

# SPIRAL PHYSICS: ALGEBRA BASED



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Spiral Physics: Algebra Based

Paul D'Alessandris



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## 01. Introduction

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### Model building in physics

Physicists do not study the real world. Physicists build models of the physical world, and then study these models in the hopes of gaining insight into how the world operates. This is the same thing that engineers do – they make assumptions to help simplify a complex problem.

What you will do throughout this course is build models, of increasing complexity, of the real world and by closely examining these models you may gain insight into how the world operates. Some of the early models you will examine will be obviously limited, but keep in mind that even the most advanced physicists are merely model-builders, and the models they study are as obviously superficial to them as the models we will study are to us.

Model building is necessary because of the overwhelming complexity of the real world. To attempt to study a real phenomenon, with all of its many details intact, is extremely difficult. Moreover, models often allow you to focus on the important aspects of a phenomenon, without the distracting details.

For example, a model of reality that everyone is familiar with is the road map. Imagine trying to drive to a strange address across town using a road map that included every driveway and alleyway! Although these details do exist, a model that tried to encompass all of these details would be less useful than one in which everyone's driveway was omitted. Thus, it is possible to omit detail, to be a poorer reflection of reality, yet to be a better, more effective, and more useful model. A useful model for driving across town would ignore driveways but probably include most, if not all, streets. However, if your task was to drive across the state, not only should the driveways be omitted, but so should the vast majority of side streets; probably state and federal highways should be the only roads on the map. Thus, a good model is closely tied to the task at hand. What can be a very useful model for one task can be useless for another.

Thus, when we build models where the effects of friction are neglected, or the shape of an object is ignored, it is not the case that this is a deficient model of the situation. It may well be the case that if these details were included, some important features of the scenario would be masked by the complexity. Simplifications made in constructing models of reality are not always limitations to the usefulness of the model, often they are the key to building a useful and productive model.

### Units

In this course we will exclusively use the *International System of Units (SI Units)*. These are also more commonly referred to as the metric system. In this system, all times are measured in seconds (s), all positions in meters (m), and all masses in kilograms (kg). For the sake of clarity, I will not include units during every step of a calculation, but it can safely be assumed that these standard units are in use throughout all calculations.

As many of you are engineering majors, we will specify all answers in engineering notation. Engineering notation is a version of scientific notation in which the power of ten is always a multiple of three (e.g, 10,000 N is written as  $10 \times 10^3$  N or 10 kN). Whenever possible it is preferable to use the SI prefixes in lieu of writing the powers of 10. Common SI prefixes are shown below:

$10^{12}$	tera	T	Trillion	1,000,000,000,000
$10^9$	giga	G	Billion	1,000,000,000
$10^6$	mega	M	Million	1,000,000
$10^3$	kilo	k	Thousand	1,000
$10^0$	--	--	One	1
$10^{-3}$	milli	m	Thousandth	0.001
$10^{-6}$	micro	$\mu$	Millionth	0.000001

$10^{-9}$	nano	n	Billionth	0.000000001
$10^{-12}$	pico	p	Trillionth	0.000000000001

## Active reading

If you do not currently have a pen or pencil in your hand, pick one up.

As you read this text, you should be writing down explanatory notes to yourself, questions to be asked in class, and any flashes of insight you may have. Don't be afraid to write in the text. I promise that actively wrestling with the ideas on the page rather than passively reading the words on the page will make a huge difference in your understanding. If you find yourself reading page after page of the text without spontaneously thinking of questions, either you are not really digesting the material or you should be in a more advanced class.

In addition, the concepts and principles of physics are complex, even the ones that appear to be simple. (If the principles really were simple, it would not have taken humankind thousands of years to understand the motion of a simple falling object!) They will become clear to you only after careful study. With this in mind, the text is not meant to be read (and written in) once. It should be re-read (and re-written in!) as you work through the various activities included. Hopefully, as you complete the activities, the concepts will come into better focus.

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## 02. Model Specifics

### 1. Model Specifics

1. The object moves along a perfectly straight line
2. The object is acted on by constant forces
3. The object's size and shape are unimportant
4. The object is classical

### Model Specifics

For our first pass through the study of mechanics, we will make a set of simplifying assumptions leading to a model of reality termed the one-dimensional, constant-force, particle model. We will assume that:

#### The object moves along a perfectly straight line

Thus, if we are studying the motion of a car we will assume that the road is perfectly smooth with no bumps, the road is perfectly straight with no curves, and the driver is perfectly steady at the wheel with no swaying to the right or left.

#### The object is acted on by constant forces

For example, when a basketball bounces off the ground the force of the ground on the basketball is:

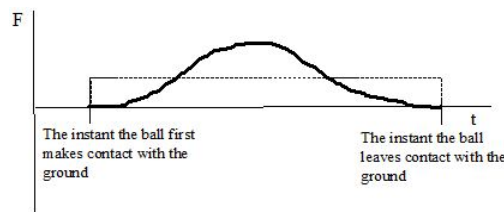
- zero before the basketball strikes the ground,
- non-zero and increasing in magnitude as the basketball deforms in its contact with the ground,
- non-zero and decreasing in magnitude as the basketball springs back into shape while still in contact with the ground,
- and once again zero as the ball leaves contact with the ground.

In reality, the force between the basketball and the ground is quite complicated.

However, in our model we will approximate the force on the basketball due to the ground as:

- zero before the basketball strikes the ground,
- non-zero and constant in magnitude the entire time the ball is in contact with the ground,
- and once again zero as the ball leaves contact with the ground.

The difference between reality and our model is illustrated by the graph below. The dark line represents the actual force between the ground and the ball, the dashed line represents how we will approximate the force in our current model.



In addition, you will soon learn that the total force acting on an object is proportional to its [acceleration](#). Thus, in this model, we will approximate all accelerations as constant.

#### The object's size and shape are unimportant

In physics, the word particle refers to a hypothetical entity having no size and therefore no shape. The statement that objects will be thought of as particles simply means that we will ignore any effects their actual size and shape could have on the scenarios we investigate.

#### The object is classical

For this entire course, we will assume the objects under investigation are much larger than an atom and move much slower than the speed of light.

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## 03. Kinematics

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- 01. Position, Velocity, Acceleration in 1D
- 02. Drawing Motion Diagrams in 1D
- 03. Drawing Motion Graphs in 1D
- 04. Tabulating Motion Information - Motion Graphs
- 05. Analyzing a More Complex Motion
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## 01. Concepts and Principles

### 1. Position, Velocity & Acceleration in 1D

1. Position ( $r$ ) Units: meters ( $m$ )
2. Velocity ( $v$ ) Units: meters per second ( $m/s$ )
3. Acceleration ( $a$ ) Units: meters per second per second ( $m/s^2$ )

### Position, Velocity & Acceleration in 1D

Kinematics is the formal language physicists use to describe motion. The need for a formal language is evidenced by a simple experiment: drop an object from about shoulder height and ask two people to independently describe the motion of the object. Chances are that the descriptions will not be in perfect agreement, even though both observers described the same motion. Obviously, a more formal way of describing motion is necessary to eliminate this type of descriptive ambiguity. Kinematics is the formal method of describing motion.

Three parameters are carefully defined and used by physicists to describe motion. Specifying these three parameters at all times forms a complete description of the motion of an object.

#### Position ( $r$ ) Units: meters ( $m$ )

The position of an object is its location relative to a well-defined coordinate system at a particular instant of time. Without a specified coordinate system, position is a meaningless concept. A coordinate system is comprised of *azero*, a specified *positive direction*, and a *scale*.

For example, in the hypothetical experiment in which the object was dropped from shoulder height, a coordinate system could have been defined in which the zero position was at ground level, the positive direction was up, and the scale used was meters. Using this coordinate system, the position of the object could have been specified at any particular instant of time. Of course, choosing the zero at the location at which the object was dropped, the positive direction as down, and the scale in feet is also perfectly acceptable. It doesn't matter what you choose as a coordinate system, only that you explicitly choose one. Depending on the coordinate system chosen, the position of an object can be positive, negative, or zero. We will use the symbol  $r$  to designate position.

#### Velocity ( $v$ ) Units: meters per second ( $m/s$ )

Although the word velocity is often used loosely in everyday conversation, its meaning in physics is specific and well-defined. To physicists, the velocity is the rate at which the position is changing. The velocity can be specified at any particular instant of time.

For example, if the position is changing quickly the velocity is large, and if the position is not changing the velocity is zero. A mathematical way to represent this definition is

$$v = \frac{\Delta r}{\Delta t} = \frac{r_{\text{final}} - r_{\text{initial}}}{t_{\text{final}} - t_{\text{initial}}} \quad (1)$$

where  $\Delta r$  represents the change in position and  $\Delta t$  the corresponding change in time. (Actually, this is the average velocity of the object over the time interval  $\Delta t$ , but as the time interval becomes smaller and smaller, the value of this expression becomes closer and closer to the actual rate at which the position is changing at one particular instant of time.)

Since the final position of the object ( $r_{\text{final}}$ ) may be either positive, negative, or zero, and either larger, smaller, or the same as the initial position ( $r_{\text{initial}}$ ), the velocity may be positive, negative, or zero. The sign of the velocity depends on the coordinate system chosen to define the position. A positive velocity simply means that the object is moving in the positive direction, as defined by the coordinate system, while a negative velocity means the object is traveling in the other direction.

#### Acceleration ( $a$ ) Units: meters per second per second ( $m/s^2$ )

Again, although the word acceleration is often used loosely in everyday conversation, its meaning in physics is specific and well-defined. To physicists, the acceleration is the rate at which the velocity is changing. Again, the acceleration can be specified at any particular instant of time.

For example, if the velocity is changing quickly the acceleration is large, and if the velocity is not changing the acceleration is zero. If an object has non-zero acceleration, it does *not* mean that the object is speeding up. It simply means that the velocity is changing. Moreover, even if an object has a *positive* acceleration, it does not mean that the object is speeding up. A positive acceleration means that the change in the velocity points in the positive direction. (I can almost guarantee you will experience confusion about this. Take some time to think about the preceding statement right now.)

A mathematical way to represent acceleration is

$$a = \frac{\Delta v}{\Delta t} = \frac{v_{\text{final}} - v_{\text{initial}}}{t_{\text{final}} - t_{\text{initial}}} \quad (2)$$

Again, this is actually the average acceleration of the object over the time interval  $\Delta t$ , but as the time interval becomes smaller and smaller, the value of this expression becomes closer and closer to the actual rate at which the velocity is changing at one particular instant of time.

Since  $v_{\text{final}}$  may be either positive, negative, or zero, and either larger, smaller, or the same as  $v_{\text{initial}}$ , the acceleration may be positive, negative, or zero. The algebraic sign of the acceleration depends on the coordinate system chosen to define the position. A negative acceleration means that the change in the velocity points in the negative direction. For example, the velocity could be in the positive direction and the object slowing down *or* the velocity could be in the negative direction and the object speeding up. Both of these scenarios would result in a negative acceleration. Conversely, a positive acceleration means that the change in the velocity points in the positive direction.

Kinematics is the correct use of the parameters position, velocity, and acceleration to describe motion. Learning to use these three terms correctly can be made much easier by learning a few tricks of the trade. These tricks, or analysis tools, are detailed in the following section.

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## 02. Drawing Motion Diagrams in 1D

### 1. Drawing Motion Diagrams (Qualitative)

1. Determining the position from a motion diagram
2. Determining the velocity from a motion diagram
3. Determining the acceleration from a motion diagram

### Drawing Motion Diagrams (Qualitative)

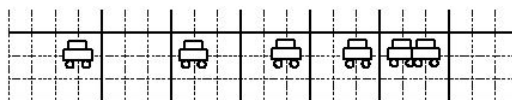
The words used by physicists to describe the motion of objects are defined above. However, learning to use these terms correctly is more difficult than simply memorizing definitions. An extremely useful tool for bridging the gap between a normal, conversational description of a situation and a physicists' description is the motion diagram. A motion diagram is the first step in translating a verbal description of a phenomenon into a physicists' description.

Start with the following verbal description of a physical situation:

*The driver of an automobile traveling at 15 m/s, noticing a red-light 30 m ahead, applies the brakes of her car until she stops just short of the intersection.*

### Determining the position from a motion diagram

A motion diagram can be thought of as a multiple-exposure photograph of the physical situation, with the image of the object exposed onto the film at equal time intervals. (You may have seen photographs of this type taken with a strobe-light.) For example, a multiple-exposure photograph of the situation described above would look something like this:



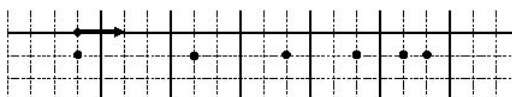
Note that the images of the automobile are getting closer together near the end of its motion because the car is traveling a smaller distance between the equally-timed exposures.

In general, in drawing motion diagrams it is better to represent the object as simply a dot, unless the actual shape of the object conveys some interesting information. Thus, a better motion diagram would be:



Since the purpose of the motion diagram is to help us describe the car's motion, a coordinate system is necessary. Remember, to define a coordinate system you must choose a zero, define a positive direction, and select a scale. We will always use meters as our position scale in this course, so you must only select a zero and a positive direction. Remember, there is no correct answer. Any coordinate system is as correct as any other.

The choice below indicates that the initial **position** of the car is the origin, and positions to the right of that are positive.



We can now describe the position of the car. The car starts at position zero and then has positive, increasing positions throughout the remainder of its motion.

### Determining the velocity from a motion diagram

Since velocity is the change in position of the car during a corresponding time interval, and we are free to select the time interval as the time interval between exposures on our multiple-exposure photograph, the velocity is simply the change in the position of the car "between dots." Thus, the arrows (vectors) on the motion diagram below represent the **velocity** of the car.



We can now describe the velocity of the car. Since the velocity vectors always point in the positive direction, the velocity is always positive. The car starts with a large, positive velocity which gradually declines until the velocity of the car is zero at the end of its motion.

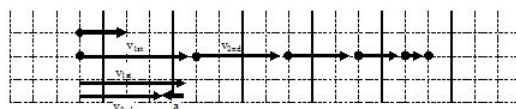
### Determining the acceleration from a motion diagram

Since acceleration is the change in velocity of the car during a corresponding time interval, and we are free to select the time interval as the time interval between exposures on our multiple-exposure photograph, we can determine the acceleration by comparing two successive velocities. The change in these velocity vectors will represent the acceleration.

To determine the acceleration,

- select two successive velocity vectors,
- draw them starting from the same point,
- construct the vector (arrow) that connects the tip of the *first* velocity vector to the tip of the *second* velocity vector.
- The vector you have constructed represents the acceleration.

Comparing the first and second velocity vectors leads to the **acceleration** vector shown below:



Thus, the acceleration points to the left and is therefore negative. You could construct the acceleration vector at every point in time, but hopefully you can see that as long as the velocity vectors continue to point toward the right and decrease in magnitude, the acceleration will remain negative.

Thus, with the help of a motion diagram, you can extract lots of information about the position, velocity, and acceleration of an object. You are well on your way to a complete kinematic description.

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## 03. Drawing Motion Graphs in 1D

1. Drawing Motion Graphs
2. Drawing the position vs. time graph
3. Drawing the velocity vs. time graph
4. Drawing the acceleration vs. time graph

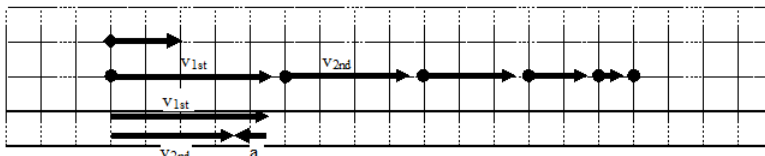
### Drawing Motion Graphs

Another useful way to describe the motion of an object is by constructing graphs of the object's position, velocity, and acceleration vs. time. A graphical representation is a very effective means of presenting information concerning an object's motion, and, moreover, it is relatively easy to construct motion graphs if you have a correct motion diagram.

Examine the same situation as before:

Template:ExampleStart

The driver of an automobile traveling at 15 m/s, noticing a red-light 30 m ahead, applies the brakes of her car until she stops just short of the intersection.

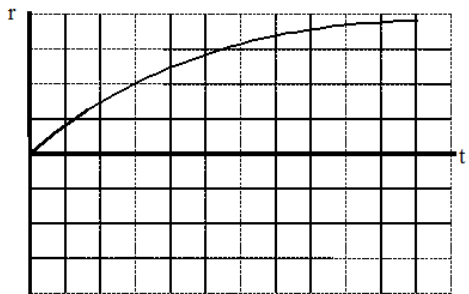


The verbal representation of the situation has already been translated into a motion diagram. A careful reading of the motion diagram allows the construction of the motion graphs.

Template:ExampleEnd

### Drawing the position vs. time graph

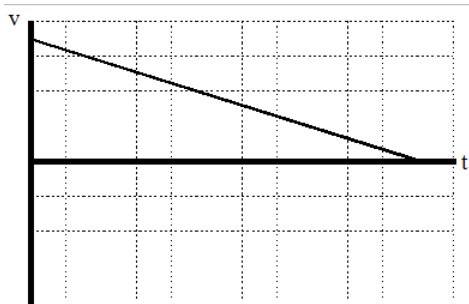
We already know, from the motion diagram, that the car starts at position zero, then has positive, increasing positions throughout the remainder of its motion. This information can be transferred onto a position vs. time graph.



Notice that the position is zero when the time is equal to zero, the position is always positive, and the position increases as time increases. Also note that in each subsequent second, the car changes its position by a smaller amount. This leads to the graph of position vs. time gradually decreasing in slope until it achieves a slope of zero. Once the car stops, the position of the car should not change.

### Drawing the velocity vs. time graph

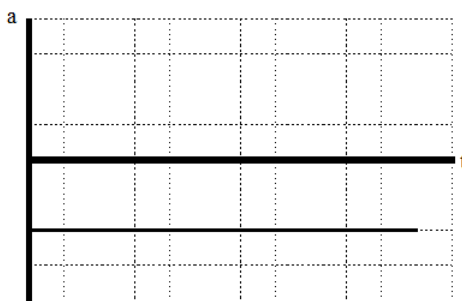
From the motion diagram, we know that the velocity of the car is always positive, starts large in magnitude, and decreases until it is zero. This information can be transferred onto a velocity vs. time graph.



