

7.2: Dimensionless Coordinates

The classical energy (Hamiltonian) is

$$E = \frac{1}{2}m\dot{x}^2 + \frac{1}{2}m\omega^2 x^2 = \frac{1}{2m}p^2 + \frac{1}{2}m\omega^2 x^2 \quad (7.2.1)$$

The quantum Hamiltonian operator is thus

$$\widehat{H} = \frac{1}{2m} \frac{1}{2m} \hat{p}^2 + \frac{1}{2}m\omega^2 x^2 = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + \frac{1}{2}m\omega^2 x^2. \quad (7.2.2)$$

And we thus have to solve Schrödinger's equation

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \phi(x) + \frac{1}{2}m\omega^2 x^2 \phi(x) = E\phi(x) \quad (7.2.3)$$

In order to treat this equation it is better to remove all the physical constants, and go over to dimensionless coordinates

$$y = \sqrt{\frac{m\omega}{\hbar}} x, \quad \epsilon = \frac{E}{\hbar\omega}. \quad (7.2.4)$$

When we substitute these new variables into the Schrödinger equation we get, using

$$\frac{d}{dx} f(y) = \frac{dy}{dx} \frac{d}{dy} f(y) = \sqrt{\frac{m\omega}{\hbar}} \frac{d}{dy} f(y), \quad (7.2.5)$$

that $(\phi(x) = u(y))$

$$-\frac{\hbar^2}{2m} \frac{m\omega}{\hbar} \frac{d^2}{dy^2} u(y) + \frac{1}{2}m\omega^2 \frac{\hbar}{m\omega} y^2 u(y) = \epsilon \hbar\omega u(y) \quad (7.2.6)$$

Cancelling the common factor

$$\hbar\omega \text{ we find} \quad \frac{d^2}{dy^2} u(y) + (2\epsilon - y^2) u(y) = 0 \quad (7.2.7)$$

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