

12.E: Time-Dependent Perturbation Theory (Exercises)

1. Consider the two-state system examined in Section 1.3.[ex8.1] Suppose that

$$\begin{aligned}\langle 1 | H_1 | 1 \rangle &= e_{11}, \\ \langle 2 | H_1 | 2 \rangle &= e_{22}, \\ \langle 1 | H_1 | 2 \rangle &= \langle 2 | H_1 | 1 \rangle^* = \frac{1}{2} \gamma \hbar \exp(i \omega t),\end{aligned}$$

where e_{11} , e_{22} , γ , and ω are real. Show that

$$\begin{aligned}i \frac{d\hat{c}_1}{dt} &= \frac{\gamma}{2} \exp[+i(\omega - \hat{\omega}_{21})t] \hat{c}_2, \\ i \frac{d\hat{c}_2}{dt} &= \frac{\gamma}{2} \exp[-i(\omega - \hat{\omega}_{21})t] \hat{c}_1,\end{aligned}$$

where $\hat{c}_1 = c_1 \exp(i e_{11} t / \hbar)$, $\hat{c}_2 = c_2 \exp(i e_{22} t / \hbar)$, and

$$\hat{\omega}_{21} = \frac{E_2 + e_{22} - E_1 - e_{11}}{\hbar}. \quad (12.E.1)$$

Hence, deduce that if the system is definitely in state 1 at time $t = 0$ then the probability of finding it in state 2 at some subsequent time, t , is

$$P_2(t) = \frac{\gamma^2}{\gamma^2 + (\omega - \hat{\omega}_{21})^2} \sin^2 \left(\left[\gamma^2 + (\omega - \hat{\omega}_{21})^2 \right]^{1/2} \frac{t}{2} \right). \quad (12.E.2)$$

2. Consider an atomic nucleus of spin- s and gyromagnetic ratio g placed in the magnetic field

$$\mathbf{B} = B_0 \mathbf{e}_z + B_1 [\cos(\omega t) \mathbf{e}_x - \sin(\omega t) \mathbf{e}_y], \quad (12.E.3)$$

where $B_1 \ll B_0$. Let $\chi_{s,m}$ be a properly normalized simultaneous eigenstate of S^2 and S_z , where \mathbf{S} is the nuclear spin. Thus, $S^2 \chi_{s,m} = s(s+1) \hbar^2 \chi_{s,m}$ and $S_z \chi_{s,m} = m \hbar \chi_{s,m}$, where $-s \leq m \leq s$. Furthermore, the instantaneous nuclear spin state is written

$$\chi = \sum_{m=-s,s} c_m(r) \chi_{s,m}, \quad (12.E.4)$$

where $\sum_{m=-s,s} |c_m|^2 = 1$.

1. Demonstrate that

$$\begin{aligned}\frac{dc_m}{dt} &= \frac{i\gamma}{2} \left([s(s+1) - m(m-1)]^{1/2} e^{i(\omega - \omega_0)t} c_{m-1} \right. \\ &\quad \left. + [s(s+1) - m(m+1)]^{1/2} e^{-i(\omega - \omega_0)t} c_{m+1} \right)\end{aligned}$$

for $-s \leq m \leq s$, where $\omega_0 = g \mu_N B_0 / \hbar$, $\gamma = g \mu_N B_1 / \hbar$, and $\mu_N = e \hbar / (2 m_p)$.

2. Consider the case $s = 1/2$. Demonstrate that if $\omega = \omega_0$ and $c_{1/2}(0) = 1$ then

$$c_{1/2}(t) = \cos(\gamma t / 2), \quad c_{-1/2}(t) = i \sin(\gamma t / 2).$$

3. Consider the case $s = 1$. Demonstrate that if $\omega = \omega_0$ and $c_1(0) = 1$ then

$$\begin{aligned}c_1(t) &= \cos^2(\gamma t / 2), \\ c_0(t) &= i\sqrt{2} \cos(\gamma t / 2) \sin(\gamma t / 2), \\ c_{-1}(t) &= -\sin^2(\gamma t / 2).\end{aligned}$$

4. Consider the case $s = 3/2$. Demonstrate that if $\omega = \omega_0$ and $c_{3/2}(0) = 1$ then

$$\begin{aligned}c_{3/2}(t) &= \cos^3(\gamma t/2), \\c_{1/2}(t) &= i\sqrt{3} \cos(\gamma t/2) \sin^2(\gamma t/2), \\c_{-1/2}(t) &= -\sqrt{3} \cos^2(\gamma t/2) \sin(\gamma t/2), \\c_{-3/2}(t) &= -i \sin^3(\gamma t/2).\end{aligned}$$

3. Demonstrate that a spontaneous transition between two atomic states of zero orbital angular momentum is absolutely forbidden. (Actually, a spontaneous transition between two zero orbital angular momentum states is possible via the simultaneous emission of two photons, but takes place at a very slow rate .)

Contributors and Attributions

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