How do we know that the slope of the line is constant? The slope of the line represents the rate at which the velocity is changing, and the rate at which the velocity is changing is termed the acceleration. Since in this model of mechanics we will only consider particles undergoing constant acceleration, the slope of a line on a velocity vs. time graph must be constant.

## Drawing the acceleration vs. time graph

From the motion diagram, the acceleration of the car can be determined to be negative at every point. Again, in this pass through mechanics, we will only be investigating scenarios in which the acceleration is constant. Thus, a correct acceleration vs. time graph is shown below.



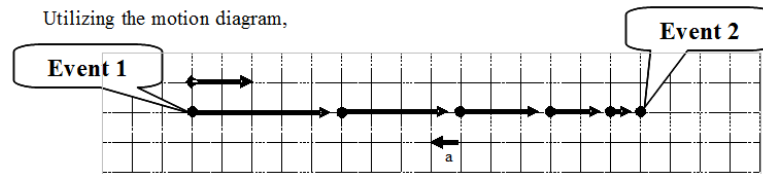
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## 04. Tabulating Motion Information - Motion Graphs

1. Tabulating Motion Information - Motion Graphs
2. Doing the Math (Quantitative)

### Tabulating Motion Information - Motion Graphs

After constructing the two qualitative representations of the motion (the motion diagram and the motion graphs), we are ready to tackle the quantitative aspects of the motion.



you can now assign numerical values to several of the kinematic variables. A glance at the situation description should indicate that information is presented about the car at two distinct events. Information is available about the car at the instant the driver applies the brakes (the velocity is given), and the instant the driver stops (the position is given). Other information can also be determined by referencing the motion diagram. To tabulate this information, you should construct a *motion table*.

Event 1: The instant the driver first applies the brakes.	Event 2: The instant the car finally stops.
$t_1 = 0 \text{ s}$	$t_2 = ?$
$r_1 = 0 \text{ m}$	$r_2 = +30 \text{ m}$
$v_1 = +15 \text{ m/s}$	$v_2 = 0 \text{ m/s}$
$a_{12} = ?$	

In addition to the information explicitly given, the velocity at the first event and the position at the second event, other information can be extracted from the problem statement and the motion diagram.

#### Example

- the position of the car at the first event is zero because the origin of the coordinate system is at this point,
- the time at the first instant can be "set to zero" by imagining a hypothetical stopwatch that is clicked on as the car begins to brake,
- and the velocity at the second event is zero because the car is stopped.

Since you are working under the assumption in this model that the acceleration is constant, the acceleration *between* the two instants in time is some unknown, constant value. To remind you that this assumption is in place, the acceleration is not labeled at the first instant,  $a_1$ , or the second instant,  $a_2$ , but rather as the acceleration *between* the two instants in time,  $a_{12}$ .

You now have a complete tabulation of all the information presented, both explicitly and implicitly, in the situation description. Moreover, you now can easily see that the only kinematic information not known about the situation is the assumed constant acceleration of the auto and the time at which it finally stops. Thus, to complete a kinematic description of the situation, these two quantities must be determined. What you may not know is that you have already been presented with the information needed to determine these two unknowns.

### Doing the Math (Quantitative)

In the concepts and principles portion of this unit, you were presented with two formal, mathematical relationships, the definitions of velocity and acceleration. In the example that you are working on, there are two unknown kinematic quantities. You should

remember from algebra that two equations are sufficient to calculate two unknowns. Thus, by applying the two definitions, you should be able to determine the acceleration of the car and the time at which it comes to rest.

Although you can simply apply the two definitions directly, normally the two definitions are rewritten, after some algebraic re-arranging, into two different relationships. This rearrangement is simply to make the algebra involved in solving for the unknowns easier. It is by no means necessary to solve the problem. In fact, the two definitions can be written in a large number of different ways, although this does not mean that there are a large number of different formulas you must memorize in order to analyze kinematic situations. There are only two independent kinematic relationships. The two kinematic relationships we will use when the acceleration is constant are:

$$v_2 = v_1 + a_{12}(t_2 - t_1)$$

and

$$r_2 = r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2$$

To finish the analysis,

<p>Event 1: The instant the driver first applies the brakes.</p> <p><math>t_1 = 0 \text{ s}</math></p> <p><math>r_1 = 0 \text{ m}</math></p> <p><math>v_1 = +15 \text{ m/s}</math></p> <p><math>a_{12} =</math></p>	<p>Event 2: The instant the car finally stops.</p> <p><math>t_2 =</math></p> <p><math>r_2 = +30 \text{ m}</math></p> <p><math>v_2 = 0 \text{ m/s}</math></p>
---	--

simply write the two kinematic relationships, input the known kinematic variables from the motion table, and solve the two relations for the two unknowns. (This process is not physics, it's algebra.)

$$\begin{aligned} v_2 &= v_1 + a_{12}(t_2 - t_1) \\ 0 &= 15 + a_{12}(t_2 - 0) \\ a_{12} &= \frac{-15}{t_2} \end{aligned} \quad \begin{aligned} r_2 &= r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2 \\ 30 &= 0 + 15(t_2 - 0) + \frac{1}{2}a_{12}(t_2 - 0)^2 \end{aligned}$$

Now substitute this expression into the other equation:

$$\begin{aligned} 30 &= 15t_2 + \frac{1}{2}a_{12}t_2^2 \\ 30 &= 15t_2 + \frac{1}{2}\left(\frac{-15}{t_2}\right)t_2^2 \\ 30 &= 15t_2 - 7.5t_2 \\ 30 &= 7.5t_2 \\ t_2 &= 4.0 \text{ s} \end{aligned}$$

Substitute this result back into the original equation:

$$\begin{aligned} a_{12} &= \frac{-15}{4} \\ a_{12} &= -3.8 \text{ m/s}^2 \end{aligned}$$

Thus, the car must have accelerated at  $3.8 \text{ m/s}^2$  in the negative direction, and stopped after 4.0 seconds. The kinematic description of the situation is complete.

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## 05. Analyzing a More Complex Motion

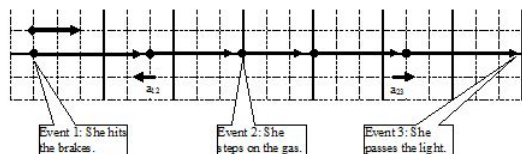
### 1. Analyzing a More Complex Motion

### Analyzing a More Complex Motion

Let's re-visit our scenario, although this time the light turns green while the car is slowing down:

The driver of an automobile traveling at 15 m/s, noticing a red-light 30 m ahead, applies the brakes of her car. When she is 10 m from the light, and traveling at 8.0 m/s, the light turns green. She instantly steps on the gas and is back at her original speed as she passes under the light.

Our first step in analyzing this motion should be to draw a motion diagram.



I've noted on the motion diagram the important events that take place during the motion. Notice that between the instant she hits the brakes and the instant she steps on the gas the acceleration is negative, while between the instant she steps on the gas and the instant she passes the light the acceleration is positive. Thus, in tabulating the motion information and applying the kinematic relations we will have to be careful not to confuse kinematic variables between these two intervals. Below is a tabulation of motion information using the coordinate system established in the motion diagram.

<p>Event 1: She hits the brakes.</p> <p><math>t_1 = 0</math> s</p> <p><math>r_1 = 0</math> m</p> <p><math>v_1 = +15</math> m/s</p> <p><math>a_{12} =</math></p>	<p>Event 2: She steps on the gas.</p> <p><math>t_2 =</math></p> <p><math>r_2 = +20</math> m</p> <p><math>v_2 = +8.0</math> m/s</p> <p><math>a_{23} =</math></p>	<p>Event 3: She passes the light</p> <p><math>t_3 =</math></p> <p><math>r_3 = +30</math> m</p> <p><math>v_3 = +15</math> m/s</p>
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First, notice that during the time interval between “hitting the brakes” and “stepping on the gas” there are two kinematic variables that are unknown. Recall that by using your two kinematic relations, you should be able to determine these values. Second, notice that during the second time interval again two variables are unknown. Once again, the two kinematic relations will allow you to determine these values. Thus, before I actually begin to *do* the algebra, I *know* the unknown variables can be determined!

First let's examine the motion between hitting the brakes and stepping on the gas:

$$\begin{aligned}
 v_2 &= v_1 + a_{12}(t_2 - t_1) \\
 8 &= 15 + a_{12}(t_2 - 0) \\
 a_{12} &= \frac{-7}{t_2}
 \end{aligned}$$

Now substitute this expression into the other equation:

$$\begin{aligned}
 r_2 &= r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2 \\
 20 &= 0 + 15(t_2 - 0) + \frac{1}{2}a_{12}(t_2 - 0)^2 \\
 20 &= 15t_2 + \frac{1}{2}a_{12}t_2^2 \\
 20 &= 15t_2 + \frac{1}{2}\left(\frac{-7}{t_2}\right)t_2^2 \\
 20 &= 15t_2 - 3.5t_2 \\
 20 &= 11.5t_2 \\
 t_2 &= 1.74\text{ s}
 \end{aligned}$$

Substitute this result back into the original equation:

$$\begin{aligned}
 a_{12} &= \frac{-7}{1.74} \\
 a_{12} &= -4.03\text{ m/s}^2
 \end{aligned}$$

Now, using these results, examine the kinematics between stepping on the gas and passing the light. Note that the initial values of the kinematic variables are denoted by '2' and the final values by '3', since we are examining the interval between event 2 and event 3.

$$v_3 = v_2 + a_{23}(t_3 - t_2)$$

$$15 = 8 + a_{23}(t_3 - 1.74)$$

$$a_{23} = \frac{7}{(t_3 - 1.74)}$$

Now substitute this expression into the other equation:

$$a_{23} = \frac{7}{(2.61 - 1.74)}$$

$$a_{12} = 8.0 \text{ m/s}^2$$

$$r_3 = r_2 + v_2(t_3 - t_2) + \frac{1}{2}a_{23}(t_3 - t_2)^2$$

$$30 = 20 + 8(t_3 - 1.74) + \frac{1}{2}a_{23}(t_3 - 1.74)^2$$

$$10 = 8(t_3 - 1.74) + \frac{1}{2}a_{23}(t_3 - 1.74)^2$$

$$10 = 8(t_3 - 1.74) + \frac{1}{2}\left(\frac{7}{(t_3 - 1.74)}\right)(t_3 - 1.74)^2$$

$$10 = 8(t_3 - 1.74) + 3.5(t_3 - 1.74)$$

$$10 = 11.5(t_3 - 1.74)$$

$$0.87 = t_3 - 1.74$$

$$t_3 = 2.61 \text{ s}$$

Substitute this result back into the original equation:

We now have a complete kinematic description of the motion.

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## 06. Symbolic Analysis

### 1. No numbers – just variables

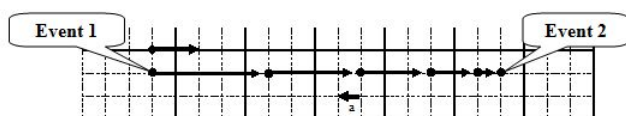
### No numbers – just variables

Working the problem out using only variables will save you time and a sore hand. The object here is to first recognize what you need to solve for. Then solve for the variable without numbers, so that you will have your equation in the form that best models what you need to use it for. Causing you to just plug in all of your numbers and pop out an answer, without having to move long decimals and complex fractions to each side of the equation. Therefore solving your problem with minimal error and using less time.

**Consider the following situation:**

*The driver of an automobile suddenly sees an obstacle blocking her lane. Determine the total distance the auto travels between seeing the obstacle and stopping ( $d$ ) as a function of the initial velocity of the car ( $v_i$ ) and the magnitude of its acceleration while stopping ( $a_s$ ).*

As always, the first step in analyzing motion is to draw a motion diagram.



Rather than calculate the stopping distance for particular values of initial velocity and acceleration, the goal of this activity is to determine, in general, how the stopping distance depends on these two parameters. If we can construct this function, we can then use the result to calculate the stopping distance for *any* car if we know its initial velocity and stopping acceleration.

Although this sounds like a different task from what we've done in the previous two examples, we will approach this task exactly the same way, by tabulating what we know about the situation,

Event 1: The instant the driver hits the brakes.

$$t_1 = 0 \text{ s}$$

$$r_1 = 0 \text{ m}$$

$$v_1 = v_i$$

$$a_{12} = -a_s$$

Event 2: The instant the car stops.

$$t_2 =$$

$$r_2 = d$$

$$v_2 = 0 \text{ m/s}$$

and then applying our two kinematic relationships:

$$v_2 = v_1 + a_{12}(t_2 - t_1) \quad r_2 = r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2$$

$$0 = v_i - a_s t_2 \quad d = 0 + v_i t_2 - \frac{1}{2}a_s t_2^2$$

Since our goal is to determine  $d$  as a function of  $v_i$  and  $a_s$ , we must eliminate  $t_2$ . To do this, solve for  $t_2$  in the left equation and substitute this expression into the right equation.

$$t_2 = \frac{v_i}{a_s} \quad d = v_i \left( \frac{v_i}{a_s} \right) - \frac{1}{2} a_s \left( \frac{v_i}{a_s} \right)^2$$

$$d = \frac{v_i^2}{a_s} - \frac{1}{2} \frac{v_i^2}{a_s}$$

$$d = \frac{v_i^2}{2a_s}$$

Thus, the stopping distance appears to be proportional to the *square* of the initial velocity and inversely proportional to the stopping acceleration. Does this make sense?

To determine if a symbolic expression is sensible, it is often useful to check *limiting cases*. A limiting case is when one of the variables in the expression takes on an extreme value, typically zero or infinity. For example, if the initial velocity of the car was zero the stopping distance would have to also be zero, since the car was never moving! Allowing  $v_i$  to equal zero in the above

expression results in a stopping distance of zero, so our expression “passes” this logical test. Another limiting case would be setting the acceleration of the car equal to zero. With no acceleration, the car should never stop. In our expression, setting the acceleration equal to zero results in an infinite stopping distance, which again agrees with commonsense. If our expression did not give the correct results in these limiting cases, we would know we made an error somewhere in the derivation (and, of course, we would then go back and find our mistake and fix it because we are good students ...).

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## 07. Addendum

1. [Algebraic Signs](#)
2. [Deriving the kinematic relationships](#)

### Algebraic Signs

Confusion about the meaning of algebraic signs is common among beginning physics students. The best way to clarify this confusion is to remember that algebraic signs are simply a mathematical way to describe direction. Instead of saying up and down, or east and west, physicists construct coordinate systems and translate the words east and west into the symbols '+' and '-', or even '-' and '+' if we choose a different coordinate system. The key to the translation is the coordinate system. A coordinate system is very similar to the English-French dictionary you might take with you on your first trip to France. When you see a '-', use your coordinate system to translate it into a verbal description of direction.

Do not fall into the common habit of translating a '-' into the word "decreasing". A negative acceleration, for example, does *NOT* imply that the object is slowing down. It implies an acceleration that points in the negative direction. It is *impossible* to determine whether an object is speeding up or slowing down by looking at the sign of the acceleration! Conversely, the word "deceleration", which does mean that an object is slowing down, does *not* give any information regarding the sign of the acceleration. I can decelerate in the positive direction as easily as I can decelerate in the negative direction.

A guide to determining the correct algebraic sign:

### Deriving the kinematic relationships

Let's construct the two independent kinematic relationships that you will use whenever the acceleration is constant. In a later chapter, we will return to the case in which the acceleration is not constant.

From the definition of acceleration:

The above relationship is **our first kinematic relationship**. The acceleration in this relationship is really the average acceleration. However, since the acceleration is constant, the average acceleration is the same as the acceleration at any instant between the initial and final state.

From the definition of velocity:

You must remember, however, that the velocity in this formula is really the average velocity of the object over the time interval selected. To keep you from having to remember this fact, we can rewrite the average velocity as the sum of the initial velocity and the final velocity divided by two:

Substituting in the first kinematic relationship for the final velocity yields:

The above relationship is **our second kinematic relationship**.

Although we could keep deriving new kinematic equations forever, it is *impossible* that any other derived equation could allow us to calculate some quantity that these equations do not allow us to calculate.

[Paul D'Alessandris](#) ([Monroe Community College](#))

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## 01. Algebraic Signs

Confusion about the meaning of algebraic signs is common among beginning physics students. The best way to clarify this confusion is to remember that algebraic signs are simply a mathematical way to describe direction. Instead of saying up and down, or east and west, physicists construct coordinate systems and translate the words east and west into the symbols '+' and '-', or even '-' and '+' if we choose a different coordinate system. The key to the translation is the coordinate system. A coordinate system is very similar to the English-French dictionary you might take with you on your first trip to France. When you see a '-', use your coordinate system to translate it into a verbal description of direction.

Do not fall into the common habit of translating a '-' into the word "decreasing". A negative acceleration, for example, does *NOT* imply that the object is slowing down. It implies an acceleration that points in the negative direction. It is *impossible* to determine whether an object is speeding up or slowing down by looking at the sign of the acceleration! Conversely, the word "deceleration", which does mean that an object is slowing down, does *not* give any information regarding the sign of the acceleration. I can decelerate in the positive direction as easily as I can decelerate in the negative direction.

A guide to determining the correct algebraic sign:

	Acceleration	Direction	Algebraic Sign
Speeding Up	+	+	+
	-	-	
Slowing Down	-	+	-
	+	-	

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## 02. Deriving the kinematic relationships

Let's construct the two independent kinematic relationships that you will use whenever the acceleration is constant. In a later chapter, we will return to the case in which the acceleration is not constant.

From the definition of acceleration:

$$a = \frac{v_f - v_i}{t_f - t_i}$$
$$v_f - v_i = a(t_f - t_i)$$
$$\boxed{v_f = v_i + a(t_f - t_i)}$$

The above relationship is **our first kinematic relationship**. The acceleration in this relationship is really the average acceleration. However, since the acceleration is constant, the average acceleration is the same as the acceleration at any instant between the initial and final state.

