

9.3: Eigenstates of S_z and S^2

Because the operators S_z and S^2 commute, they must possess simultaneous eigenstates. (See Section [smeas].) Let these eigenstates take the form [see Equations ([e8.29]) and ([e8.30])]:

$$\begin{aligned} S_z \chi_{s,m_s} &= m_s \hbar \chi_{s,m_s}, \\ S^2 \chi_{s,m_s} &= s(s+1) \hbar^2 \chi_{s,m_s}. \end{aligned}$$

Now, it is easily demonstrated, from the commutation relations ([e10.9]) and ([e10.10]), that

$$S_z (S_+ \chi_{s,m_s}) = (m_s + 1) \hbar (S_+ \chi_{s,m_s}), \quad (9.3.1)$$

and

$$S_z (S_- \chi_{s,m_s}) = (m_s - 1) \hbar (S_- \chi_{s,m_s}). \quad (9.3.2)$$

Thus, S_+ and S_- are indeed the raising and lowering operators, respectively, for spin angular momentum. (See Section [seian].) The eigenstates of S_z and S^2 are assumed to be orthonormal: that is,

$$\chi_{s,m_s}^\dagger \chi_{s',m'_s} = \delta_{ss'} \delta_{m_s m'_s}. \quad (9.3.3)$$

Consider the wavefunction $\chi = S_+ \chi_{s,m_s}$. Because we know, from Equation ([e10.11]), that $\chi^\dagger \chi \geq 0$, it follows that

$$(S_+ \chi_{s,m_s})^\dagger (S_+ \chi_{s,m_s}) = \chi_{s,m_s}^\dagger S_+^\dagger S_+ \chi_{s,m_s} = \chi_{s,m_s}^\dagger S_- S_+ \chi_{s,m_s} \geq 0, \quad (9.3.4)$$

where use has been made of Equation ([e10.7]). Equations ([e10.8]), ([e10.16]), ([e10.17]), and ([e10.20]) yield

$$s(s+1) \geq m_s(m_s+1). \quad (9.3.5)$$

Likewise, if $\chi = S_- \chi_{s,m_s}$ then we obtain

$$s(s+1) \geq m_s(m_s-1). \quad (9.3.6)$$

Assuming that $s \geq 0$, the previous two inequalities imply that

$$-s \leq m_s \leq s. \quad (9.3.7)$$

Hence, at fixed s , there is both a maximum and a minimum possible value that m_s can take.

Let $m_{s \min}$ be the minimum possible value of m_s . It follows that (see Section [slsq])

$$S_- \chi_{s,m_{s \min}} = 0. \quad (9.3.8)$$

Now, from Equation ([e10.7a]),

$$S^2 = S_+ S_- + S_z^2 - \hbar S_z. \quad (9.3.9)$$

Hence,

$$S^2 \chi_{s,m_{s \min}} = (S_+ S_- + S_z^2 - \hbar S_z) \chi_{s,m_{s \min}}, \quad (9.3.10)$$

giving

$$s(s+1) = m_{s \min}(m_{s \min} - 1). \quad (9.3.11)$$

Assuming that $m_{s \min} < 0$, this equation yields

$$m_{s \min} = -s. \quad (9.3.12)$$

Likewise, it is easily demonstrated that

$$m_{s \max} = +s. \quad (9.3.13)$$

Moreover,

$$S_- \chi_{s,-s} = S_+ \chi_{s,s} = 0. \quad (9.3.14)$$

Now, the raising operator S_+ , acting upon $\chi_{s,-s}$, converts it into some multiple of $\chi_{s,-s+1}$. Employing the raising operator a second time, we obtain a multiple of $\chi_{s,-s+2}$. However, this process cannot continue indefinitely, because there is a maximum possible value of m_s . Indeed, after acting upon $\chi_{s,-s}$ a sufficient number of times with the raising operator S_+ , we must obtain a multiple of $\chi_{s,s}$, so that employing the raising operator one more time leads to the null state. [See Equation ([e10.31]).] If this is not the case then we will inevitably obtain eigenstates of S_z corresponding to $m_s > s$, which we have already demonstrated is impossible.

It follows, from the previous argument, that

$$m_{s \max} - m_{s \min} = 2s = k, \quad (9.3.15)$$

where k is a positive integer. Hence, the quantum number s can either take positive integer or positive half-integer values. Up to now, our analysis has been very similar to that which we used earlier to investigate orbital angular momentum. (See Section [sorb].) Recall, that for orbital angular momentum the quantum number m , which is analogous to m_s , is restricted to take integer values. (See Section [slz].) This implies that the quantum number l , which is analogous to s , is also restricted to take integer values. However, the origin of these restrictions is the representation of the orbital angular momentum operators as differential operators in real space. (See Section [s8.3].) There is no equivalent representation of the corresponding spin angular momentum operators. Hence, we conclude that there is no reason why the quantum number s cannot take half-integer, as well as integer, values.

In 1940, Wolfgang Pauli proved the so-called *spin-statistics theorem* using relativistic quantum mechanics. According to this theorem, all fermions possess half-integer spin (i.e., a half-integer value of s), whereas all bosons possess integer spin (i.e., an integer value of s). In fact, all presently known fermions, including electrons and protons, possess *spin one-half*. In other words, electrons and protons are characterized by $s = 1/2$ and $m_s = \pm 1/2$.

Contributors and Attributions

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