.56)$$

$$= \frac{1}{2 + 2S} (\langle 1s_A | 1s_A \rangle + \langle 1s_A | 1s_B \rangle + \langle 1s_B | 1s_A \rangle + \langle 1s_B | 1s_B \rangle) \quad (22.3.6.57)$$

$$= \frac{1}{2 + 2S} (1.000 + 0.753 + 0.753 + 1.000) \quad (22.3.6.58)$$

$$= 0.999 \approx 1 \quad (22.3.6.59)$$

$$\langle \sigma_g | \sigma_u \rangle = \left\langle \frac{1s_A + 1s_B}{\sqrt{2+2S}} \middle| \frac{1s_A - 1s_B}{\sqrt{2-2S}} \right\rangle \quad (22.3.6.60)$$

$$= \frac{\langle 1s_A | 1s_A \rangle + \langle 1s_A | 1s_B \rangle + \langle 1s_B | 1s_A \rangle + \langle 1s_B | 1s_B \rangle}{\sqrt{2+2S}\sqrt{2-2S}} \quad (22.3.6.61)$$

$$= (1.434)(0.534)((1.000) - (0.753) + (0.753) - (1.000)) \quad (22.3.6.62)$$

$$= 0 \quad (22.3.6.63)$$

$$\langle \sigma_u | \rangle = \left\langle \frac{1s_A - 1s_B}{\sqrt{2-2S}} \middle| \frac{1s_A - 1s_B}{\sqrt{2-2S}} \right\rangle \quad (22.3.6.64)$$

$$= \frac{\langle 1s_A | 1s_A \rangle - \langle 1s_A | 1s_B \rangle - \langle 1s_B | 1s_A \rangle + \langle 1s_B | 1s_B \rangle}{2-2S} \quad (22.3.6.65)$$

$$= (2.024)((1.000) - (0.753) - (0.753) + (1.000)) \quad (22.3.6.66)$$

$$= 1.000 \quad (22.3.6.67)$$

b.

$$\langle \sigma_g | h | \sigma_g \rangle = \left\langle \frac{1s_A + 1s_B}{\sqrt{2+2S}} \middle| h \middle| \frac{1s_A + 1s_B}{\sqrt{2+2S}} \right\rangle \quad (22.3.6.68)$$

$$= \frac{\langle 1s_A | h | 1s_A \rangle + \langle 1s_A | h | 1s_B \rangle + \langle 1s_B | h | 1s_A \rangle + \langle 1s_B | h | 1s_B \rangle}{2+2S} \quad (22.3.6.69)$$

$$= (0.285)((-1.110) + (-0.968) + (-0.968) + (-1.110)) \quad (22.3.6.70)$$

$$= -1.184 \quad (22.3.6.71)$$

$$\langle \sigma_u | h | \sigma_u \rangle = \left\langle \frac{1s_A - 1s_B}{\sqrt{2-2S}} \middle| h \middle| \frac{1s_A - 1s_B}{\sqrt{2-2S}} \right\rangle \quad (22.3.6.72)$$

$$= \frac{\langle 1s_A | h | 1s_A \rangle - \langle 1s_A | h | 1s_B \rangle - \langle 1s_B | h | 1s_A \rangle + \langle 1s_B | h | 1s_B \rangle}{2-2S} \quad (22.3.6.73)$$

$$= (2.024)((-1.110) + (0.968) + (0.968) + (-1.110)) \quad (22.3.6.74)$$

$$= -0.575 \quad (22.3.6.75)$$

$$\langle \sigma_g \sigma_g | h | \sigma_g \sigma_g \rangle \equiv \langle gg | gg \rangle = \frac{1}{2+2S} \frac{1}{2+2S} \quad (22.3.6.76)$$

$$\langle (1s_A + 1s_B)(1s_A + 1s_B) | (1s_A + 1s_B)(1s_A + 1s_B) \rangle \quad (22.3.6.77)$$