From the definition of velocity:

$$v = \frac{r_f - r_i}{t_f - t_i}$$
$$r_f - r_i = v(t_f - t_i)$$
$$r_f = r_i + v(t_f - t_i)$$

You must remember, however, that the velocity in this formula is really the average velocity of the object over the time interval selected. To keep you from having to remember this fact, we can rewrite the average velocity as the sum of the initial velocity and the final velocity divided by two:

$$r_f = r_i + \left(\frac{v_i + v_f}{2}\right)(t_f - t_i)$$
$$r_f = r_i + \frac{1}{2}(v_i + v_f)(t_f - t_i)$$

Substituting in the first kinematic relationship for the final velocity yields:

$$r_f = r_i + \frac{1}{2}(v_i + v_i + a(t_f - t_i))(t_f - t_i)$$
$$r_f = r_i + \frac{1}{2}(2v_i(t_f - t_i) + a(t_f - t_i)^2)$$
$$\boxed{r_f = r_i + v_i(t_f - t_i) + \frac{1}{2}a(t_f - t_i)^2}$$

The above relationship is **our second kinematic relationship**.

Although we could keep deriving new kinematic equations forever, it is *impossible* that any other derived equation could allow us to calculate some quantity that these equations do not allow us to calculate.

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## 08. Activities

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## 01. Selected Answers For Kinematics

### Selected Answers for ModelSpiral Physics: Mechanics

#### Kinematics

- 41.  $r_2 = 1.4 \text{ m}$
- 42.  $t_2 = 1.45 \text{ m}$
- 43.  $t_3 = 3.4 \text{ s}$
- 44.  $t_3 = 4.55 \text{ s}$
- 45.  $r_3 = 36.6 \text{ m}$
- 46.  $r_4 = 2.0 \text{ km}$
- 47.  $t_4 = 370 \text{ ks}$
- 49.  $t_2 = 14.9 \text{ s}$
- 50.  $t_2 = 7.8 \text{ s}$
- 51.  $t_3 = 2.87 \text{ s}$

#### Dynamics

- 80.  $F_{\text{rope}} = 420 \text{ N}$
- 81.  $F_{\text{cushion}} = 1.91 \text{ kN}$
- 83.  $F_{\text{bottom cable}} = 11.4 \text{ N}$
- 84.  $F_{\text{cushion}} = 2.83 \text{ kN}$
- 85.  $F_{\text{ground}} = 43.7 \text{ kN}$
- 86.  $r_3 = 63.5 \text{ m}$
- 87.  $r_3 = 1.55 \text{ km}$
- 89.  $m_{\text{block}} = 240 \text{ kg}$
- 90.  $m_{\text{block}} = 26 \text{ kg}$
- 91.  $F_{\text{rope}} = 500 \text{ N}$
- 92.  $F_{\text{rope}} = 490 \text{ N}$

#### Conservation Laws

- 122. a.  $F_{\text{scale}} = 755 \text{ N}$       b.  $F_{\text{scale}} = 780 \text{ N}$
- 123. a.  $v = 116 \text{ m/s}$       b.  $t = 26.8 \text{ s}$
- 124. a.  $v = 12.5 \text{ m/s}$       b.  $F_{\text{ground}} = 43.7 \text{ kN}$
- 125.  $t_2 = 0.89 \text{ s}$        $r_2 = 1.78 \text{ m}$
- 127.  $m_{\text{student}} = 94 \text{ kg}$
- 128.  $v_2 = 6.0 \text{ m/s}$
- 129.  $v_2 = 1.5 \text{ m/s}$
- 130.  $v_{2 \text{ ship}} = 24.3 \text{ m/s}$
- 131.  $v_{\text{probe}} = 667 \text{ m/s}$
- 132.  $v_{2 \text{ platform}} = 0.28 \text{ m/s}$
- 133.  $v_{2 \text{ platform}} = 0.06 \text{ m/s}$
- 134.  $v_{2 \text{ balloon}} = 6.3 \text{ m/s}$

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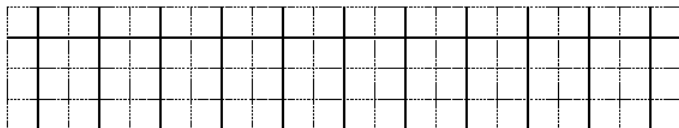
## I-24

Construct motion diagrams for the motions described below.

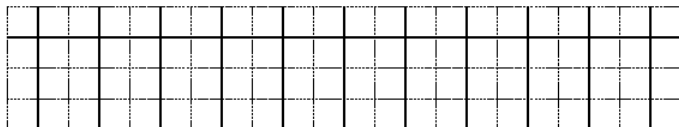
a. A subway train in Washington, D.C., starts from rest and accelerates at  $2.0 \text{ m/s}^2$  for 12 seconds.



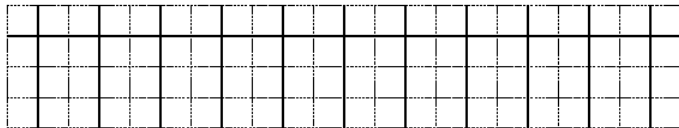
b. A pole-vaulter, just before touching the cushion on which she lands after a jump, is falling downward at a speed of  $10 \text{ m/s}$ . The pole-vaulter sinks about  $2.0 \text{ m}$  into the cushion before stopping.



c. The driver of a car traveling at  $35 \text{ m/s}$  suddenly sees a police car. The driver attempts to reach the speed limit of  $25 \text{ m/s}$  by accelerating at  $2.5 \text{ m/s}^2$ . The driver has a reaction time of  $0.55 \text{ s}$ . (The reaction time is the time between first seeing the police car and pressing the brake.)



d. A car, initially at rest, accelerates toward the west at  $2.0 \text{ m/s}^2$ . At the same time that the car starts moving, a truck,  $350 \text{ m}$  west of the car and moving at  $16 \text{ m/s}$  toward the east, starts to move slower, accelerating at  $1.0 \text{ m/s}^2$ . The car and truck pass safely.



## I-25

Construct motion diagrams for the motions described below.

a. A child is hanging from a rope by her hands. She exerts a burst of strength and 2.0 s later is traveling at 1.4 m/s up the rope.



b. An automobile comes to rest after skidding 35 m. The car's acceleration while skidding is known to be  $6.0 \text{ m/s}^2$ .



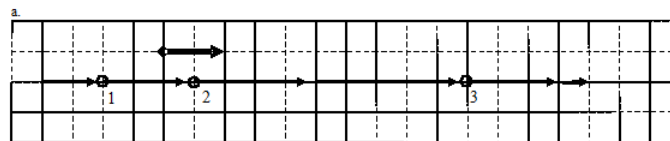
c. Tired of walking up the stairs, an engineering student designs an ingenious device for reaching her third floor dorm room. A block is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the block is released, the student is pulled up to her dorm room, 8.0 m off the ground, in a time of 1.8 s.



d. A car, initially traveling at 20 m/s to the east, accelerates toward the west at  $2.0 \text{ m/s}^2$ . At the same time that the car starts moving, a truck, 60 m west of the car and moving at 16 m/s toward the east, starts to move faster, accelerating at  $1.0 \text{ m/s}^2$ . It's a one-lane road and both drivers are too busy fiddling with their CD players to notice each other.

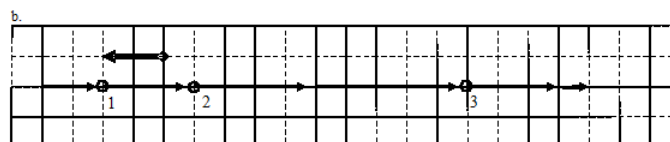


For each of the motion diagrams below, determine the algebraic sign (+, - or zero) of the position, velocity, and acceleration of the object at the location of the three open circles. Describe an actual motion that could be represented by each motion diagram.



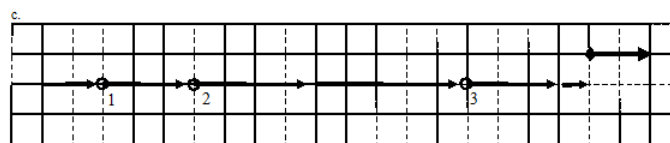
Description:

	1	2	3
$r$			
$v$			
$a$			



Description:

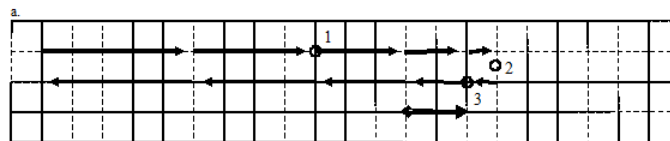
	1	2	3
$r$			
$v$			
$a$			



Description:

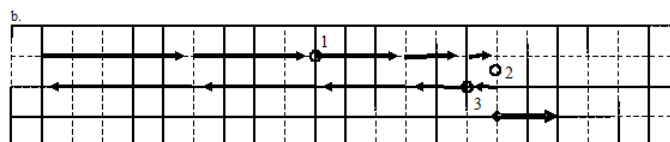
	1	2	3
$r$			
$v$			
$a$			

For each of the motion diagrams below, determine the algebraic sign (+, - or zero) of the position, velocity, and acceleration of the object at the location of the three open circles. Describe an actual motion that could be represented by each motion diagram.



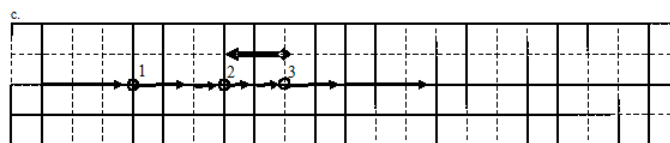
Description:

	1	2	3
r			
v			
a			



Description:

	1	2	3
r			
v			
a			

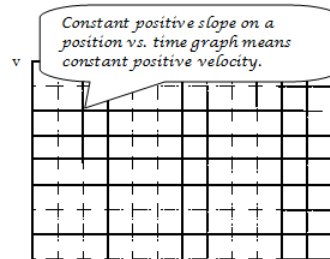
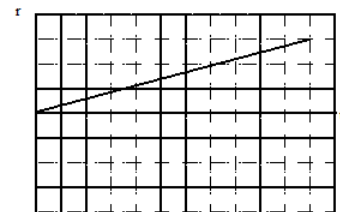


Description:

	1	2	3
r			
v			
a			

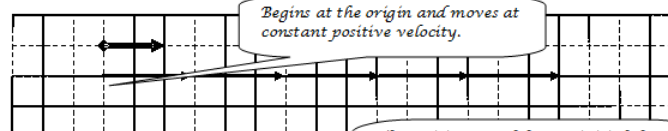
For each of the position vs. time graphs below, construct a corresponding motion diagram and velocity vs. time graph.

a. Motion Graphs



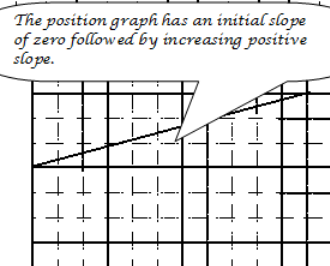
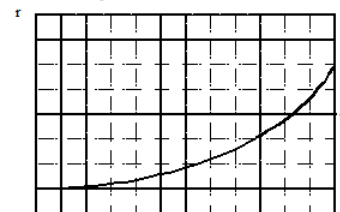
Constant positive slope on a position vs. time graph means constant positive velocity.

Motion Diagram



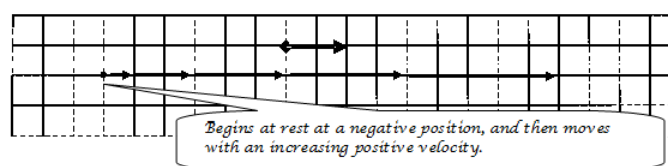
Begins at the origin and moves at constant positive velocity.

b. Motion Graphs



The position graph has an initial slope of zero followed by increasing positive slope.

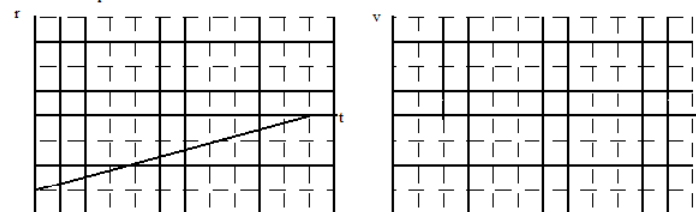
Motion Diagram



Begins at rest at a negative position, and then moves with an increasing positive velocity.

For each of the position vs. time graphs below, construct a corresponding motion diagram and velocity vs. time graph.

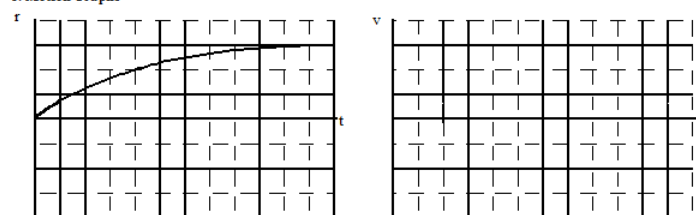
a. Motion Graphs



Motion Diagram



b. Motion Graphs

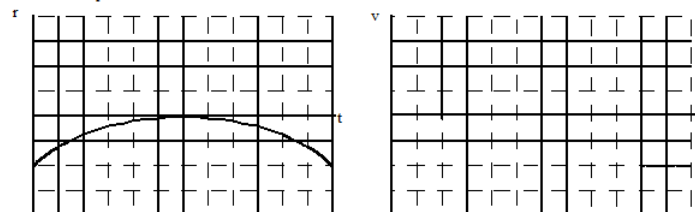


Motion Diagram

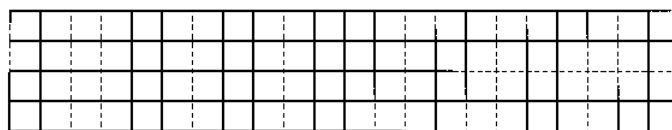


For each of the graphs below, construct a corresponding graph and motion diagram.

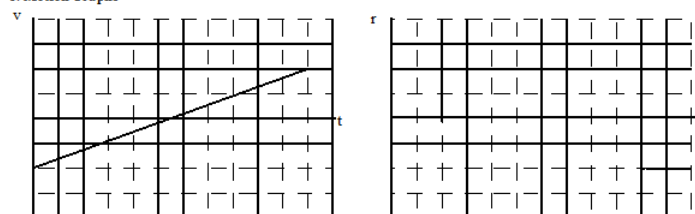
a. Motion Graphs



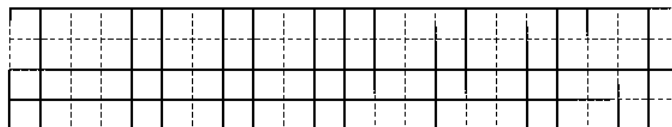
Motion Diagram



b. Motion Graphs



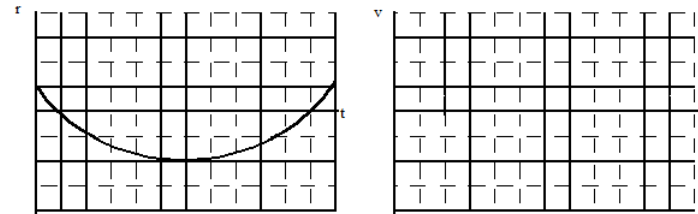
Motion Diagram



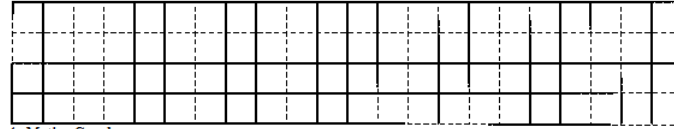


For each of the graphs below, construct a corresponding graph and motion diagram.

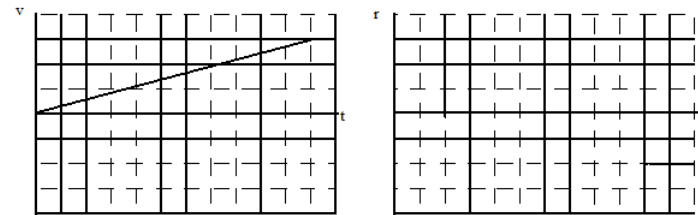
a. Motion Graphs



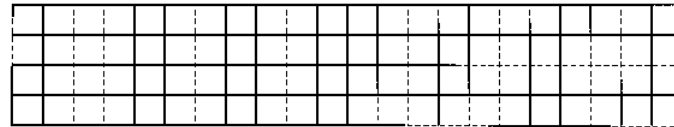
Motion Diagram



b. Motion Graphs

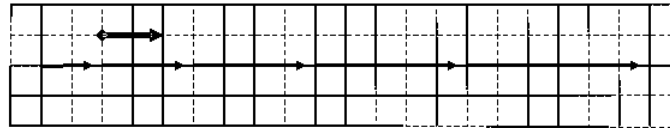


Motion Diagram

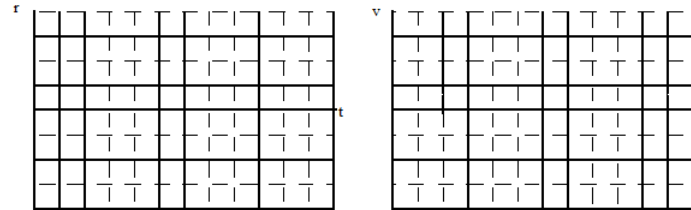


For each of the motion diagrams below, construct the corresponding motion graphs.

