

12.13: Forbidden Transitions

Atomic transitions which are forbidden by the electric dipole selection rules ([\[e13.133\]](#)) and ([\[e13.134\]](#)) are unsurprisingly known as *forbidden transitions*. It is clear from the analysis in Section [1.8](#) that a forbidden transition is one for which the matrix element $\langle f | \epsilon \cdot \mathbf{p} | i \rangle$ is zero. However, this matrix element is only an approximation to the true matrix element for radiative transitions, which takes the form $\langle f | \epsilon \cdot \mathbf{p} \exp(i \mathbf{k} \cdot \mathbf{r}) | i \rangle$. Expanding $\exp(i \mathbf{k} \cdot \mathbf{r})$, and keeping the first two terms, the matrix element for a forbidden transition becomes

$$\langle f | \epsilon \cdot \mathbf{p} \exp(i \mathbf{k} \cdot \mathbf{r}) | i \rangle \simeq i \langle f | (\epsilon \cdot \mathbf{p}) (\mathbf{k} \cdot \mathbf{r}) | i \rangle. \quad (12.13.1)$$

Hence, if the residual matrix element on the right-hand side of the previous expression is non-zero then a “forbidden” transition can take place, albeit at a much reduced rate. In fact, in Section [1.9](#), we calculated that the typical rate of an electric dipole transition is

$$w_{i \rightarrow f} \sim \alpha^3 \omega_{if}. \quad (12.13.2)$$

Because the transition rate is proportional to the square of the radiative matrix element, it is clear that the transition rate for a forbidden transition enabled by the residual matrix element ([\[e13.146\]](#)) is smaller than that of an electric dipole transition by a factor $(kr)^2$. Estimating r as the Bohr radius, and k as the wavenumber of a typical spectral line of hydrogen, it is easily demonstrated that

$$w_{i \rightarrow f} \sim \alpha^5 \omega_{if} \quad (12.13.3)$$

for such a transition. Of course, there are some transitions (in particular, the $2S \rightarrow 1S$ transition) for which the true radiative matrix element $\langle f | \epsilon \cdot \mathbf{p} \exp(i \mathbf{k} \cdot \mathbf{r}) | i \rangle$ is zero. Such transitions are absolutely forbidden.

Finally, it is fairly obvious that excited states which decay via forbidden transitions have much longer life-times than those which decay via electric dipole transitions. Because the natural width of a spectral line is inversely proportional to the life-time of the associated decaying state, it follows that spectral lines associated with forbidden transitions are generally much sharper than those associated with electric dipole transitions.

Contributors and Attributions

- [Richard Fitzpatrick](#) (Professor of Physics, The University of Texas at Austin)

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