$$= \frac{1}{(2+2S)^2} (\langle AA | AA \rangle + \langle AA | AB \rangle + \langle AA | BA \rangle + \langle AA | BB \rangle + \langle AB | AA \rangle + \langle AB | AB \rangle + \langle AB | BA \rangle + \langle AB | BB \rangle + \langle BA | AA \rangle + \langle BA | AB \rangle + \langle BA | BA \rangle + \langle BA | BB \rangle + \langle BB | AA \rangle + \langle BB | BA \rangle + \langle BB | BA \rangle + \langle BB | BB \rangle) \quad (22.3.6.78)$$

$$= (0.081)((0.625) + (0.426) + (0.426) + (0.323) + (0.426) + (0.504) + (0.323) + (0.426) + (0.426) + (0.323) + (0.504) + (0.426) + (0.426) + (0.625)) \quad (22.3.6.79)$$

$$= 0.564 \quad (22.3.6.80)$$

$$= 0.564 \quad (22.3.6.81)$$

$$= 0.564 \quad (22.3.6.82)$$

$$\langle uu | uu \rangle = \frac{1}{2-2S} \frac{1}{2-2S} \quad (22.3.6.83)$$

$$\langle (1s_A - 1s_B)(1s_A - 1s_B) | (1s_A - 1s_B)(1s_A - 1s_B) \rangle \quad (22.3.6.84)$$

$$= \frac{1}{(2-2S)^2} (\langle AA | AA \rangle - \langle AA | AB \rangle - \langle AA | BA \rangle \langle AA | BB \rangle - \langle BA | AA \rangle + \langle BA | AB \rangle + \langle BA | BA \rangle - \langle BA | BB \rangle + \langle BA | AA \rangle + \langle BA | AB \rangle + \langle BA | BA \rangle - \langle BA | BB \rangle + \langle BB | AA \rangle - \langle BB | AB \rangle - \langle BB | BA \rangle + \langle BB | BB \rangle) \quad (22.3.6.85)$$

$$= (4.100)((0.625) - (0.426) - (0.426) + (0.323) - (0.426) + (0.504) + (0.323) - (0.426) - (0.426) + (0.323) + (0.504) - (0.426) + (0.323) - (0.426) - (0.426) + (0.625)) \quad (22.3.6.86)$$

$$= 0.582 \quad (22.3.6.87)$$

$$= 0.582 \quad (22.3.6.88)$$

$$= 0.582 \quad (22.3.6.89)$$

$$\begin{aligned}
 \langle gg|uu \rangle &= \frac{1}{2+2S} \frac{1}{2-2S} && (22.3.6.90) \\
 &= \frac{1}{(1s_A+1s_B)(1s_A+1s_B)(1s_A-1s_B)(1s_A-1s_B)} && (22.3.6.91) \\
 &= \frac{1}{2+2S} \frac{1}{2-2S} && (22.3.6.92) \\
 &= \frac{(\langle AA|AA \rangle - \langle AA|AB \rangle - \langle AA|BA \rangle + \langle AA|BB \rangle + \langle AB|AA \rangle - \langle AB|AB \rangle - \langle AB|BA \rangle + \langle AB|BB \rangle + \langle BA|AA \rangle - \langle BA|AB \rangle - \langle BA|BA \rangle + \langle BA|BB \rangle + \langle BB|AA \rangle - \langle BB|AB \rangle - \langle BB|BA \rangle + \langle BB|BB \rangle)}{2+2S} && (22.3.6.93) \\
 &= \frac{(0.285)(2.024)((0.625) - (0.426) - (0.426) + (0.323) + (0.426) - (0.504) - (0.323) + (0.426) + (0.426) - (0.323) - (0.504) - (0.426) + (0.323) - (0.426) - (0.426) + (0.625))}{2+2S} && (22.3.6.94) \\
 &= 0.140 && (22.3.6.95) \\
 \langle gu|gu \rangle &= \frac{1}{2+2S} \frac{1}{2-2S} && (22.3.6.96) \\
 &= \frac{1}{(1s_A+1s_B)(1s_A-1s_B)(1s_A+1s_B)(1s_A-1s_B)} && (22.3.6.97) \\
 &= \frac{1}{2+2S} \frac{1}{2-2S} && (22.3.6.98) \\
 &= \frac{(\langle AA|AA \rangle - \langle AA|AB \rangle + \langle AA|BA \rangle - \langle AA|BB \rangle - \langle AB|AA \rangle + \langle AB|AB \rangle - \langle AB|BA \rangle + \langle AB|BB \rangle + \langle BA|AA \rangle - \langle BA|AB \rangle + \langle BA|BA \rangle - \langle BA|BB \rangle - \langle BB|AA \rangle + \langle BB|AB \rangle - \langle BB|BA \rangle + \langle BB|BB \rangle)}{2+2S} && (22.3.6.99) \\
 &= \frac{(0.285)(2.024)((0.625) - (0.426) + (0.426) - (0.323) - (0.426) + (0.504) - (0.323) + (0.426) + \dots \dots + (0.426) - (0.323) + (0.504) - (0.426) - (0.323) + (0.426) - (0.426) + (0.625))}{2+2S} && (22.3.6.100) \\
 &= 0.557 && (22.3.6.101) \\
 &= 0.557 && (22.3.6.102) \\
 &= 0.557 && (22.3.6.103) \\
 &= 0.557 && (22.3.6.104) \\
 &= 0.557 && (22.3.6.105)
 \end{aligned}$$