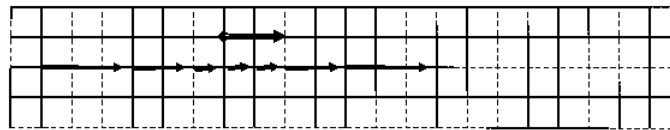
a. Motion Diagram



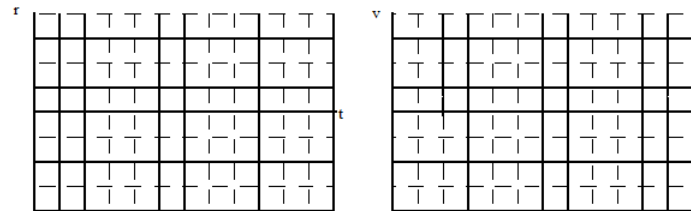
Motion Graphs



b. Motion Diagram

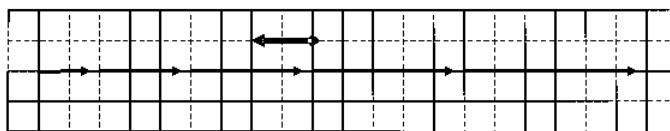


Motion Graphs

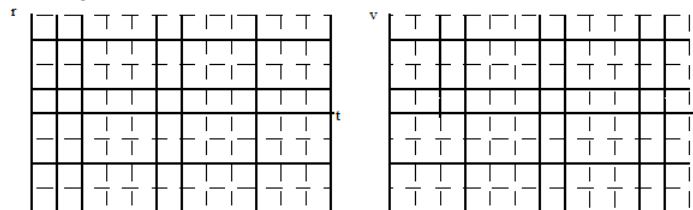


For each of the motion diagrams below, construct the corresponding motion graphs.

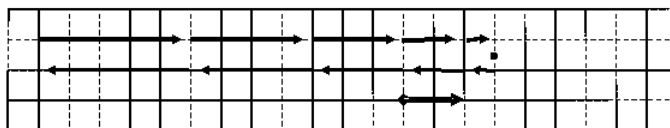
a. Motion Diagram



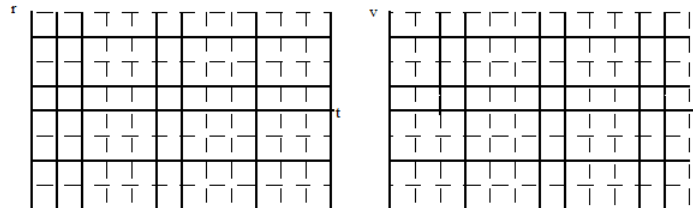
Motion Graphs

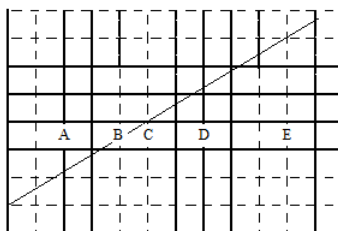


b. Motion Diagram



Motion Graphs





An object's motion is represented by the position vs. time graph at the top of the page.

a. Rank the object's position at the lettered times.

Largest 1. E 2. D 3. C 4. B 5. A Smallest

b. Rank the object's velocity at the lettered times.

Largest 1. A B C D E 2. \_\_\_\_\_ 3. \_\_\_\_\_ 4. \_\_\_\_\_ 5. \_\_\_\_\_ Smallest

c. Rank the object's acceleration at the lettered times.

Largest 1. A B C D E 2. \_\_\_\_\_ 3. \_\_\_\_\_ 4. \_\_\_\_\_ 5. \_\_\_\_\_ Smallest

Constant slope means constant velocity (and zero acceleration).

An object's motion is represented by the velocity vs. time graph at the top of the page.

d. Rank the object's position at the lettered times.

Largest 1. E 2. A 3. D 4. B 5. C Smallest

e. Rank the object's velocity at the lettered times.

Largest 1. E 2. D 3. C 4. B 5. A Smallest

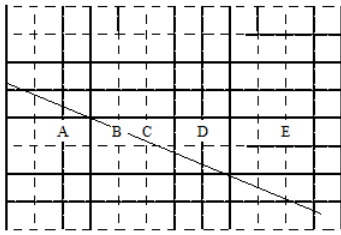
f. Rank the object's acceleration at the lettered times.

Largest 1. A B C D E 2. \_\_\_\_\_ 3. \_\_\_\_\_ 4. \_\_\_\_\_ 5. \_\_\_\_\_ Smallest

Constant slope means constant acceleration.

The motion diagram for the object is sketched above. Notice that regardless of where the origin is located, the turn-around point (C) is the smallest position.

## I-35



An object's motion is represented by the position vs. time graph at the top of the page.

a. Rank the object's position at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

b. Rank the object's velocity at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

c. Rank the object's acceleration at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

An object's motion is represented by the velocity vs. time graph at the top of the page.

d. Rank the object's position at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

e. Rank the object's velocity at the lettered times.

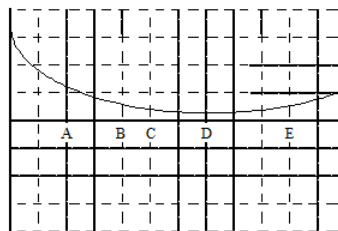
Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

f. Rank the object's acceleration at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

I-35

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An object's motion is represented by the position vs. time graph at the top of the page.

a. Rank the object's position at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

b. Rank the object's velocity at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

c. Rank the object's acceleration at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

An object's motion is represented by the velocity vs. time graph at the top of the page.

d. Rank the object's position at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

e. Rank the object's velocity at the lettered times.

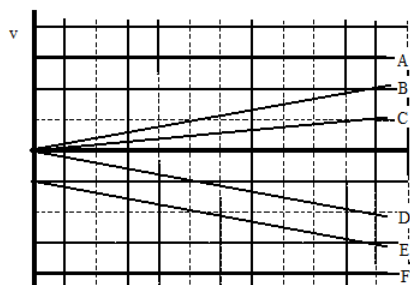
Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

f. Rank the object's acceleration at the lettered times.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest

# I-37

Below are velocity vs. time graphs for six different objects.



Rank these graphs on the basis of the distance traveled by each object.

Largest 1. F 2. A 3. E 4. B D 5. C 6. \_\_\_\_\_ Smallest

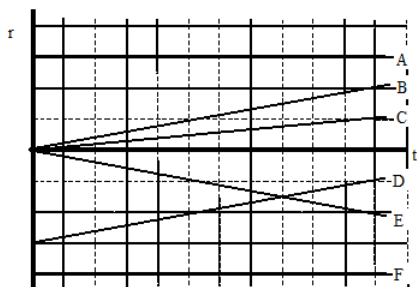
\_\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

*The faster you travel, the more distance you travel. The direction you are headed is not important. Therefore, F travels the largest distance because it is always moving the fastest, followed by A. E is consistently traveling faster than D, so it covers a larger distance than D. B and D always travel at the same speed (although in opposite directions) so they cover the same distance, therefore they are ranked as equal. C travels the slowest so it covers the least distance.*

## I-38

Below are position vs. time graphs for six different objects.



a. Rank these graphs on the basis of the velocity of the object. Rank positive velocities as larger than negative velocities.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. Rank these graphs on the basis of the acceleration of the object. Rank positive accelerations as larger than negative accelerations.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

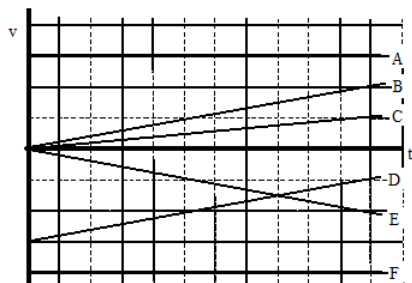
I-38

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## I-39

Below are velocity vs. time graphs for six different objects.



a. Rank these graphs on the basis of the final position of each object. Rank positive positions as larger than negative positions.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. Rank these graphs on the basis of the acceleration of the object. Rank positive accelerations as larger than negative accelerations.

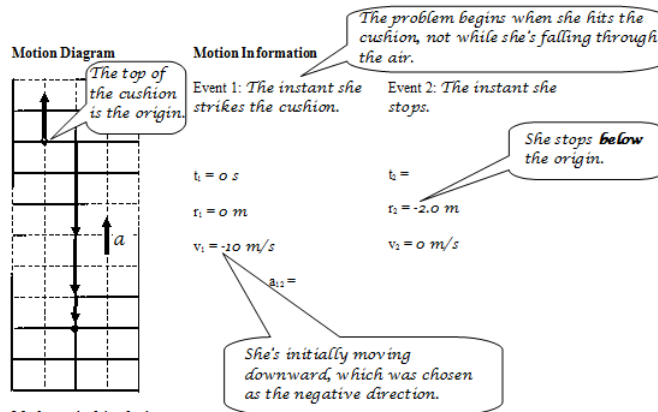
Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

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A pole-vaulter, just before touching the cushion on which she lands after a jump, is falling downward at a speed of 10 m/s. The pole-vaulter sinks about 2.0 m into the cushion before stopping.



#### Mathematical Analysis

$$v_2 = v_1 + a_{12}(t_2 - t_1)$$

$$0 = -10 + a_{12}(t_2 - 0)$$

$$a_{12} = \frac{10}{t_2}$$

Now substitute this expression into the other equation:

$$a_{12} = \frac{10}{0.4}$$

$$a_{12} = 25 \text{ m/s}^2$$

The acceleration is positive, as it should be since the jumper is moving downward and slowing down.

$$r_2 = r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2$$

$$-2 = 0 - 10(t_2 - 0) + \frac{1}{2}a_{12}(t_2 - 0)^2$$

$$-2 = -10t_2 + \frac{1}{2}a_{12}t_2^2$$

$$-2 = -10t_2 + \frac{1}{2}\left(\frac{10}{t_2}\right)t_2^2$$

$$-2 = -10t_2 + 5t_2$$

$$-2 = -5t_2$$

$$t_2 = 0.4 \text{ s}$$

Substitute this result back into the original equation:

## I-41

A child is hanging from a rope by her hands. She exerts a burst of strength and 2.0 s later is traveling at 1.4 m/s up the rope.

### Motion Diagram



### Motion Information

Event 1:

Event 2:

$t_1 =$

$t_2 =$

$r_1 =$

$r_2 =$

$v_1 =$

$v_2 =$

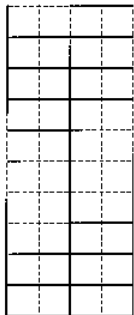
$a_{12} =$

### Mathematical Analysis

## I-42

An elevator is moving downward at 4.0 m/s when someone presses the emergency stop button. The elevator comes to rest after traveling 2.9 m.

### Motion Diagram



### Motion Information

Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

$a_{12} =$

Event 2:

$t_2 =$

$r_2 =$

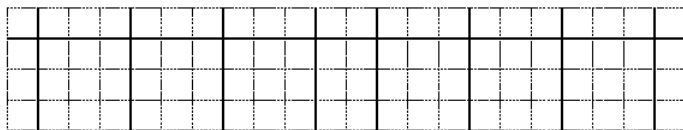
$v_2 =$

### Mathematical Analysis

## I-43

The driver of a car traveling at 16 m/s suddenly sees a truck that has entered from a side street and blocks the car's path. The car's maximum magnitude acceleration is  $6.0 \text{ m/s}^2$  and the driver has a reaction time of 0.75 s. (The reaction time is the time between first seeing the truck and pressing the brake.) The driver stops just in time to avoid an accident.

### Motion Diagram



### Motion Information

Event 1:

Event 2:

Event 3:

$t_1 =$

$t_2 =$

$t_3 =$

$r_1 =$

$r_2 =$

$r_3 =$

$v_1 =$

$v_2 =$

$v_3 =$

$a_{12} =$

$a_{23} =$

### Mathematical Analysis

I-43

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## I-44

The driver of a car traveling at 35 m/s suddenly sees a police car. The driver attempts to reach the speed limit of 25 m/s by accelerating at  $2.5 \text{ m/s}^2$ . The driver has a reaction time of 0.55 s. (The reaction time is the time between first seeing the police car and pressing the brake.)

### Motion Diagram



### Motion Information

Event 1:

Event 2:

Event 3:

$t_1 =$

$t_2 =$

$t_3 =$

$r_1 =$

$r_2 =$

$r_3 =$

$v_1 =$

$v_2 =$

$v_3 =$

$a_{12} =$

$a_{23} =$

### Mathematical Analysis

## I-45

An automobile, initially traveling at 15 m/s, begins to slow down as it approaches a red light. After traveling 15 m, and slowing to 3.0 m/s, the light turns green and the driver steps on the gas and accelerates for 2.4 seconds until she reaches her original speed.

### Motion Diagram



### Motion Information

Event 1:

Event 2:

Event 3:

$t_1 =$

$t_2 =$

$t_3 =$

$r_1 =$

$r_2 =$

$r_3 =$

$v_1 =$

$v_2 =$

$v_3 =$

$a_{12} =$

$a_{23} =$

### Mathematical Analysis

## I-46

A subway train in Washington, D.C., starts from rest and accelerates at  $2.0 \text{ m/s}^2$  for  $12 \text{ s}$ . The train travels at a constant speed for  $65 \text{ s}$ . The speed of the train then decreases for  $25 \text{ s}$  until it reaches the next station.

### Motion Diagram



### Motion Information

Event 1:

Event 2:

Event 3:

Event 4:

$t_1 =$

$t_2 =$

$t_3 =$

$t_4 =$

$r_1 =$

$r_2 =$

$r_3 =$

$r_4 =$

$v_1 =$

$v_2 =$

$v_3 =$

$v_4 =$

$a_{12} =$

$a_{23} =$

$a_{34} =$

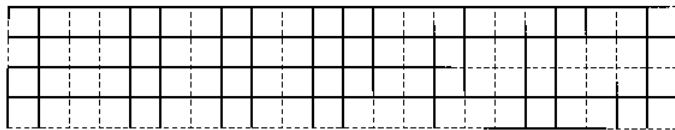
### Mathematical Analysis



## I-47

A rocket ship is launched from rest from a space station. Its destination is  $1.0 \times 10^{11}$  m away. The ship is programmed to accelerate at  $7.4 \text{ m/s}^2$  for 12 hours. After 12 hours, the ship will travel at constant velocity until it comes within  $1.0 \times 10^6$  m of its destination. Then, it will fire its retrorockets to land safely.

### Motion Diagram



### Motion Information

Event 1:                      Event 2:                      Event 3:                      Event 4:

$t_1 =$                        $t_2 =$                        $t_3 =$                        $t_4 =$

$r_1 =$                        $r_2 =$                        $r_3 =$                        $r_4 =$

$v_1 =$                        $v_2 =$                        $v_3 =$                        $v_4 =$

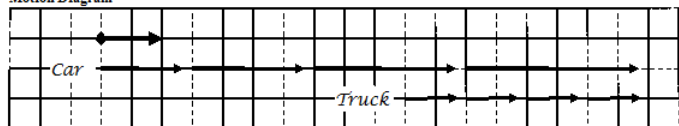
$a_{12} =$                        $a_{23} =$                        $a_{34} =$

### Mathematical Analysis

# I-48

The driver of a car traveling at 16 m/s sees a truck 20 m ahead traveling at a constant speed of 12 m/s. The car starts without delay to accelerate at 4.0 m/s<sup>2</sup> in an attempt to rear-end the truck. The truck driver is too busy talking on his cellular phone to notice the car.

## Motion Diagram



## Motion Information

Object: Car

Event 1: Car begins to accelerate.

Event 2: Collision!

$$t_1 = 0 \text{ s}$$

$$t_2 =$$

$$r_1 = 0 \text{ m}$$

$$r_2 =$$

$$v_1 = 16 \text{ m/s}$$

$$v_2 =$$

$$a_{12} = +4 \text{ m/s}^2$$

Object: Truck

Event 1: Car begins to accelerate.

$$t_1 = 0 \text{ s}$$

$$t_2 =$$

$$r_1 = 20 \text{ m}$$

$$r_2 =$$

$$v_1 = 12 \text{ m/s}$$

$$v_2 =$$

$$a_{12} = 0 \text{ m/s}^2$$

## Mathematical Analysis

At first glance, there appear to be six unknowns in the motion table. This should concern you since you only have four equations (the two kinematic equations applied to the car and the same two applied to the truck). However, since the car and truck collide at event 2,  $t_2$  and  $r_2$  for the car and truck must be equal at this event. Thus, the only four variables are  $t_2$ ,  $r_2$ ,  $v_{2\text{car}}$ , and  $v_{2\text{truck}}$ . These can be determined by the four kinematic equations. Specifically, set the position equation for the car equal to the position equation for the truck and solve for  $t_2$ :

$$\begin{aligned} r_{2\text{car}} &= r_{2\text{truck}} \\ 0 + 16(t_2 - 0) + \frac{1}{2}(4)(t_2 - 0)^2 &= 20 + 12(t_2 - 0) + \frac{1}{2}(0)(t_2 - 0)^2 \\ 16t_2 + 2t_2^2 &= 20 + 12t_2 \\ 0 &= 20 - 4t_2 - 2t_2^2 \end{aligned}$$

Using the quadratic formula,  $t_2 = 2.32 \text{ s}$ . Plugging this back into either position equation yields,

$$\begin{aligned} r_{2\text{car}} &= 0 + 16(2.32 - 0) + \frac{1}{2}(4)(2.32 - 0)^2 \\ r_{2\text{car}} &= 47.9 \text{ m} \end{aligned}$$

Solving the two velocity equations gives:

$$\begin{aligned} v_{2\text{car}} &= 16 + 4(2.32) & v_{2\text{truck}} &= 12 + 0(2.32) \\ v_{2\text{car}} &= 25.3 \text{ m/s} & v_{2\text{truck}} &= 12 \text{ m/s} \end{aligned}$$

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## I-49

A car, initially at rest, accelerates toward the west at  $2.0 \text{ m/s}^2$ . At the same time that the car starts, a truck, 350 m west of the car and moving at 16 m/s toward the east, starts to move slower, accelerating at  $1.0 \text{ m/s}^2$ . The car and truck pass safely.

