

4.4: Barrier Penetration

In order to understand quantum mechanical tunnelling in fission it makes sense to look at the simplest fission process: the emission of a He nucleus, so called α radiation (Figure 4.4.1).

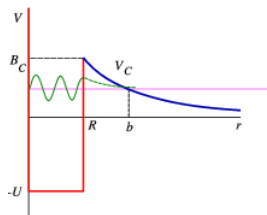


Figure 4.4.1: The potential energy for alpha decay

Suppose there exists an α particle inside a nucleus at an (unbound) energy > 0 . Since it isn't bound, why doesn't it decay immediately? This must be tunnelling. In Figure 4.4.1) we have once again shown the nuclear binding potential as a square well (red curve), but we have included the Coulomb tail (blue curve),

$$V_{\text{Coulomb}}(r) = \frac{(Z-2)2e^2}{4\pi\epsilon_0 r}.$$

The height of the barrier is exactly the coulomb potential at the boundary, which is the nuclear radius, $R_C = 1.2A^{1/3} \text{ fm}$, and thus $B_C = 2.4(Z-2)A^{-1/3}$. The decay probability across a barrier can be given by the simple integral expression $P = e^{-2\gamma}$, with

$$\begin{aligned} \gamma &= \frac{(2\mu_\alpha)^{1/2}}{\hbar} \int_{R_C}^b [V(r) - E_\alpha]^{1/2} dr \\ &= \frac{(2\mu_\alpha)^{1/2}}{\hbar} \int_{R_C}^b \left[\frac{2(Z-2)e^2}{4\pi\epsilon_0 r} - E_\alpha \right]^{1/2} dr \\ &= \frac{2(Z-2)e^2}{2\pi\epsilon_0 \hbar v} [\arccos(E_\alpha/B_C) - (E_\alpha/B_C)(1 - E_\alpha/B_C)], \end{aligned}$$

where v is the velocity associated with E_α . In the limit that $B_C \gg E_\alpha$ we find

$$P = \exp \left[-\frac{2(Z-2)e^2}{2\epsilon_0 \hbar v} \right].$$

This shows how sensitive the probability is to Z and v !

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