Note, that $\langle gg|gu \rangle = \langle uu|ug \rangle = 0$ from symmetry considerations, but this can be easily verified. For example,

$$\begin{aligned}
 \langle gg|gu \rangle &= \frac{1}{\sqrt{2+2S}} \frac{1}{\sqrt{(2-2S)^3}} && (22.3.6.106) \\
 &= \frac{1}{\sqrt{2+2S}} \frac{1}{\sqrt{(2-2S)^3}} \frac{1}{(1s_A+1s_B)(1s_A+1s_B)(1s_A-1s_B)(1s_A-1s_B)} && (22.3.6.107) \\
 &= \frac{1}{\sqrt{2+2S}} \frac{1}{(2-2S)^3} \cdot && (22.3.6.108) \\
 &= \frac{(\langle AA|AA \rangle - \langle AA|AB \rangle + \langle AA|BA \rangle - \langle AA|BB \rangle + \langle AB|AA \rangle - \langle AB|AB \rangle + \langle AB|BA \rangle - \langle AB|BB \rangle + \langle BA|AA \rangle - \langle BA|AB \rangle + \langle BA|BA \rangle - \langle BA|BB \rangle + \langle BB|AA \rangle - \langle BB|AB \rangle + \langle BB|BA \rangle - \langle BB|BB \rangle)}{\sqrt{2+2S} (2-2S)^3} && (22.3.6.109) \\
 &= \frac{(0.534)(2.880)((0.625) - (0.426) + (0.426) - (0.323) + (0.426) - (0.504) + (0.323) - (0.426) + \dots \dots + (0.426) - (0.323) + (0.504) - (0.426) + (0.323) - (0.426) + (0.426) - (0.625))}{\sqrt{2+2S} (2-2S)^3} && (22.3.6.110) \\
 &= 0.000 && (22.3.6.111) \\
 &= 0.000 && (22.3.6.112) \\
 &= 0.000 && (22.3.6.113) \\
 &= 0.000 && (22.3.6.114) \\
 &= 0.000 && (22.3.6.115)
 \end{aligned}$$

c. We can now set up the configuration interaction Hamiltonian matrix. The elements are evaluated by using the Slater-Condon rules as shown in the text.

$$\begin{aligned}
 H_{11} &= \langle \sigma_g \alpha \sigma_g \beta | H | \sigma_g \alpha \sigma_g \beta \rangle && (22.3.6.116) \\
 &= 2f_{\sigma_g \sigma_g} + g_{\sigma_g \sigma_g \sigma_g \sigma_g} && (22.3.6.117) \\
 &= 2(-1.184) + 0.564 = -1.804 && (22.3.6.118) \\
 H_{21} &= \langle \sigma_g \alpha \sigma_g \beta | H | \sigma_u \alpha \sigma_u \beta \rangle && (22.3.6.119) \\
 &= g_{\sigma_g \sigma_g \sigma_u \sigma_u} && (22.3.6.120) \\
 &= 0.140 && (22.3.6.121) \\
 H_{22} &= \langle \sigma_u \alpha \sigma_u \beta | H | \sigma_u \alpha \sigma_u \beta \rangle && (22.3.6.122) \\
 &= 2f_{\sigma_u \sigma_u} + g_{\sigma_u \sigma_u \sigma_u \sigma_u} && (22.3.6.123) \\
 &= 2(-0.575) + 0.582 = -0.568 && (22.3.6.124)
 \end{aligned}$$

