

## 21.2: Finding the Effective Potential Generated by the Oscillating Force

As stated above, our system is a particle of mass  $m$  moving in one dimension in a time-independent potential  $V(x)$  and subject to a rapidly oscillating force  $f = f_1 \cos \omega t + f_2 \sin \omega t$ .

The oscillation's strength and frequency are such that the particle only moves a small distance in  $V(x)$  during one cycle, and the oscillation is much faster than any oscillation possible in the potential alone.

The equation of motion is

$$m\ddot{x} = -dV/dx + f \quad (21.2.1)$$

The particle will follow a path

$$x(t) = X(t) + \xi(t) \quad (21.2.2)$$

where  $\xi(t)$  describes rapid oscillations about a smooth path  $X(t)$ , and the average value  $\overline{\xi(t)}$  of  $\xi(t)$  over a period  $2\pi/\omega$  is zero.

Expanding to first order in  $\xi$ ,

$$m\ddot{X} + m\ddot{\xi} = -\frac{dV}{dx} - \xi \frac{d^2V}{dx^2} + f(X, t) + \xi \frac{\partial f}{\partial X} \quad (21.2.3)$$

This equation has smooth terms and rapidly oscillating terms on both sides, and we can equate them separately. The leading oscillating terms are

$$m\ddot{\xi} = f(X, t) \quad (21.2.4)$$

We've dropped the terms on the right of order  $\xi$ , but kept  $\ddot{\xi}$ , because  $\ddot{\xi} \sim \omega^2 \xi \gg \xi$ .

So to leading order in the rapid oscillation,

$$\xi = -f/m\omega^2 \quad (21.2.5)$$

Now, averaging the full equation of motion with respect to time (smoothing out the jiggle, matching the slow-moving terms), the  $m\ddot{\xi}$  on the left and the  $f(X, t)$  on the right both disappear (but cancel each other anyway), the  $\xi d^2V/dx^2$  term averages to zero on the assumption that the variation of  $d^2V/dx^2$  over a cycle of the fast oscillation is negligible, but we cannot drop the average

$$\overline{\xi \frac{\partial f}{\partial X}} = -\frac{1}{m\omega^2} \overline{f \frac{\partial f}{\partial X}} = -\frac{1}{m\omega^2} \nabla_X \overline{f^2} \quad (21.2.6)$$

Incorporating this nonzero term, we have an equation of "slow motion"

$$m\ddot{X} = -dV_{\text{eff}}/dX \quad (21.2.7)$$

where, using  $|\dot{\xi}| = |f|/m\omega$ ,

$$V_{\text{eff}} = V + \overline{f^2}/2m\omega^2 = V + \frac{1}{2} m \overline{\dot{\xi}^2} \quad (21.2.8)$$

The effective potential is the original plus a term proportional to the kinetic energy of the oscillation.

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