### Motion Diagram



### Motion Information

Object:		Object:	
Event 1:	Event 2:	Event 1:	Event 2:
$t_1 =$	$t_2 =$	$t_1 =$	$t_2 =$
$r_1 =$	$r_2 =$	$r_1 =$	$r_2 =$
$v_1 =$	$v_2 =$	$v_1 =$	$v_2 =$
$a_{12} =$		$a_{12} =$	

### Mathematical Analysis

## I-50

A car, initially traveling at 20 m/s to the east, accelerates toward the west at  $2.0 \text{ m/s}^2$ . At the same time that the car starts, a truck, 60 m west of the car and moving at 16 m/s toward the east, starts to move faster, accelerating at  $1.0 \text{ m/s}^2$ . It's a one-lane road and both drivers are too busy fiddling with their CD players to notice each other.

### Motion Diagram



### Motion Information

Object:		Object:	
Event 1:	Event 2:	Event 1:	Event 2:
$t_1 =$	$t_2 =$	$t_1 =$	$t_2 =$
$r_1 =$	$r_2 =$	$r_1 =$	$r_2 =$
$v_1 =$	$v_2 =$	$v_1 =$	$v_2 =$
$a_{12} =$		$a_{12} =$	

### Mathematical Analysis



## I-52

A pole-vaulter lands on a cushion after a vault. Determine her acceleration as she sinks into the cushion ( $a_{\text{cushion}}$ ) as a function of her velocity when she hits the cushion ( $v_i$ ) and the distance she sinks into the cushion ( $d$ ).

### Motion Diagram



### Motion Information

Event 1:

Event 2:

$t_i =$

$t_f =$

$r_i =$

$r_f =$

$v_i =$

$v_f =$

$a_{1,2} =$

### Mathematical Analysis

### Questions

If  $v_i = 0$  m/s, what should  $a_{\text{cushion}}$  equal? Does your function agree with this observation?

If  $d = 0$  m, what should  $a_{\text{cushion}}$  equal? Does your function agree with this observation?

What would result in a larger magnitude acceleration, hitting the cushion twice as fast or sinking one-half as far into the cushion?

## I-53

The driver of an automobile suddenly sees an obstacle blocking his lane. Determine the total distance the car travels between seeing the obstacle and stopping ( $d$ ) as a function of his reaction time ( $t_{12}$ ), the initial velocity of the car ( $v_1$ ), and the magnitude of the car's acceleration while stopping ( $a$ ).

### Motion Information

Event 1:                      Event 2:                      Event 3:

$t_1 =$                        $t_2 =$                        $t_3 =$

$r_1 =$                        $r_2 =$                        $r_3 =$

$v_1 =$                        $v_2 =$                        $v_3 =$

$a_{12} =$                        $a_{23} =$

### Mathematical Analysis

#### Questions

If  $v_1 = 0$  m/s, what should  $d$  equal? Does your function agree with this observation?

If  $a = 0$  m/s<sup>2</sup>, what should  $d$  equal? Does your function agree with this observation?

Does the car travel a greater distance during the reaction phase or the braking phase? Estimate necessary quantities assuming a panic stop from highway speeds.

I-53

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## I-54

Two automobiles are involved in a race over level ground. Both cars begin at rest. The Audi accelerates at  $A$  for  $T$  seconds and then travels at constant velocity for the rest of the race. The Buick accelerates at  $A/2$  for  $2T$  seconds and then travels at constant velocity for the rest of the race. Determine the distance ( $d$ ) between the two cars, at the instant  $t = 2T$ , as a function of  $A$  and  $T$ .

### Motion Information

Object:			Object:		
Event 1:	Event 2:	Event 3:	Event 1:	Event 2:	Event 3:
$t_1 =$	$t_2 =$	$t_3 =$	$t_1 =$	$t_2 =$	$t_3 =$
$r_1 =$	$r_2 =$	$r_3 =$	$r_1 =$	$r_2 =$	$r_3 =$
$v_1 =$	$v_2 =$	$v_3 =$	$v_1 =$	$v_2 =$	$v_3 =$
$a_{12} =$	$a_{23} =$		$a_{12} =$	$a_{23} =$	

### Mathematical Analysis

#### Questions

If  $A = 0 \text{ m/s}^2$ , what should  $d$  equal? Does your function agree with this observation?

Assume the race takes longer than  $2T$  to finish. What is the difference in speed between the two cars as they cross the finish line?

Assume the race takes longer than  $2T$  to finish. What is the distance between the two cars as they cross the finish line?



## I-55

Two automobiles are involved in a game of chicken. The two cars start from rest at opposite ends of a long, one-lane road. The Audi accelerates at  $A$  for  $T$  seconds and then travels at constant velocity. The Buick accelerates at  $A/2$  for  $2T$  seconds and then travels at constant velocity. Determine the elapsed time when the cars collide ( $t_{\text{collide}}$ ) as a function of the length of the road ( $d$ ),  $A$  and  $T$ . The collision takes place after both cars have reached constant velocity.

### Motion Information

Object:			Object:		
Event 1:	Event 2:	Event 3:	Event 1:	Event 2:	Event 3:
$t_1 =$	$t_2 =$	$t_3 =$	$t_1 =$	$t_2 =$	$t_3 =$
$r_1 =$	$r_2 =$	$r_3 =$	$r_1 =$	$r_2 =$	$r_3 =$
$v_1 =$	$v_2 =$	$v_3 =$	$v_1 =$	$v_2 =$	$v_3 =$
$a_{12} =$	$a_{21} =$		$a_{12} =$	$a_{21} =$	

### Mathematical Analysis

#### Questions

If  $A = 0 \text{ m/s}^2$ , what should  $t_{\text{collide}}$  equal? Does your function agree with this observation?

If  $d = \infty$ , what should  $t_{\text{collide}}$  equal? Does your function agree with this observation?

What is the minimum length of road needed to guarantee the two cars collide at top speed?

I-55

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## 04. Dynamics

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- [01. Concepts and Principles](#)
  - 02. Drawing Free-Body Diagrams
  - 03. Calculating the Force of Gravity near the Surface of the Earth
  - 04. Analyzing a More Complex Scenario
  - 05. Hints and Suggestions
  - 06. Activities
- 

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## 01. Concepts and Principles

### 1. Newton's Laws

1. Newton's First Law
2. Newton's Second Law
3. Newton's Third Law

## Newton's Laws

### Newton's First Law

Dynamics is the study of the cause of motion, or more precisely, the cause of *changes* in motion. In the late 1600's, Isaac Newton hypothesized that motion does not require a cause, but that *changes* in motion require causes. An object experiences a change in motion only when it interacts with some aspect of its surroundings. This bold hypothesis, referred to as Newton's first law of motion, is summarized by the idea that an object will maintain its state of motion, whether at rest or traveling at high speed, unless acted upon by some aspect of its surroundings.

Using the kinematic terminology developed in the last unit, this means that an object's velocity (state of motion) is constant unless it interacts with some outside agent. An external interaction is not necessary for an object to move, it is only necessary if the object's velocity changes. Thus, what is *caused* is not velocity, but acceleration. This concept is one of the most subtle, and complex, in all of physics.

**An object will maintain its state of motion, whether at rest or traveling at high speed, unless acted upon by some aspect of its surroundings**

### Newton's Second Law

Newton also hypothesized that the sum total of all interactions with the external environment, which he termed *forces*, is directly proportional to the acceleration of the object. Moreover, the proportionality constant between the sum of all forces acting on an object and the acceleration of the object measures the "resistance" of an object to changes in its motion. This resistance to changes in motion is termed the *inertia*.

For example, an object with great inertia (quantified by a large proportionality constant) responds to the application of forces with a relatively small acceleration. An object with little inertia (a small proportionality constant) responds to the application of the same forces with a relatively large acceleration. The amount of inertia an object has is measured by the *inertial mass* of the object.

In summary, this relationship, known as Newton's Second Law of motion, and can be written mathematically as:

$$\Sigma F = ma$$

The capital Greek letter sigma,  $\Sigma$ , is used as a shorthand to remind you to sum all of the forces acting on the object. The sum of all of the forces acting on an object will be referred to as the *total force* acting on the object.

### Newton's Third Law

Newton's third great contribution to the study of dynamics was his vision of force, defined to be the interaction between an object and some aspect of its surroundings. Newton theorized that since objects *interact* with other objects in their surroundings, always in pairs, a certain symmetry exists in nature. The distinction between the actor and the acted-upon is arbitrary. It would be just as easy to switch focus and consider the object in the surroundings as the acted-upon and the original object of interest the actor.

If nature exhibits this symmetry, then the force that one object exerts on another must *always* be equal in magnitude to the force that the second object exerts on the first. To speak of one object as exerting a force on another is to speak of only one-half of the picture. This idea, known as Newton's Third Law of motion, is of central importance in the study of forces. In summary, objects *interact* with each other, and equal magnitude forces are exerted on each of the two objects interacting. A simplistic way of picturing this is the idea that you cannot touch something without being touched, and moreover, that the harder you touch, the harder you will be touched in return.

Investigating the dynamics of a situation involves the identification of all interactions an object experiences with other objects in its surroundings. To help in the identification of these interactions, and to use this information to better describe the ultimate motion of

the object, a number of useful analysis tools are detailed below.

**Forces come in pairs and are always equal in magnitude, but opposite in direction**

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## 02. Drawing Free-Body Diagrams

1. 1. Select the object you would like to study
2. 2. Draw a picture of the object of interest free from all other objects.
3. 3. Indicate on the diagram all interactions of the object with its environment.
  1. Non-Contact Interactions
  2. Contact Interactions

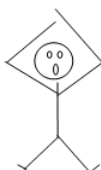
The free-body diagram is by far the most important analysis tool for determining the interactions between an object and its surroundings. There are three distinct steps to creating a free-body diagram. Let's walk through the steps for the situation described below:

*A child pulls herself up a rope using only her hands.*

### 1. Select the object you would like to study

In this example, it is probably safe to assume that the object we would like to study is the child. However, depending on what we are investigating it may be the rope or even the ceiling we are interested in. Selecting the correct object to represent by a free-body diagram is a crucial step, especially in more complicated situations. With practice you will develop a knack for selecting the correct object to represent.

### 2. Draw a picture of the object of interest *free from all other objects*.



Notice that the rope does not appear in the diagram. As the name free-body implies, the object is drawn free of all external constraints.

### 3. Indicate on the diagram all interactions of the object with its environment.

Now comes the most difficult part of constructing a free-body diagram. It is crucial not to miss an interaction. If an interaction is overlooked, then the total of the forces will be incorrect, and the acceleration will be incorrect, and your entire analysis will be incorrect.

Also, only the portion of the interaction that acts *on the girl* should be indicated on a free-body diagram of the girl. For example, she is interacting with the rope. The rope's action *on the girl* will be indicated, not the action of the girl *on the rope*.

To aid in the search for interactions, we will divide the types of interactions that the girl can be part of into two types, non-contact and contact.

#### Non-Contact Interactions

Non-contact interactions include all interactions that can occur between the girl and objects in her surroundings that do not require direct physical contact between the two objects. Non-contact interactions include the interaction of the girl with the gravitational and electromagnetic fields in her vicinity. (How these fields are created and how they can effect the girl will slowly be incorporated into our physics model.)

At the current level of complexity, however, the only non-contact interaction you need worry about is the interaction of the girl with the gravitational field created by the earth, which we will simply term the force of gravity. The direction of this force is down, toward the center of the earth.

#### Contact Interactions

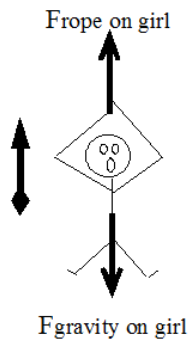
Contact interactions occur at every point on the girls' body in which she is in direct physical contact with an external object. The most obvious of these is the rope. The girl is in contact with the rope, so the rope and girl exert forces on each other. These forces

are equal in magnitude. Remember, however, that it is only the force exerted *on* the girl that is indicated on a free-body diagram of the girl. The location of this force is at the girls' hands, and the direction of this force is up. (The direction of this force cannot be down, because that would imply that the rope is *pushing* the girl, as opposed to *pulling* her. It is impossible for a rope to push someone, unless it is a very stiff rope. Very stiff ropes will be called rods.)

The only other objects to actually make contact with the girl are air molecules. The air molecules interact with the girl on all sides, each exerting a small force directly inward, perpendicular to the girl's body. Although each of these forces is very small, their sum is not always small. For example, if the girl were falling freely from an airplane the vast numbers of air molecules colliding with the girl from underneath, versus the rather small number colliding from above, and the strength of these collisions, would add to a very large force acting upward on the girl. This force could easily be equal in magnitude to the force of gravity on the girl. The force of air molecules on an object, referred to as air resistance, is often ignored in analyzing scenarios simply because of the difficulty of dealing with the complexity.

However, in many cases the effects of the air molecules are negligible compared to the other forces acting on the object. This is the case with the girl climbing the rope. The forces exerted by the air molecules are probably very close to being uniformly distributed around the girl's surface. Thus, for every air molecule pushing her to the right, there is probably an air molecule pushing her to the left. These forces will add to a total force very close to zero.

A correct free-body diagram for the girl is shown below:



Since a coordinate system is crucial for translating motion diagrams and free-body diagrams into mathematical relationships, a coordinate system has been added to the free-body diagram. It is always a good idea to use the same coordinate system for both the free-body diagram and the motion diagram.

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## 03. Calculating the Force of Gravity near the Surface of the Earth

1. Calculating the Force of Gravity near the Surface of the Earth
2. Applying Newton's Second Law

### Calculating the Force of Gravity near the Surface of the Earth

In addition to creating the three laws of motion mentioned earlier, Newton also postulated the Law of Universal Gravitation. This law states that every object of mass in the universe creates a gravitational field, and every object of mass in the universe senses and interacts with every other objects' field. That's an awful lot of forces! To try to identify and estimate the magnitude of all of these forces on an object near the surface of the earth would be a lifelong task.

Luckily, the strength of the gravitational field depends on the mass of the object producing the field. The more massive the object, the stronger the field. Thus, since the earth is much more massive than any other nearby object, when creating free-body diagrams for objects near the surface of the earth we can safely include just the gravitational force due to the earth, ignoring all the other, relatively small, gravitational forces.

The magnitude of the gravitational field of a massive object, denoted  $g$ , depends on the mass of the object, the distance from the center of the object, and a constant called, appropriately, the gravitational constant,  $G$ . The relationship is:

$$g = \frac{GM}{d^2} \quad [\text{N/kg}]$$

Near the surface of the earth, the gravitational field has a magnitude of approximately 9.8 N/kg. Although the gravitational field strength varies with the distance from the surface of the earth, we will ignore this slight variation unless explicitly told to include its effects.

The gravitational force felt by a massive object in the presence of a gravitational field is given by the product of the object's mass and the magnitude of the gravitational field at the location of the object:

$$F_{\text{gravity}} = mg \quad [\text{N}]$$

Forces are measured in Newtons [N], where,

$$N = \frac{\text{kg} \cdot \text{m}}{\text{s}^2}$$

### Applying Newton's Second Law

Let's return to the scenario under investigation and make some quantitative information more explicit. Then, we can attempt to further investigate the situation using Newton's Second Law.

***A 30 kg child pulls herself up a rope at an approximately constant speed using only her hands.***

Newton's Second Law states:

$$\Sigma F = ma$$

$\Sigma F$  refers to the sum of all of the forces acting on the girl, the force of the rope (which is positive in our coordinate system) and the force of gravity (which is negative in our coordinate system). Thus,

$$F_{\text{rope}} - F_{\text{gravity}} = ma$$

Since  $m = 30 \text{ kg}$ , and  $a = 0 \text{ m/s}^2$  (since she climbs at constant speed), the equation becomes:

$$F_{\text{rope}} - F_{\text{gravity}} = 0$$

By Newton's relationship for the force of gravity:

$$F_{gravity} = mg$$

$$F_{gravity} = (30kg)(9.8 \frac{N}{kg})$$

$$F_{gravity} = 294N$$

Therefore:

$$F_{rope} - 294N = 0$$

$$F_{rope} = 294N$$

Thus, Newton's Second Law allows us to determine the force with which the rope pulls on the girl. Of course, by Newton's Third Law, the force with which the girl pulls on the rope is equal in magnitude, so the girl exerts a 294 N force on the rope.

If the girl had not climbed the rope at approximately constant speed her acceleration would have to be determined, either from an explicit mention in the description or through using the kinematic relations developed in the last unit, and then inserted into Newton's Second Law. If her acceleration had been directed upwards (positive) the force of the rope on the girl would have had to be larger. If her acceleration had been directed downwards (negative) the force of the rope on the girl would have had to be smaller.

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## 04. Analyzing a More Complex Scenario

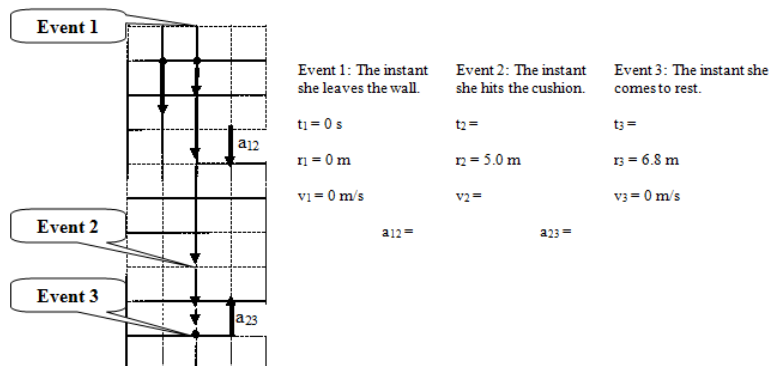
### 1. Analyzing a More Complex Scenario

### Analyzing a More Complex Scenario

Before you start analyzing dynamics scenarios on your own, let's work our way through a more complex scenario.

To practice falling, a 55 kg pole-vaulter falls from rest off of a wall 5.0 m above a foam cushion. The pole-vaulter sinks about 1.8 m into the cushion before stopping.

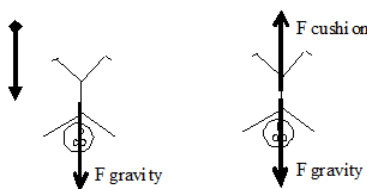
Before we begin analyzing the forces acting on this pole-vaulter, I think we should try to get a handle on the kinematics of the situation. Therefore, our first step in analyzing this situation is to draw a motion diagram and tabulate motion information.