d. Solving this eigenvalue problem:

$$\begin{aligned}
 \begin{vmatrix} -1.804 - \epsilon & 0.140 \\ 0.140 & -0.568 - \epsilon \end{vmatrix} &= 0 && (22.3.6.125) \\
 (-1.804 - \epsilon)(-0.568 - \epsilon) - (0.140)^2 &= 0 && (22.3.6.126) \\
 1.025 + 1.804\epsilon + 0.568\epsilon + \epsilon^2 - 0.0196 &= 0 && (22.3.6.127) \\
 \epsilon^2 + 2.372\epsilon + 1.005 &= 0 && (22.3.6.128) \\
 \epsilon &= \frac{-2.372 \pm \sqrt{(2.372)^2 - 4(1)(1.005)}}{(2)(1)} && (22.3.6.129) \\
 &= -1.186 \pm 0.634 && (22.3.6.130) \\
 &= -1.820, \text{ and } -0.552. && (22.3.6.131)
 \end{aligned}$$

Solving for the coefficients:

$$\begin{bmatrix} -1,804 - \varepsilon & 0.140 \\ 0.140 & -0.568 - \varepsilon \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (22.3.6.132)$$

For the first eigenvalue this becomes:

$$\begin{bmatrix} -1.804 + 1.820 & 0.140 \\ 0.140 & -0.568 + 1.820 \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (22.3.6.133)$$

$$\begin{bmatrix} 0.016 & 0.140 \\ 0.140 & 1.252 \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (22.3.6.134)$$

$$(0.140)(C_1) + (1.252)(C_2) = 0 \quad (22.3.6.135)$$

$$C_1 = -8.943C_2 \quad (22.3.6.136)$$

$$C_1^2 + C_2^2 = 1 \text{ (from normalization)} \quad (22.3.6.137)$$

$$(-8.943C_2)^2 + C_2^2 = 1 \quad (22.3.6.138)$$

$$80.975C_2^2 = 1 \quad (22.3.6.139)$$

$$C_2 = 0.111, C_1 = -0.994 \quad (22.3.6.140)$$

For the second eigenvalue this becomes:

$$\begin{bmatrix} -1.804 + 0.552 & 0.140 \\ 0.140 & -0.568 + 0.552 \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (22.3.6.141)$$

$$\begin{bmatrix} -1.252 & 0.140 \\ 0.140 & -0.016 \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \quad (22.3.6.142)$$

$$(-1.252)(C_1) + (0.140)(C_2) = 0 \quad (22.3.6.143)$$

$$C_1 = 0.112C_2 \quad (22.3.6.144)$$

$$C_1^2 + C_2^2 = 1 \text{ (from normalization)} \quad (22.3.6.145)$$

$$(0.112C_2)^2 + C_2^2 = 1 \quad (22.3.6.146)$$

$$1.0125C_2^2 = 1 \quad (22.3.6.147)$$

$$C_2 = 0.994, C_1 = 0.111 \quad (22.3.6.148)$$

e. The polarized orbitals, R_{\pm} , are given by:

$$R_{\pm} = \sigma_g \pm \sqrt{\frac{C_2}{C_1}} \sigma_u \quad (22.3.6.149)$$

$$R_{\pm} = \sigma_g \pm \sqrt{\frac{0.111}{0.994}} \sigma_u \quad (22.3.6.150)$$

$$R_{\pm} = \sigma_g \pm 0.334\sigma_u \quad (22.3.6.151)$$

$$R_+ = \sigma_g + 0.334\sigma_u \text{ (left polarized)} \quad (22.3.6.152)$$

$$R_- = \sigma_g - 0.334\sigma_u \text{ (right polarized)} \quad (22.3.6.153)$$



R_+ Left Polarized



R_- Right Polarized

This page titled 22.3.6: vi. Problem Solutions is shared under a CC BY-NC-SA 4.0 license and was authored, remixed, and/or curated by Jack Simons via source content that was edited to the style and standards of the LibreTexts platform.