

## 9.1: Spin Operators

Because spin is a type of angular momentum, it is reasonable to suppose that it possesses similar properties to orbital angular momentum. Thus, by analogy with Section [s8.2], we would expect to be able to define three operators— $S_x$ ,  $S_y$ , and  $S_z$ —that represent the three Cartesian components of spin angular momentum. Moreover, it is plausible that these operators possess analogous commutation relations to the three corresponding orbital angular momentum operators,  $L_x$ ,  $L_y$ , and  $L_z$ . [See Equations ([e8.6])–([e8.8]).] In other words,

$$\begin{aligned}[S_x, S_y] &= i\hbar S_z, \\ [S_y, S_z] &= i\hbar S_x, \\ [S_z, S_x] &= i\hbar S_y.\end{aligned}$$

We can represent the magnitude squared of the spin angular momentum vector by the operator

$$S^2 = S_x^2 + S_y^2 + S_z^2. \quad (9.1.1)$$

By analogy with the analysis in Section [s8.2], it is easily demonstrated that

$$[S^2, S_x] = [S^2, S_y] = [S^2, S_z] = 0. \quad (9.1.2)$$

We thus conclude (see Section [smeas]) that we can simultaneously measure the magnitude squared of the spin angular momentum vector, together with, at most, one Cartesian component. By convention, we shall always choose to measure the  $z$ -component,  $S_z$ .

By analogy with Equation ([e8.13]), we can define raising and lowering operators for spin angular momentum:

$$S_{\pm} = S_x \pm i S_y. \quad (9.1.3)$$

If  $S_x$ ,  $S_y$ , and  $S_z$  are Hermitian operators, as must be the case if they are to represent physical quantities, then  $S_{\pm}$  are the Hermitian conjugates of one another: that is,

$$(S_{\pm})^{\dagger} = S_{\mp}. \quad (9.1.4)$$

Finally, by analogy with Section [s8.2], it is easily demonstrated that

$$\begin{aligned}S_+ S_- &= S^2 - S_z^2 + \hbar S_z, \\ S_- S_+ &= S^2 - S_z^2 - \hbar S_z, \\ [S_+, S_z] &= -\hbar S_+, \\ [S_-, S_z] &= +\hbar S_-.\end{aligned}$$

### Contributors and Attributions

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