Notice that between the instant she leaves the wall and the instant she hits the cushion the acceleration is positive (down), while between the instant she hits the cushion and the instant she comes to rest the acceleration is negative (up). Thus, when applying the kinematic relationships and Newton's Second Law we will have to be careful not to confuse variables between these two intervals.

What should jump out at you is the fact that this kinematic scenario can not be solved! There are *five* unknown kinematic quantities and only *four* kinematic equations. Something else is needed in order to complete the kinematic description. Let's look at the forces acting on the pole-vaulter to see if we can figure out another piece of kinematic information.

Between the first two instants, the only force acting on the pole-vaulter is the force of gravity. Once she hits the cushion, however, there are two forces acting on the pole-vaulter, the force of gravity and the force of the cushion. Correct free-body diagrams for these *two distinct phases of her motion* are given below.



The diagram on the left corresponds to the first time interval and the diagram on the right to the second time interval. For *each* of these free-body diagrams, I will apply Newton's Second Law:

$$\begin{aligned} \Sigma F &= ma_{12} \\ + F_{\text{gravity}} &= (55\text{kg})a_{12} \end{aligned} \qquad \begin{aligned} \Sigma F &= ma_{23} \\ + F_{\text{gravity}} - F_{\text{cushion}} &= (55\text{kg})a_{23} \end{aligned}$$

But

$$\begin{aligned} F_{\text{gravity}} &= mg \\ F_{\text{gravity}} &= (55\text{kg})(9.8 \frac{\text{N}}{\text{kg}}) \\ F_{\text{gravity}} &= 539\text{N} \end{aligned}$$

$$\begin{aligned} 539\text{N} &= (55\text{kg})a_{12} \\ a_{12} &= 9.8\text{m/s}^2 \end{aligned} \qquad \begin{aligned} 539\text{N} - F_{\text{cushion}} &= (55\text{kg})a_{23} \end{aligned}$$

Thus, from Newton's Second Law, we know that the acceleration *during the fall* is  $9.8 \text{ m/s}^2$ . (We still don't know what the acceleration was during the impact portion of the motion.) Substituting this value back into the motion table yields:

Event 1: The instant she leaves the wall.	Event 2: The instant she hits the cushion.	Event 3: The instant she comes to rest.
$t_1 = 0 \text{ s}$	$t_2 =$	$t_3 =$
$r_1 = 0 \text{ m}$	$r_2 = 5.0 \text{ m}$	$r_3 = 6.8 \text{ m}$
$v_1 = 0 \text{ m/s}$	$v_2 =$	$v_3 = 0 \text{ m/s}$
$a_{12} = 9.8 \text{ m/s}^2$	$a_{23} =$	

This is now solvable, using strictly kinematics, for the four remaining unknowns. Try to do the math on your own, and compare your result to:

Event 1: The instant she leaves the wall.	Event 2: The instant she hits the cushion.	Event 3: The instant she comes to rest.
$t_1 = 0 \text{ s}$	$t_2 = 1.0 \text{ s}$	$t_3 = 1.36 \text{ s}$
$r_1 = 0 \text{ m}$	$r_2 = +5.0 \text{ m}$	$r_3 = +6.8 \text{ m}$
$v_1 = 0 \text{ m/s}$	$v_2 = +9.9 \text{ m/s}$	$v_3 = 0 \text{ m/s}$
$a_{12} = 9.8 \text{ m/s}^2$	$a_{23} = -27 \text{ m/s}^2$	

We now have a complete kinematic description of the motion.

Returning to Newton's Second Law for the impact portion of the motion,

$$\begin{aligned}
 539 \text{ N} - F_{\text{cushion}} &= (55 \text{ kg})a_{23} \\
 539 \text{ N} - F_{\text{cushion}} &= (55 \text{ kg})(-27 \text{ m/s}^2) \\
 539 \text{ N} - F_{\text{cushion}} &= -1485 \text{ N} \\
 F_{\text{cushion}} &= 2024 \text{ N}
 \end{aligned}$$

The cushion exerts a force of about 2000 N on the pole-vaulter to stop her fall.

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## 05. Hints and Suggestions

### 1. The Magnitude of the Gravitational Field

1. 'g' is never negative.
2. 'g' is not an acceleration.

### 2. Newton's Third Law

## The Magnitude of the Gravitational Field

Quite often, students make a pair of mistakes when dealing with the magnitude of the gravitational field,  $g$ .

### 'g' is never negative.

Since  $g$  is the *magnitude* of the gravitational field, it cannot be a negative number. As a magnitude, it *does not have a direction associated with it*! Resist all temptation to replace 'g' with the value "-9.8 N/kg"!

Part of the confusion lies with the fact that the gravitational *field* does have an associated direction. The gravitational field of the earth is directed downward toward the center of the earth. Even so, the gravitational field is *not* negative. Negative only makes sense relative to a coordinate system, and since you are always free to choose any system you want, the gravitational field is just as likely to be oriented in the positive as the negative direction.

### 'g' is not an acceleration.

Often times, students have the idea that 'g' is the "acceleration due to gravity". However, as I sit here in a chair writing this book, the force of gravity *is* acting on me and I am most definitely *not* accelerating at  $9.8 \text{ m/s}^2$ ! In fact, the force of gravity has acted on me for every second of my life and only very rarely have I accelerated at  $9.8 \text{ m/s}^2$ . 'g' measures the strength of the gravitational field. As such, it is related to the gravitational force, which, like all forces, can give rise to accelerations. However, it is the *total* force acting on an object that determines its acceleration, not simply the force of gravity.

It is true that the units of 'g', N/kg, are also the units of acceleration, since a Newton is defined to be a  $\text{kg m/s}^2$ . It is also true that in a *very specific scenario*, when the only force acting on an object is the force of gravity, the magnitude of the object's acceleration is numerically equal to 'g'. However, there are also very specific scenarios in which the acceleration of an object is numerically equal to  $4.576 \text{ m/s}^2$ , or  $62.31452 \text{ m/s}^2$ . The strength of physics is its ability to analyze diverse scenarios with the same small set of tools, not to develop specialized tools tailored to every different specific scenario. Newton's Second Law will always allow you to determine an object's acceleration, whether the force of gravity acts alone or not.

## Newton's Third Law

Many physics students have heard the saying, "For every action there is an equal and opposite re-action." I was forced to memorize this statement in a middle-school science class, and was told it was called Newton's Third Law. I had no idea what it meant then, and I still am not very clear about what it has to do with Newton's Third Law.

Template:ExampleStart

As a test of your understanding of Newton's Third Law, try to answer the following question:

As you drive along the highway, a mosquito splats against your car windshield. During the collision between the mosquito and the car:

- a. the force on the mosquito was greater in magnitude than the force on the car.
- b. the force on the car was greater in magnitude than the force on the mosquito.
- c. the force on the mosquito was equal in magnitude to the force on the car.
- d. it is impossible to determine the relative sizes of the forces without more information.

As strange as it may seem, the correct answer is 'c'. The forces exerted on the mosquito and the car are equal in magnitude. In the terminology used in this chapter, the mosquito and car interact (probably an unpleasant interaction for the mosquito), and in an interaction the two agents involved exert equal forces on each other.

However, obviously *something* is different about the interaction from the mosquito's perspective. What is different is not the force acting on the mosquito but rather its acceleration. Although the forces acting on the mosquito and car are the same, the mosquito's

acceleration is *much greater* than the car's acceleration because the mosquito's mass is *much smaller* than the car's mass. The acceleration of the car is so small that it is not even noticed by the driver, while the acceleration of the mosquito is certainly noticed by the mosquito!

$$F_{car} = F_{mosquito}$$
$$M_{car} \cdot a_{car} = m_{mosquito} \cdot A_{mosquito}$$

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## 06. Activities

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## 01. Selected Answers For Dynamics

### Selected Answers for ModelSpiral Physics: Mechanics

#### Kinematics

- 41.  $r_2 = 1.4 \text{ m}$
- 42.  $t_2 = 1.45 \text{ s}$
- 43.  $t_3 = 3.4 \text{ s}$
- 44.  $t_3 = 4.55 \text{ s}$
- 45.  $r_3 = 36.6 \text{ m}$
- 46.  $r_4 = 2.0 \text{ km}$
- 47.  $t_4 = 370 \text{ ks}$
- 49.  $t_2 = 14.9 \text{ s}$
- 50.  $t_2 = 7.8 \text{ s}$
- 51.  $t_3 = 2.87 \text{ s}$

#### Dynamics

- 80.  $F_{\text{rope}} = 420 \text{ N}$
- 81.  $F_{\text{cushion}} = 1.91 \text{ kN}$
- 83.  $F_{\text{bottom cable}} = 11.4 \text{ N}$
- 84.  $F_{\text{cushion}} = 2.83 \text{ kN}$
- 85.  $F_{\text{ground}} = 43.7 \text{ kN}$
- 86.  $r_3 = 63.5 \text{ m}$
- 87.  $r_3 = 1.55 \text{ km}$
- 89.  $m_{\text{block}} = 240 \text{ kg}$
- 90.  $m_{\text{block}} = 26 \text{ kg}$
- 91.  $F_{\text{rope}} = 500 \text{ N}$
- 92.  $F_{\text{rope}} = 490 \text{ N}$

#### Conservation Laws

- 122. a.  $F_{\text{scale}} = 755 \text{ N}$       b.  $F_{\text{scale}} = 780 \text{ N}$
- 123. a.  $v = 116 \text{ m/s}$       b.  $t = 26.8 \text{ s}$
- 124. a.  $v = 12.5 \text{ m/s}$       b.  $F_{\text{ground}} = 43.7 \text{ kN}$
- 125.  $t_2 = 0.89 \text{ s}$        $r_2 = 1.78 \text{ m}$
- 127.  $m_{\text{student}} = 94 \text{ kg}$
- 128.  $v_2 = 6.0 \text{ m/s}$
- 129.  $v_2 = 1.5 \text{ m/s}$
- 130.  $v_{2 \text{ ship}} = 24.3 \text{ m/s}$
- 131.  $v_{\text{probe}} = 667 \text{ m/s}$
- 132.  $v_{2 \text{ platform}} = 0.28 \text{ m/s}$
- 133.  $v_{2 \text{ platform}} = 0.06 \text{ m/s}$
- 134.  $v_{2 \text{ balloon}} = 6.3 \text{ m/s}$

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## I-68

Construct free-body diagrams for the objects described below.

a. When throwing a ball vertically upward, my hand moves through a distance of about 1.0 m before the ball leaves my hand. The 0.80 kg ball reaches a maximum height of about 20 m above my hand.

while the ball is in my hand

after the ball leaves my hand



b. To practice falling, a 55 kg pole-vaulter falls off of a wall 6.0 m above a 2.0 m thick foam cushion resting on the ground. However, he misses the cushion.

while falling through the air

while being stopped by the ground



c. A decorative light fixture in an elevator consists of a 2.0 kg light suspended by a cable from the ceiling of the elevator. From this light, a separate cable suspends a second 0.80 kg light.

the top light

the bottom light



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## I-69

Construct free-body diagrams for the objects described below.

a. A 4000 kg rocket's engine produces a thrust of 70,000 N for 15 s. The rocket is fired vertically upward.

while the engine is firing

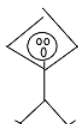


after the engine turns off



b. Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching his third floor dorm room. An 84 kg block is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the 84 kg block is released, the student is pulled up to his dorm room.

the student



the block



c. A 1.0 kg block is stacked on top of a 2.0 kg block on the floor of an elevator moving downward at constant speed.

the 1.0 kg block



the 2.0 kg block



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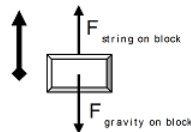


## I-70

A block hangs from the ceiling of an elevator by a string. For each of the following situations, circle the correct relationship symbol between the force of the string on the block and the force of gravity on the block and explain your reasoning.

a. The elevator is at rest.

$F_{\text{string on block}}$       >      ☒      <      ?       $F_{\text{gravity on block}}$   
 Explanation:



Since the block is not accelerating, the two forces acting on it must be equal in magnitude.

b. The elevator is moving upward at a constant speed.

$F_{\text{string on block}}$       >      ☒      <      ?       $F_{\text{gravity on block}}$   
 Explanation:

Since the block is still not accelerating, the two forces acting on it must be equal in magnitude.

c. The elevator is moving downward at a decreasing speed.

$F_{\text{string on block}}$       >      ☐ =      <      ?       $F_{\text{gravity on block}}$   
 Explanation:

Since the block is accelerating upward, the force directed upward (the force of the string) must be larger than the force directed downward (the force of gravity).

d. The elevator is moving upward at an increasing speed.

$F_{\text{string on block}}$       >      ☐ =      <      ?       $F_{\text{gravity on block}}$   
 Explanation:

The block is accelerating upward, so the force directed upward must be larger than the force directed downward.

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## I-71

A man stands on a bathroom scale inside of an elevator. For each of the following situations, circle the correct relationship symbol between the force of the scale on the man and the force of gravity on the man and explain your reasoning.

a. The elevator is at rest.

$F_{\text{scale on man}}$        $>$        $=$        $<$        $?$        $F_{\text{gravity on man}}$   
Explanation:

b. The elevator is moving downward at a constant speed.

$F_{\text{scale on man}}$        $>$        $=$        $<$        $?$        $F_{\text{gravity on man}}$   
Explanation:

c. The elevator is moving downward at an increasing speed.

$F_{\text{scale on man}}$        $>$        $=$        $<$        $?$        $F_{\text{gravity on man}}$   
Explanation:

d. The elevator is moving upward at a decreasing speed.

$F_{\text{scale on man}}$        $>$        $=$        $<$        $?$        $F_{\text{gravity on man}}$   
Explanation:

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## I-72

Two blocks are stacked on top of each other on the floor of an elevator. For each of the following situations, circle the correct relationship symbol between the two forces and explain your reasoning.

a. The elevator is moving downward at a constant speed.

$F_{\text{bottom block on top block}}$  > = < ?  $F_{\text{top block on bottom block}}$   
Explanation:

$F_{\text{bottom block on top block}}$  > = < ?  $F_{\text{gravity on top block}}$   
Explanation:

b. The elevator is moving downward at an increasing speed.

$F_{\text{bottom block on top block}}$  > = < ?  $F_{\text{top block on bottom block}}$   
Explanation:

$F_{\text{bottom block on top block}}$  > = < ?  $F_{\text{gravity on top block}}$   
Explanation:

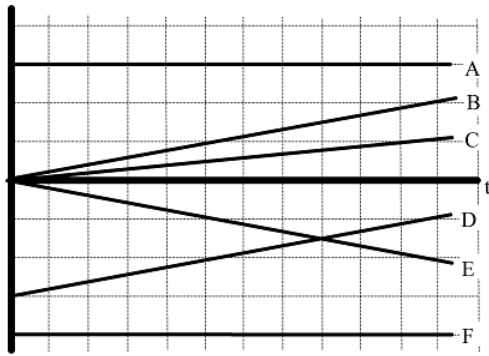
c. The elevator is moving upward at a decreasing speed.

$F_{\text{bottom block on top block}}$  > = < ?  $F_{\text{top block on bottom block}}$   
Explanation:

$F_{\text{bottom block on top block}}$  > = < ?  $F_{\text{gravity on top block}}$   
Explanation:

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## I-73



a. If the graph is of position vs. time, rank these graphs on the basis of the total force acting on the object. Rank a positive total force as larger than a negative total force.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. If the graph is of velocity vs. time, rank these graphs on the basis of the total force acting on the object. Rank a positive total force as larger than a negative total force.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

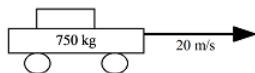
Explain the reason for your ranking:

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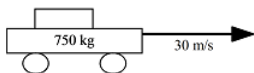
## I-74

Below are six automobiles traveling at constant velocity. The automobiles have different masses and velocities. Rank these automobiles on the basis of the total force acting on them.

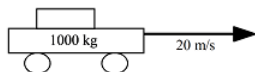
**A**



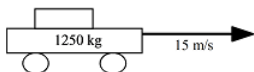
**B**



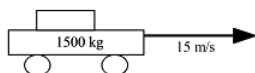
**C**



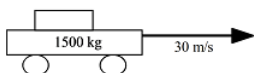
**D**



**E**



**F**



Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.


Explain the reason for your ranking:

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## I-75

Below are six identical baseballs shortly after being thrown. At the instant shown, the baseball's velocity is indicated, along with the distance the ball has traveled and the elapsed time since leaving the thrower's hand. Rank these baseballs on the basis of the force of the thrower's hand currently acting on them.


**A**

10 m  
0.2 s  30 m/s


**B**

15 m  
0.2 s  50 m/s


**C**

10 m  
0.2 s  40 m/s


**D**

20 m  
0.4 s  40 m/s

**E**

15 m  
0.4 s  30 m/s

**F**

5.0 m  
0.1 s  50 m/s

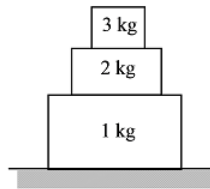
Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest  
\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

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## I-76

Below are three blocks stacked on top of each other at rest. Rank the magnitude of the forces referred to below from largest to smallest.



- A** The force of the 3 kg block on the 2 kg block
- B** The force of the 2 kg block on the 3 kg block
- C** The force of the 3 kg block on the 1 kg block
- D** The force of the 1 kg block on the 3 kg block
- E** The force of the 2 kg block on the 1 kg block
- F** The force of the 1 kg block on the 2 kg block
- G** The force of the 1 kg block on the floor
- H** The force of the floor on the 1 kg block

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

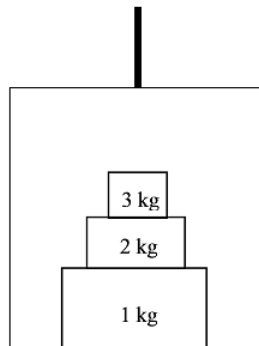
\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

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## I-77

Below are three blocks stacked on top of each other inside of an elevator moving upward at increasing speed. Rank the magnitude of the forces referred to below from largest to smallest.



