

## 10.6: Example- Copper

The real and imaginary parts  $(n_\theta + i\kappa_\theta) = \sqrt{\epsilon_r - \sin^2 \theta}$  have been plotted in Figure (10.6.8) as a function of the angle of incidence,  $\theta$ , for room temperature copper and for a wavelength of  $\lambda = 0.5145$  microns (see Table(10.1)). As can be seen from the figure, the angular dependence of the indices  $n_\theta$ ,  $\kappa_\theta$  is not very pronounced. For a lossy material such as copper that has a complex dielectric constant the reflectivity,  $E_R/E_0$ , is complex; that is, the phase shift between the incident wave and reflected wave electric vectors is neither  $0^\circ$  (in phase) nor  $180^\circ$  (out of phase). The real and imaginary parts of the reflectivity have been plotted in Figure (10.6.9) as a function of the angle of incidence for S-polarized 0.5145 micron light incident on room temperature copper; the absolute value of the reflectivity has been plotted in Figure (10.6.10).

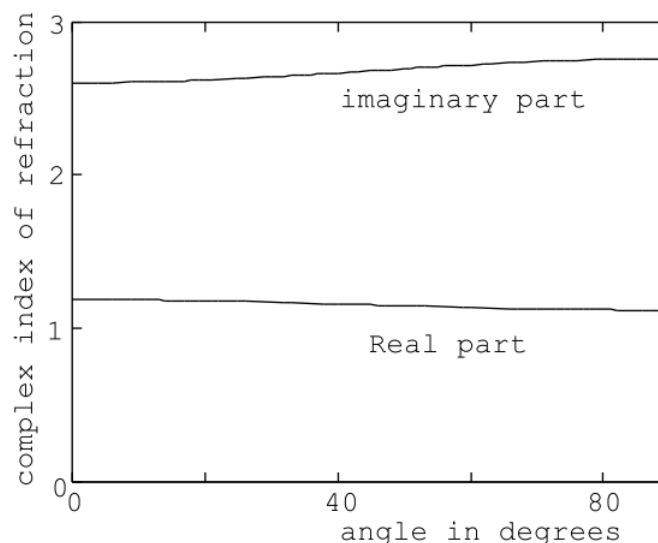


Figure 10.6.8: The dependence of the complex index of refraction,  $n_\theta + i\kappa_\theta$ , upon angle of incidence of the incident wave calculated for copper at room temperature and for an incident wavelength of  $\lambda = 0.5145 \mu\text{m}$ . The normal component of the wave-vector in copper is given by  $k_z = \left(\frac{\omega}{c}\right)(n_\theta + i\kappa_\theta)$ . At this wavelength the relative dielectric constant for copper is  $\epsilon_r = (-5.34 + i6.19)$ ,  $n=1.19$ , and  $\kappa=2.60$ .

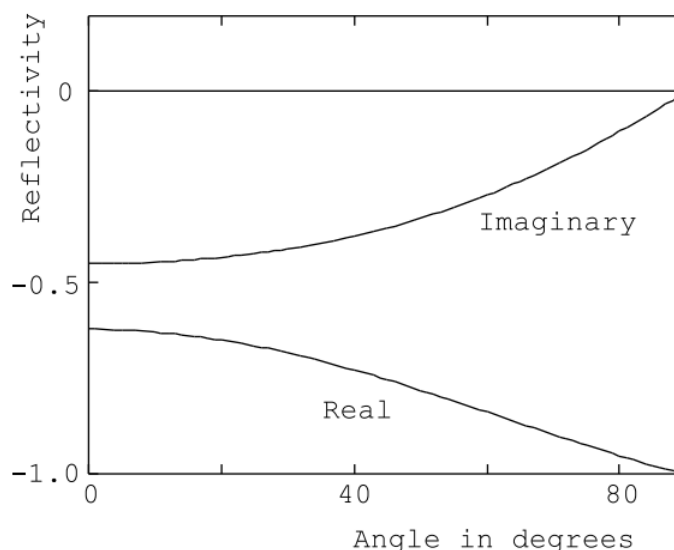


Figure 10.6.9: The real and imaginary part of the reflectivity of copper,  $E_R/E_0$ , as a function of angle of incidence for a wavelength  $\lambda = 0.5145 \mu\text{m}$  and S-polarized radiation. Copper at room temperature;  $\epsilon_r = (-5.34 + i6.19)$ ,  $n=1.19$ , and  $\kappa=2.60$ .

