

22.6: Frequency Multiples

The above analysis is for frequencies not very far from ω_0 . But nonlinear terms can cause resonance to occur at frequencies which are rational multiples of ω_0 . Landau shows that a small $\frac{1}{3}\alpha x^3$ in the potential (so an additional force αx^2 in the equation of motion) can generate a resonance near $\gamma = \frac{1}{2}\omega_0$. We've only considered a quartic addition to the potential, $\frac{1}{4}\beta x^4$, a force βx^3 , we can show that gives a resonance near $\gamma = \frac{1}{3}\omega_0$, and presumably this is the small bump near the beginning of the curves above for large driving strength.

We have $\ddot{x} + 2\lambda\dot{x} + \omega_0^2 x = (f/m) \cos \gamma t - \beta x^3$

We'll write $x = x^{(0)} + x^{(1)} + \dots$

Let's define $x^{(0)}$ by

$$\ddot{x}^{(0)} + 2\lambda\dot{x}^{(0)} + \omega_0^2 x^{(0)} = (f/m) \cos \gamma t \quad (22.6.1)$$

So $x^{(0)} = b \cos(\gamma t + \delta)$. Then

$$\begin{aligned} \ddot{x}^{(1)} + 2\lambda\dot{x}^{(1)} + \omega_0^2 x^{(1)} &= -\beta \left(x^{(0)}\right)^3 \\ &= -\beta b^3 \cos^3(\gamma t + \delta) \\ &= -\beta b^3 \left[\frac{3}{4} \cos(\gamma t + \delta) + \frac{1}{4} \cos(3\gamma t + \delta) \right] \end{aligned}$$

Then, for $\gamma = \frac{1}{3}\omega_0$, the second term, $-\beta b^3 \frac{1}{4} \cos(3\gamma t + \delta) = -\beta b^3 \frac{1}{4} \cos(\omega_0 t + \delta)$, will have a resonant response, although it is proportional to the (small) amplitude cubed. Similar arguments work for other fractional frequencies.

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