- A** The force of the 3 kg block on the 2 kg block
- B** The force of the 2 kg block on the 3 kg block
- C** The force of the 3 kg block on the 1 kg block
- D** The force of the 1 kg block on the 3 kg block
- E** The force of the 2 kg block on the 1 kg block
- F** The force of the 1 kg block on the 2 kg block
- G** The force of the 1 kg block on the floor of the elevator
- H** The force of the floor of the elevator on the 1 kg block

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

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## I-78

For each of the collisions illustrated below, circle the correct relationship symbol between the force of car A on car B and the force of car B on car A and explain your reasoning.

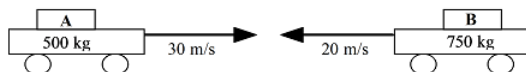
a.



$F_{\text{car A on car B}}$  > = < ?  $F_{\text{car B on car A}}$

Explanation:

b.



$F_{\text{car A on car B}}$  > = < ?  $F_{\text{car B on car A}}$

Explanation:

c.



$F_{\text{car A on car B}}$  > = < ?  $F_{\text{car B on car A}}$

Explanation:

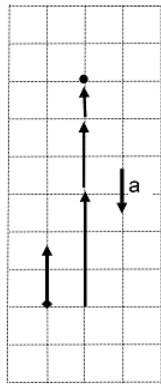
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# I-79

A 100 kg man concerned about his weight decides to weigh himself in an elevator. He stands on a bathroom scale in an elevator that is moving upward at 3.0 m/s. As the elevator reaches his floor, it slows to a stop over a distance of 2.0 m.

## Motion Diagram

### Free-Body Diagram



## Motion Information

Event 1: The instant the elevator begins to slow.

$$t_1 = 0 \text{ s}$$

$$r_1 = 0 \text{ m}$$

$$v_1 = 3.0 \text{ m/s}$$

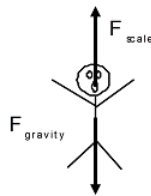
$$a_{12} =$$

Event 2: The instant it stops.

$$t_2 =$$

$$r_2 = +2.0 \text{ m}$$

$$v_2 = 0 \text{ m/s}$$



## Mathematical Analysis

Since there are only two unknown kinematic quantities, we can determine them by our two kinematic equations.

$$v_2 = v_1 + a_{12}(t_2 - t_1)$$

$$0 = 3 + a_{12}t_2 - 0$$

$$a_{12} = \frac{-3}{t_2}$$

Now substitute this expression into the other equation:

$$a_{12} = \frac{-3}{1.33}$$

$$a_{12} = -2.25 \text{ m/s}^2$$

$$r_2 = r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2$$

$$2 = 3t_2 + \frac{1}{2}a_{12}t_2^2$$

$$2 = 3t_2 + \frac{1}{2}\left(\frac{-3}{t_2}\right)t_2^2$$

$$2 = 3t_2 - 1.5t_2$$

$$2 = 1.5t_2$$

$$t_2 = 1.33 \text{ s}$$

Substitute this result back into the original equation:

Now apply Newton's Second Law to the man:

$$\Sigma F = ma$$

$$F_{\text{scale}} - F_{\text{gravity}} = (100)(-2.25)$$

$$F_{\text{scale}} - (100)(9.8) = -225$$

$$F_{\text{scale}} = 755 \text{ N}$$

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## I-80

A 40 kg child is hanging from a rope by her hands. She exerts a burst of strength and 2.0 s later is traveling at 1.4 m/s up the rope.

### Motion Diagram

#### Free-Body Diagram



Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

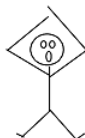
$a_{12} =$

Event 2:

$t_2 =$

$r_2 =$

$v_2 =$



### Mathematical Analysis

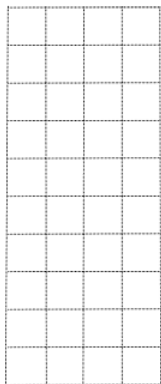
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## I-81

A 55 kg pole-vaulter, just before touching the cushion on which she lands after a jump, is falling downward at a speed of 10 m/s. The pole-vaulter sinks about 2.0 m into the cushion before stopping.

### Motion Diagram

#### Free-Body Diagram



Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

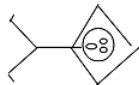
$a_{12} =$

Event 2:

$t_2 =$

$r_2 =$

$v_2 =$



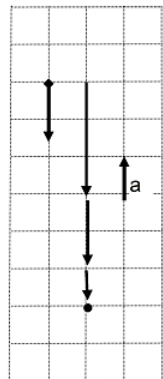
### Mathematical Analysis

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## I-82

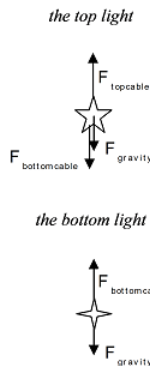
A decorative light fixture in an elevator consists of a 2.0 kg light suspended by a cable from the ceiling of the elevator. From this light, a separate cable suspends a second 0.80 kg light. The elevator is moving downward at 4.0 m/s when someone presses the emergency stop button. The elevator comes to rest in 1.2 seconds.

### Motion Diagram Free-Body Diagrams



### Motion Information

Event 1: The stop button is pressed.	Event 2: The elevator stops.
$t_1 = 0 \text{ s}$	$t_2 = 1.2 \text{ s}$
$r_1 = 0 \text{ m}$	$r_2 =$
$v_1 = 4.0 \text{ m/s}$	$v_2 = 0 \text{ m/s}$
	$a_{12} =$



### Mathematical Analysis

Since there are only two unknown kinematic quantities, we can determine them by our two kinematic equations. Note that both lights have the same kinematic description.

$$0 = 4 + a_{12}(1.2 - 0)$$

$$a_{12} = -3.33 \text{ m/s}^2$$

$$r_2 = 0 + 4(1.2 - 0) + \frac{1}{2}(-3.33)(1.2 - 0)^2$$

$$r_2 = 2.4 \text{ m}$$

Now apply Newton's Second Law to the two lights:

<p>bottom light</p> $-F_{\text{bottomcable}} + F_{\text{gravity}} = ma$ $-F_{\text{bottomcable}} + (0.8)(9.8) = (0.8)(-3.33)$ $F_{\text{bottomcable}} = 10.5 \text{ N}$	<p>top light</p> $-F_{\text{topcable}} + F_{\text{bottomcable}} + F_{\text{gravity}} = ma$ $-F_{\text{topcable}} + 10.5 + (2.0)(9.8) = (2.0)(-3.33)$ $F_{\text{topcable}} = 36.8 \text{ N}$
---	---

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## I-83

A decorative light fixture in an elevator consists of a 2.0 kg light suspended by a cable from the ceiling of the elevator. From this light, a separate cable suspends a second 0.80 kg light. The elevator is moving downward at 4.0 m/s when someone presses the emergency stop button. During the stop, the upper cable snaps. The elevator engineer says that the cable could withstand a force of 40 N without breaking.

### Motion Diagram

#### Free-Body Diagrams



### Motion Information

Event 1:

Event 2:

the top light

$t_1 =$

$t_2 =$



$r_1 =$

$r_2 =$

$v_1 =$

$v_2 =$

$a_{12} =$

the bottom light



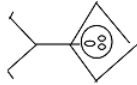

### Mathematical Analysis

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## I-84

To practice falling, a 55 kg pole-vaulter falls off of a wall 6.0 m above a foam cushion. The pole-vaulter sinks about 1.4 m into the cushion before stopping.

### Motion Diagram Free-Body Diagrams

	Event 1:	Event 2:	Event 3:	
				<i>before hitting cushion</i>
	$t_1 =$	$t_2 =$	$t_3 =$	
	$r_1 =$	$r_2 =$	$r_3 =$	
	$v_1 =$	$v_2 =$	$v_3 =$	
	$a_{12} =$	$a_{23} =$		<i>after hitting cushion</i>
				

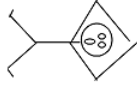

### Mathematical Analysis

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## I-85

To practice falling, a 55 kg pole-vaulter falls off of a wall 6.0 m above a 2.0 m thick foam cushion resting on the ground. However, he misses the cushion. The pole-vaulter sinks about 0.10 m into the ground before stopping.

### Motion Diagram Free-Body Diagrams

	Event 1:	Event 2:	Event 3:	
	$t_1 =$	$t_2 =$	$t_3 =$	<i>before hitting cushion</i>
	$r_1 =$	$r_2 =$	$r_3 =$	
	$v_1 =$	$v_2 =$	$v_3 =$	<i>after hitting ground</i>
	$a_{12} =$	$a_{23} =$		

### Mathematical Analysis

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## I-86

When throwing a 0.80 kg ball vertically upward, my hand moves through a distance of about 1.0 m before the ball leaves my hand. The ball leaves my hand at 35 m/s.

### Motion Diagram      Motion Information

#### Free-Body Diagrams

	Event 1:	Event 2:	Event 3:	
				<i>before ball leaves hand</i>
	$t_1 =$	$t_2 =$	$t_3 =$	
	$r_1 =$	$r_2 =$	$r_3 =$	
	$v_1 =$	$v_2 =$	$v_3 =$	
	$a_{12} =$	$a_{23} =$		<i>after ball leaves hand</i>
				

#### Mathematical Analysis

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## 1-87

A 4000 kg rocket's engine produces a thrust of 70,000 N for 15 s. The rocket is fired vertically upward.

## Motion Diagram      Motion Information

### Free-Body Diagrams

[illegible]

Event 1:

Event 2:

Event 3:

*before engine turns off*

$$t_1 =$$
$$t_2 =$$
$$t_3 =$$
$$\mathbf{r}_1 =$$
$$\mathbf{r}_2 =$$
$$\Gamma_3 =$$
$$\mathbf{v}_1 =$$
 $V_2 =$  $V_3 =$ 
$$\mathbf{a}_{12} =$$
$$\mathbf{a}_{23} =$$


*after engine turns off*



## Mathematical Analysis

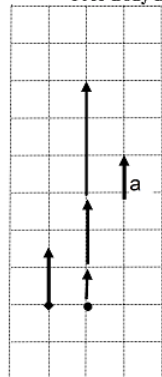
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# I-88

Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching his third floor dorm room. An 84 kg block is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the 84 kg block is released, the student is pulled up to his dorm room, 8.0 m off the ground.

## Motion Diagram

### Free-Body Diagrams



## Motion Information

Object: Student

Event 1: Block is released. Event 2: Student reaches room.

$$t_1 = 0 \text{ s}$$

$$t_2 =$$

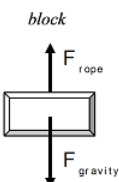
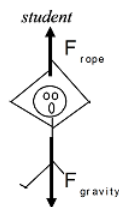
$$r_1 = 0 \text{ m}$$

$$r_2 = +8.0 \text{ m}$$

$$v_1 = 0 \text{ m/s}$$

$$v_2 =$$

$$a_{12} =$$



## Mathematical Analysis

Since there are three kinematic variables, we will have to analyze the forces first:

student

$$F_{\text{rope}} - F_{\text{gravity}} = ma$$

$$F_{\text{rope}} - (80)(9.8) = 80a_{\text{student}}$$

block

$$F_{\text{rope}} - F_{\text{gravity}} = ma$$

$$F_{\text{rope}} - (84)(9.8) = 84a_{\text{block}}$$

Because they are tied together, the acceleration of the student and the acceleration of the block are equal in magnitude, but opposite in direction.

Therefore,  $a_{\text{block}} = -a_{\text{student}}$ .

$$F_{\text{rope}} - 784 = 80a_{\text{student}}$$

$$F_{\text{rope}} = 80a_{\text{student}} + 784$$

$$F_{\text{rope}} - 823 = 84(-a_{\text{student}})$$

$$(80a_{\text{student}} + 784) - 823 = -84a_{\text{student}}$$

$$164a_{\text{student}} = 39$$

$$a_{\text{student}} = 0.24 \text{ m/s}^2$$

We can now complete the kinematic description of the student's motion:

$$8 = 0 + 0(t_2 - 0) + \frac{1}{2}(0.24)(t_2 - 0)^2$$

$$t_2 = 8.18 \text{ s}$$

$$v_2 = 0 + 0.24(8.18 - 0)$$

$$v_2 = 1.96 \text{ m/s}$$

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## I-89

Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching his third floor dorm room. A block is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the block is released, the student is pulled up to his dorm room, 8 m off the ground, in a time of 1.8 s.

### Motion Diagram

#### Free-Body Diagrams



Object:

Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

$a_{12} =$

### Motion Information

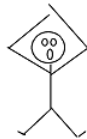
Event 2:

$t_2 =$

$r_2 =$

$v_2 =$

student



block



### Mathematical Analysis

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## I-90

Tired of walking down the stairs, an 80 kg engineering student designs an ingenious device for reaching the ground from his third floor dorm room. A block, at rest on the ground, is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the student steps out of the window, she falls the 8 m to the ground in a time of 1.8 s.

### Motion Diagram

#### Free-Body Diagrams



Object:

Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

$a_{12} =$

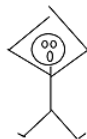
Event 2:

$t_2 =$

$r_2 =$

$v_2 =$

student



block

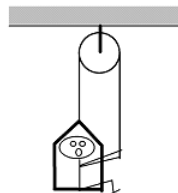


### Mathematical Analysis

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## I-91

A 60 kg student lifts herself from rest to a speed of 1.5 m/s in 2.1 s. The chair has a mass of 35 kg.



### Motion Diagram

#### Free-Body Diagrams



Object:

Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

$a_{12} =$

Event 2:

$t_2 =$

$r_2 =$

$v_2 =$

*student*



*chair*

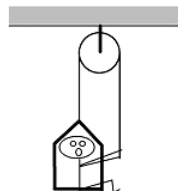


### Mathematical Analysis

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## I-92

A 60 kg student lowers herself down 40 m at a constant speed of 1.0 m/s. The chair has a mass of 35 kg.



### Motion Diagram

#### Free-Body Diagrams



Object:

Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

$a_{12} =$

Event 2:

$t_2 =$

$r_2 =$

$v_2 =$

*student*



*chair*



### Mathematical Analysis

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## I-93

A man of mass  $m$ , concerned about his weight, decides to weigh himself in an elevator. He stands on a bathroom scale in an elevator moving upward at speed  $v$ . As the elevator reaches his floor, it slows to a stop over a distance,  $d$ . Determine the reading on the bathroom scale ( $F_{\text{scale}}$ ) as a function of  $m$ ,  $v$ ,  $d$ , and  $g$ .

### Motion Diagram

#### Free-Body Diagram



### Motion Information

Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

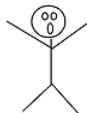
$a_{12} =$

Event 2:

$t_2 =$

$r_2 =$

$v_2 =$



### Mathematical Analysis

### Questions

If  $v = 0$  m/s, what should  $F_{\text{scale}}$  equal? Does your function agree with this observation?

If  $d = \infty$ , what should  $F_{\text{scale}}$  equal? Does your function agree with this observation?

For what stopping distance,  $d$ , would the bathroom scale read 0 N? Would the scale also read 0 N for this stopping distance if the elevator was initially moving downward?

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## I-94

A falling pole-vaulter of mass  $m$  lands on a cushion at speed  $v$ . The pole-vaulter sinks a distance  $d$  into the cushion before stopping. Determine the force exerted on the pole-vaulter due to the cushion ( $F_{\text{cushion}}$ ) as a function of  $m$ ,  $v$ ,  $d$ , and  $g$ .

### Motion Diagram

#### Free-Body Diagram



### Motion Information

Event 1:

$t_1 =$

$r_1 =$

$v_1 =$

$a_{12} =$

Event 2:

$t_2 =$

$r_2 =$

$v_2 =$



### Mathematical Analysis

### Questions

If  $v = 0$  m/s, what should  $F_{\text{cushion}}$  equal? Does your function agree with this observation?

If  $d = 0$  m, what should  $F_{\text{cushion}}$  equal? Does your function agree with this observation?

What would be worse for the pole-vaulter, hitting the cushion at twice her original speed or sinking half of the original distance into the cushion?

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## I-95

A rocket of mass  $m$  is fired vertically upward from rest. The rocket's engine produces a thrust of constant magnitude  $F_{\text{thrust}}$  for  $t_{\text{thrust}}$  seconds. Determine the maximum height reached by the rocket ( $H$ ) as a function of  $F_{\text{thrust}}$ ,  $t_{\text{thrust}}$ ,  $m$ , and  $g$ .

### Motion Information Body Diagrams

### Free-

Event 1:                  Event 2:                  Event 3:                  *before engine turns off*      *after engine turns off*

$t_1 =$                            $t_2 =$                            $t_3 =$

$r_1 =$                            $r_2 =$                            $r_3 =$

$v_1 =$                            $v_2 =$                            $v_3 =$

$a_{12} =$                            $a_{23} =$



### Mathematical Analysis

### Questions

If  $g = 0 \text{ m/s}^2$ , what should  $H$  equal? Does your function agree with this observation?

If  $F_{\text{thrust}} = mg$ , what should  $H$  equal? Does your function agree with this observation?

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## I-96

A rocket of mass  $m$  is fired vertically upward from rest. The rocket's engine produces a thrust of constant magnitude  $F_{\text{thrust}}$  for  $t_{\text{thrust}}$  seconds. Determine the time it takes the rocket to reach its apex ( $t_{\text{apex}}$ ) as a function of  $F_{\text{thrust}}$ ,  $t_{\text{thrust}}$ ,  $m$ , and  $g$ .

### Motion Information Body Diagrams

### Free-

Event 1:                  Event 2:                  Event 3:                  *before engine turns off*      *after engine turns off*

$t_1 =$                            $t_2 =$                            $t_3 =$

$r_1 =$                            $r_2 =$                            $r_3 =$

$v_1 =$                            $v_2 =$                            $v_3 =$

$a_{12} =$                            $a_{23} =$



### Mathematical Analysis

### Questions

If  $g = 0 \text{ m/s}^2$ , what should  $t_{\text{apex}}$  equal? Does your function agree with this observation?