Similarly, the real and the imaginary parts of the ratio  $H_R/H_0$  have been plotted in Figure (10.6.11) as a function of the angle of incidence for P-polarized 0.5145 micron light incident on copper; the absolute value of this ratio is shown in Figure (10.6.12). The

reflection coefficient for P-polarized radiation is given by  $R_p = E_R/E_0$  but this is very closely related to the ratio  $H_R/H_0$  because  $E_0 = Z_0 H_0$  and  $E_R = -Z_0 H_R$ , where  $Z_0 = 377$  Ohms, the impedance of free space. Notice that the real part of the reflectivity for P-polarized light vanishes at an angle of incidence of approximately  $69^\circ$ ; the phase of the reflected light at that angle is shifted by  $90^\circ$  relative to the incident light. The phase shift between reflected and incident light is much less pronounced for S-polarized light; approximately  $15^\circ$  for an angle of incidence of  $69^\circ$ .

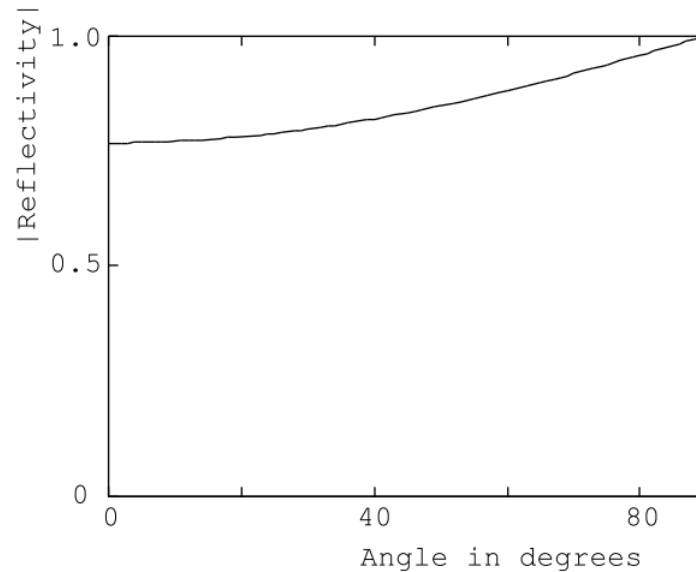


Figure 10.6.10: The absolute value of the reflectivity of copper,  $|E_R/E_0|$ , as a function of angle of incidence for a wavelength  $\lambda = 0.5145 \mu\text{m}$  and S-polarized radiation. Copper at room temperature;  $\epsilon_r = (-5.34 + i6.19)$ ,  $n = 1.19$ , and  $\kappa = 2.60$ .

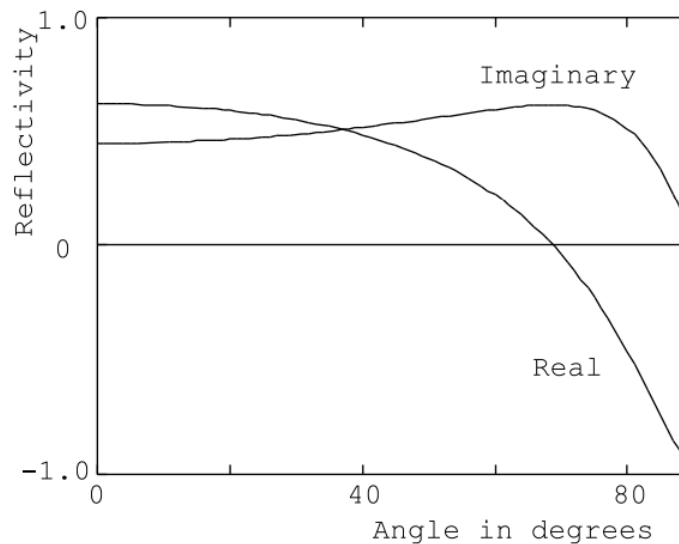


Figure 10.6.11: The real and imaginary parts of the complex ratio  $H_R/H_0$  for copper as a function of angle of incidence for a wavelength  $\lambda = 0.5145 \mu\text{m}$  and for P-polarized radiation. Copper at room temperature;  $\epsilon_r = (-5.34 + i6.19)$ ,  $n = 1.19$ , and  $\kappa = 2.60$ .

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