

## 19.7: Finding the Eigenvalues

The eigenvalues are found by operating on the eigenvector we just found with the matrix, meaning the  $N$  dimensional generalization of

$$-m\Omega^2 \begin{pmatrix} 1 \\ e^{ik_n a} \\ e^{ik_n 2a} \\ e^{ik_n 3a} \end{pmatrix} = \begin{pmatrix} -2\kappa & \kappa & 0 & \kappa \\ \kappa & -2\kappa & \kappa & 0 \\ 0 & \kappa & -2\kappa & \kappa \\ \kappa & 0 & \kappa & -2\kappa \end{pmatrix} \begin{pmatrix} 1 \\ e^{ik_n a} \\ e^{ik_n 2a} \\ e^{ik_n 3a} \end{pmatrix} \quad (19.7.1)$$

Applying the matrix to the column vector

$$\left(1, e^{ik_n a}, e^{2ik_n a}, e^{3ik_n a}, \dots, e^{i(N-1)k_n a}\right)^T \quad (19.7.2)$$

, and cancelling out the common  $e^{ik_n a}$  factor, we have

$$-m\Omega_n^2 = \kappa (e^{ik_n a} + e^{-ik_n a} - 2) \quad (19.7.3)$$

(Of course, this same result comes from every row.)

The complete set of eigenvalues is given by inserting in the above expression

$$k_n = 2\pi n / Na, \quad n = 0, 1, 2, \dots, N-1 \text{ so } e^{ik_n a} = e^{2\pi i n / N} \quad (19.7.4)$$

so  $n = 0$  is displacement of the system as a whole, as is  $n = N$ .

Wavenumber values  $k_n$  beyond  $n = N$  repeats the eigenstates we already have, since

$$e^{ik_{N+n} a} = e^{i \frac{2\pi(N+n)a}{Na}} = e^{2\pi i} e^{2\pi i n / N} = e^{2\pi i n / N} = e^{ik_n a} \quad (19.7.5)$$

$k$  are restricted to

$$\Omega_n = 2\sqrt{\frac{\kappa}{m}} \sin\left(\frac{k_n a}{2}\right) = 2\sqrt{\frac{\kappa}{m}} \sin\left(\frac{n\pi}{N}\right) \quad (19.7.6)$$

$$0 \leq k < 2\pi/a \quad (19.7.7)$$

or equivalently

$$-\pi/a < k \leq \pi/a \quad (19.7.8)$$

The eigenvalue equation is

$$\Omega_n^2 = 2(\kappa/m) (1 - \cos k_n a) \quad (19.7.9)$$

or

$$\Omega_n = 2\sqrt{\frac{\kappa}{m}} \sin\left(\frac{k_n a}{2}\right) = 2\sqrt{\frac{\kappa}{m}} \sin\left(\frac{n\pi}{N}\right) \quad (19.7.10)$$

To see the dynamics of this eigenstate

$$\left(1, e^{ik_n a}, e^{2ik_n a}, e^{3ik_n a}, \dots, e^{ik_n(N-1)a}\right) \quad (19.7.11)$$

, we need to multiply by the time dependence  $e^{i\Omega_n t}$ , then finally take the real part of the solution:

$$(\cos \Omega_n t, \cos(k_n a + \Omega_n t), \cos(2k_n a + \Omega_n t), \cos(3k_n a + \Omega_n t), \dots, \cos((N-1)k_n a + \Omega_n t)) \quad (19.7.12)$$

Notice that in the continuum limit, meaning large  $N$  and small  $a$ , the atom displacement as a function of position has the form  $\cos(kx + \Omega t)$  in other words we're looking at a sinusoidal wave disturbance with wavenumber  $k_n$  here.

Now,  $-k_n$  is also a solution, but that is the same as  $n' = N - n$  so one must be careful not to overcount. The two frequencies  $\pm\Omega_n$  correspond to waves going in opposite directions.

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