

## 9.1: The Program of Renormalization

A statistical mechanical system is defined by a set of *degrees of freedom* and by a set of *coupling constants*  $\{K_\alpha\}$ . The degrees of freedom can be discrete, such as Ising spins  $\sigma_i$ , or continuous, such as a field  $\phi(\mathbf{r})$ . Additionally, each such system possesses a *microscopic length scale*  $\ell$ . For discrete, lattice-based systems, this length scale is simply the lattice spacing:  $\ell = a$ . For continuous systems, we can define a microscopic length scale by imposing a *cutoff*  $\Lambda$  on the wavevectors we integrate over in all Fourier transforms. That is, we replace

$$\int \frac{d^d k}{(2\pi)^d} F(\mathbf{k}) \longrightarrow \int \frac{d^d k}{(2\pi)^d} F(\mathbf{k}) g_\Lambda(\mathbf{k}), \quad (9.1.1)$$

where  $F(\mathbf{k})$  is any function and  $g_\Lambda(\mathbf{k})$  is the cutoff function. The simplest such case to imagine is a sharp cutoff which is isotropic in wavevector,  $g_\Lambda(\mathbf{k}) = \Theta(\Lambda - |\mathbf{k}|)$ . Other cutoff schemes, however, are possible, including ‘soft cutoffs’ where  $g_\Lambda(\mathbf{k})$  is smooth. The microscopic length scale is then  $\ell \sim \Lambda^{-1}$ , which is the smallest distance in real space over which the system can independently fluctuate.

The idea behind renormalization is that we can successively winnow degrees of freedom from a system in some exact or approximate way, and in so doing we generate a new version of the system, at a different length scale  $\ell' > \ell$ , and with different couplings  $\{K'_\alpha\}$ . We then iterate this procedure. The result is a set of equations which tell us how the couplings behave under a change of the microscopic length scale. As we shall see, the fixed points of this procedure – where couplings do not change under a change of length scale – are critical points. Such a fixed point is defined by a set of couplings  $\{K_\alpha^*\}$ .

If we denote by  $\mathcal{R}_b$  the renormalization procedure

$$\mathcal{R}_b(\ell, \{K_\alpha\}) = (\ell', \{K'_\alpha\}), \quad (9.1.2)$$

where  $\ell' = b\ell$ , then we have the composition law  $\mathcal{R}_b \mathcal{R}_{b'} = \mathcal{R}_{b+b'}$ . The set of transformations  $\{\mathcal{R}_b\}$  is collectively referred to as the *renormalization group* (RG) because of this mathematical structure. It is somewhat of a misnomer, however, since the transformations are only defined for  $b \geq 1$ , which means that there is no inverse operation, and hence no true group structure<sup>1</sup>. Nevertheless we shall use the RG terminology because it has become universally accepted in the literature.

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