If  $F_{\text{thrust}} = mg$ , what should  $t_{\text{apex}}$  equal? Does your function agree with this observation?

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## 05. Conservation

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- 01. Concepts and Principles
- 02. Analysis Tools
- [03. Analysis Tools 2](#)
- [04. Analysis Tools 3](#)
- [05. Activities](#)

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## 01. Concepts and Principles

1. What is the Conservation Law?
2. The Impulse-Momentum Relation
3. The Work-Energy Relation

### What is the Conservation Law?

In general, a conservation law is a statement that a certain quantity does not change over time. If you know how much of this quantity you have today, you can be assured that the exact same amount of the quantity will be available tomorrow. A famous (at least to physicists) explanation of the nature of a conservation law was given by Richard Feynman.

*Imagine your child has a set of 20 wooden blocks. Every day before bedtime, you gather up your child's blocks to put them away. As you gather up the blocks, you keep count in your head. Once you reach 20, you know you have found all of the blocks and it is unnecessary for you to search any longer. This is because the number of blocks is conserved. It is the same today as it was yesterday.*

*If one day you only find 18 blocks, you know to keep looking until you find the missing 2 blocks. Also, with experience, you discover the typical hiding places for the blocks. You know to check under the sofa, or behind the curtains.*

*If your child is rambunctious, you may even have to look outside of the room. Perhaps he threw a block or two out of the window. Even though blocks can disappear from inside of the room, and appear out in the yard, if you search everywhere you will always find the 20 blocks.*

Physicists have discovered a number of quantities that behave exactly like the number of wooden blocks. We will examine two of these quantities, energy and momentum, below.

### The Impulse-Momentum Relation

While Newton's Second Law directly relates the total force that acts on an object at a specific time to the object's acceleration at that exact same time, conservation laws relate the amount of a certain quantity present at one time to the amount present at a later time.

The first conserved quantity we will investigate is *momentum*. Of course, just because momentum is conserved doesn't mean that the momentum of any particular object is always constant. The momentum of a single object, like the number of blocks in the playroom, can change. Just as blocks can be thrown out of the window of the playroom, the momentum of a single object can be changed by applying *impulse* to it. The relationship between impulse and momentum is, conceptually,

$$\text{initial momentum} + \text{impulse} = \text{final momentum}$$

Momentum:  $P = mv$  [kg m/s]

Impulse:  $J = Ft$  [N sec]

Mathematically this is written as:

$$mv_i + \sum(F(\Delta t)) = mv_f \quad (1)$$

The product of mass and velocity is the momentum of the object, typically denoted  $P$ , and the product of force and the time interval over which the force acts is impulse,  $J$ . The summation symbol,  $\sum$ , indicates that you must sum all of the impulses acting on the object to find the change in momentum.

In short, if no impulse is applied to an object, its momentum will remain constant. However, if an impulse is applied to the object, its momentum will change by an amount exactly equal to the impulse applied. This momentum does not disappear; however, it is simply transferred to the object *supplying* the impulse. In this sense, impulse is the transfer of momentum from one object to another, analogous to tossing blocks out of the playroom and into the yard.

### The Work-Energy Relation

The second conserved quantity we will investigate is *energy*. Energy can be defined as the ability to do work and has units of Joules (J) where a  $1 \text{ J} = 1 \text{ N m}$ . Just like momentum, or wooden blocks, the conservation of energy doesn't mean that the energy of any

particular object is always constant. The energy of a single object can be changed by doing *work* to it. The relationship between work and energy is, conceptually,

$$\text{initial energy} + \text{work} = \text{final energy}$$

The similarity between momentum and energy is not complete, however. While there is only one form of momentum (i.e., one hiding place for momentum “blocks”) there are several forms of energy. These different forms of energy will be introduced as you progress through more and more complicated models of the physical world. For now, the only “hiding place” I want to discuss is *kinetic energy*. In terms of kinetic energy, the above conceptual relationship between work and energy becomes, expressed mathematically,

$$\frac{1}{2}mv_i^2 + \sum (|F| |\Delta r| \cos\theta) = \frac{1}{2}mv_f^2 \quad (2)$$

Unlike anything we’ve studied up to this point, the work-energy relation is a *scalar* equation. This will become especially important when we study objects moving in more than one dimension. For now, all this means is that in the expression for work,  $|F| |\Delta r| \cos\theta$ , we should use the *magnitude* of the force and the *magnitude* of the change in position. This product is then multiplied by  $\cos\theta$  where  $\theta$  is defined to be the angle between the applied force and the displacement of the object. If the force and displacement are in the same direction  $\theta = 0^\circ$  and the work is positive (the object gains energy). If the force and displacement are in the opposite direction  $\theta = 180^\circ$ , and the work is negative (the object loses energy). Note that the actual directions of the force and the displacement are unimportant, only their directions *relative to each other* affect the work.

In general, if no work is done to an object, its total (not simply kinetic) energy will remain constant. However, if work is done to the object, its total energy will change by an amount exactly equal to the work done. In this sense, work is the transfer of energy from one object to another, again analogous to tossing blocks out of the playroom and into the yard.

---

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## 02. Analysis Tools

1. Applying the Impulse-Momentum Relation to a Single Object
2. Applying the Work-Energy Relation to a Single Object

### Applying the Impulse-Momentum Relation to a Single Object

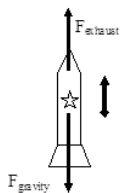
Let's investigate the following scenario:

A 0.35 kg model rocket is fitted with an engine that produces a thrust of 11.8 N for 1.8 s. The rocket is launched vertically upward.

To apply the impulse-momentum relation, you must clearly specify the initial and final events at which you will tabulate the momentum. Let's choose:

Event 1: The instant the engine is ignited.

Event 2: The instant the engine shuts off.



Applying impulse-momentum to the rocket during this time interval yields:

$$mv_i + \sum F(\Delta t) = mv_f$$

$$.35(0) + F_{exhaust}(1.8) - F_{gravity}(1.8) = .35v_2$$

$$0 + (11.8)(1.8) - (.35)(9.8)(1.8) = .35v_2$$

$$v_2 = 43.0 \text{ m/s}$$

where each applied force is independently multiplied by the time interval over which it acts. (Note that the rocket's engine does not produce a force on the rocket! The engine produces a downward force on the hot exhaust gases emitted from the engine and these hot gases exert an equal magnitude force back up on the rocket. That is why the force on the rocket is labeled as  $F_{exhaust}$  rather than  $F_{engine}$ .) Thus, the rocket is traveling at 43.0 m/s at the instant the engine shuts off.

Of course, there is no reason why we had to analyze the rocket's motion between the two instants of time we selected above. We could have selected the events:

Event 1: The instant the engine is ignited.

Event 2: The instant the rocket reaches maximum height.

During this time interval, the force of the exhaust gases only act on the rocket for a *portion* of the entire time interval. Noting that the rocket's velocity when it reaches its maximum height is zero, impulse-momentum would look like this:

$$mv_i + \sum F(\Delta t) = mv_f$$

$$.35(0) + F_{exhaust}(1.8) - F_{gravity}(\Delta t) = (0.35)(0)$$

$$0 + (11.8)(1.8) - (.35)(9.8)(\Delta t) = 0$$

$$\Delta t = 6.19 \text{ s}$$

Thus, the rocket is in the air for 6.19 s before reaching its maximum height.

### Applying the Work-Energy Relation to a Single Object

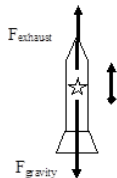
The work-energy relation also has many uses for investigating physical scenarios. For example, let's look again at our model rocket:

A 0.35 kg model rocket is fitted with an engine that produces a thrust of 11.8 N for 1.8 s. The rocket is launched vertically upward.

To apply the work-energy relation, you must clearly specify the initial and final events at which you will tabulate the energy. Let's choose:

Event 1: The instant the engine is ignited.

Event 2: The instant the engine shuts off.



Assuming we've already analyzed this scenario using impulse-momentum, what additional information can we extract using work-energy?

$$\begin{aligned}\frac{1}{2}mv_i^2 + \Sigma |F| |\Delta r| \cos \phi &= \frac{1}{2}mv_f^2 \\ \frac{1}{2}(0.35)(0)^2 + |F_{\text{exhaust}}| |\Delta r| \cos 0 + |F_{\text{gravity}}| |\Delta r| \cos 180 &= \frac{1}{2}(0.35)(43.2)^2 \\ 0 + 11.8(\Delta r)(1) + (0.35)(9.8)(\Delta r)(-1) &= 327 \\ 11.8\Delta r - 3.43\Delta r &= 327 \\ \Delta r &= 39.1\text{m}\end{aligned}$$

Thus, the rocket rises to a height of 39.1 m before the engines shuts off.

What if we apply work-energy between the following two events:

Event 1: The instant the engine is ignited.

Event 2: The instant the rocket reaches its maximum height.

During this time interval, the force of the exhaust gases only act on the rocket for a *portion* of the entire displacement, namely 39 m, while the force of gravity acts over the entire displacement.

$$\begin{aligned}\frac{1}{2}mv_i^2 + \Sigma |F| |\Delta r| \cos \phi &= \frac{1}{2}mv_f^2 \\ \frac{1}{2}(0.35)(0)^2 + (11.8)(39) \cos 0 + (0.35)(9.8)(\Delta r) \cos 180 &= \frac{1}{2}(0.35)(0)^2 \\ 0 + 460 - 3.43\Delta r &= 0 \\ \Delta r &= 134\text{m}\end{aligned}$$

Thus, the maximum height reached by the rocket is 134 m.

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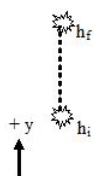
## 03. Analysis Tools 2

1. Gravitational Potential Energy
2. Applying Work-Energy with Gravitational Potential Energy

### Gravitational Potential Energy

In any situation in which an object changes its height above the surface of the earth, the force of gravity does work on the object. It is possible to calculate this work in general, and to rewrite the work-energy relation in such a way as to incorporate the effects of this work. This is referred to as constructing a *potential energy function* for the work done by gravity.

Let's imagine an object of mass,  $m$ , located an initial height,  $h_i$ , above the zero of a vertical coordinate system, with the upward direction designated positive. It moves to a final height of  $h_f$ .



To calculate the work done by gravity on this object:

$$\begin{aligned} W_{gravity} &= |F||\Delta r|\cos\phi \\ W_{gravity} &= (mg)(h_f - h_i)\cos 180 \\ W_{gravity} &= -mgh_f + mgh_i \end{aligned}$$

The “ $mgh$ ” terms are referred to as *gravitational potential energy*. Thus, the work done by gravity can be thought of as changing the gravitational potential energy of the object. Let's plug the above result into the work-energy relation:

$$\begin{aligned} \frac{1}{2}mv_i^2 + \Sigma|F||\Delta r|\cos\phi &= \frac{1}{2}mv_f^2 \\ \frac{1}{2}mv_i^2 + \Sigma|F||\Delta r|\cos\phi + W_{gravity} &= \frac{1}{2}mv_f^2 \\ \frac{1}{2}mv_i^2 + \Sigma|F||\Delta r|\cos\phi - mgh_f + mgh_i &= \frac{1}{2}mv_f^2 \\ \frac{1}{2}mv_i^2 + mgh_i + \Sigma|F||\Delta r|\cos\phi &= \frac{1}{2}mv_f^2 + mgh_f \end{aligned}$$

Therefore, this final relation:

$$\text{initial energy} + \text{work} = \text{final energy}$$

$$\frac{1}{2}mv_i^2 + mgh_i + \Sigma|F||\Delta r|\cos\phi = \frac{1}{2}mv_f^2 + mgh_f$$

can be used in place of the standard work-energy relation provided:

1. You do not include the force of gravity a second time by calculating the work done by gravity. Basically, in this relationship gravity is no longer thought of as a force that does work on objects but rather as a source of potential energy.
2. You calculate the initial and final heights,  $h_i$  and  $h_f$ , using a coordinate system in which the upward direction is positive.

### Applying Work-Energy with Gravitational Potential Energy

Let's use the work-energy relation, with gravitational potential energy terms, to re-analyze the previous scenario:

A 0.35 kg model rocket is fitted with an engine that produces a thrust of 11.8 N for 1.8 s. The rocket is launched vertically upward.

We apply work-energy between the following two events:

Event 1: The instant the engine is ignited.

Event 2: The instant the rocket reaches its maximum height.

During this time interval, the force of the exhaust gases only act on the rocket for a *portion* of the entire displacement, namely 39 m. Remember, the force of gravity *does not do work* in this way of modeling nature, rather the gravitational energy of the rocket changes as it changes its elevation.

Setting the initial elevation of the rocket equal to zero:

$$\begin{aligned}\frac{1}{2}mv_i^2 + mgh_i + \Sigma|F||\Delta r|\cos\phi &= \frac{1}{2}mv_f^2 + mgh_f \\ \frac{1}{2}(0.35)(0)^2 + (0.35)(9.8)(0) + (11.8)(39)\cos 0 &= \frac{1}{2}(0.35)(0)^2 + (0.35)(9.8)h_f \\ 0 + 0 + 460 &= 0 + 3.43h_f \\ h_f &= 134m\end{aligned}$$

results in, of course, the same maximum height reached by the rocket.

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## 04. Analysis Tools 3

1. Applying the Impulse-Momentum Relation to a Collision
2. Applying the Work-Energy Relation to the Same Collision

### Applying the Impulse-Momentum Relation to a Collision

Probably the most useful application of the impulse-momentum relation is in the study of collisions. For example:

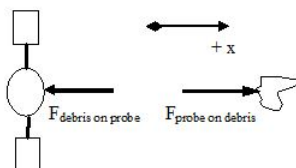
Far from the earth, a 250 kg space probe, moving at 5 km/s, collides head-on with a 65 kg piece of space debris initially at rest. The debris becomes entangled in the probe's solar collectors.

Let's choose:

**Event 1: The instant before the collision.**

**Event 2: The instant the debris and probe reach a common velocity.**

The free-body diagrams for the two objects during this time interval are shown below.



Applying the impulse-momentum relation to each object separately yields:

<i>Probe</i>	<i>Debris</i>
$mv_i + \Sigma F(\Delta t) = mv_f$	$mv_i + \Sigma F(\Delta t) = mv_f$
$250(5000) - F_{\text{debris on probe}}(\Delta t) = 250v_2$	$60(0) + F_{\text{probe on debris}}(\Delta t) = 60v_2$
$1250000 - F_{\text{debris on probe}}(\Delta t) = 250v_2$	$F_{\text{probe on debris}}(\Delta t) = 60v_2$

Notice that the final velocities of the two objects are the same, because they remain joined together following the collision. Also, the  $\Delta t$ 's are the same because the time interval over which the force of the debris acts on the probe must be the same as the time interval over which the force of the probe acts on the debris. In fact, these two forces must be equal to each other in magnitude by Newton's Third Law.

Thus, the impulses must cancel if the two equations are added together:

$$\begin{aligned}
 1250000 - F_{\text{debris on probe}}(\Delta t) &= 250v_2 \\
 \underline{F_{\text{probe on debris}}(\Delta t) &= 60v_2} \\
 1250000 &= 310v_2 \\
 v_2 &= 4032 \text{ m/s}
 \end{aligned}$$

The probe slows to a speed of 4032 m/s (and the debris changes direction and accelerates to a speed of 4032 m/s) via the collision. Thus, even though we do not know the magnitude of the force involved, or the duration of the collision, we can calculate the final velocities of the two objects colliding. This is because the forces involved comprise an interaction, and by Newton's Third Law forces that comprise an interaction are always equal in magnitude and opposite in direction.

In fact, in problems involving collisions (or explosions, which to physicists are merely collisions played backward in time!), you should almost always apply the impulse-momentum relation to the interacting objects because the forces involved comprise an interaction. Thus, by adding your equations together, these terms will always add to zero. This will often allow you to determine the final velocities of the colliding objects.

In conclusion, I should point out that the probe loses momentum during the collision and that the debris gains the exact same amount of momentum (check the numbers to verify this statement). The momentum is transferred from the probe to the debris through the action of the impulse the probe and debris exert on each other. The momentum transfer from the probe to the debris is analogous to throwing a wooden block from the playroom into the yard: The playroom now has one less block and the yard has one more!

## Applying the Work-Energy Relation to the Same Collision

Let's return to the collision scenario discussed above and attempt to investigate it using work-energy.

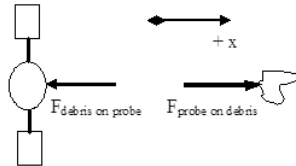
Far from the earth, a 250 kg space probe, moving at 5 km/s, collides head-on with a 65 kg piece of space debris initially at rest. The debris becomes entangled in the probe's solar collectors.

Let's choose:

**Event 1: The instant before the collision.**

**Event 2: The instant the debris and probe reach a common velocity.**

The free-body diagrams for the two objects during this time interval are shown below.



Applying the work-energy relation to each object separately yields:

Probe	Debris
$v_i^2 + \sum  F  \Delta r \cos \phi = \frac{1}{2} m v_f^2$	$\frac{1}{2} m v_i^2 + \sum  F  \Delta r \cos \phi = \frac{1}{2} m v_f^2$
$50(5000)^2 + (F_{\text{debris}})(\Delta r_{\text{probe}}) \cos 180 = \frac{1}{2} (250)(v_{2\text{probe}})^2$	$\frac{1}{2} (60)(0)^2 + (F_{\text{probe}})(\Delta r_{\text{debris}}) \cos 0 = \frac{1}{2} (60)(v_{2\text{debris}})^2$
$x10^9 - F_{\text{debris}} (\Delta r_{\text{probe}}) = 125 v_{2\text{probe}}^2$	$F_{\text{probe}} (\Delta r_{\text{debris}}) = 30 v_{2\text{debris}}^2$

The final velocities of the two objects are the same, because they remain joined together following the collision, and the two forces are the same by Newton's Third Law. **However, these two equations can not be added together and solved because the two distances over which the forces act,  $\Delta r_{\text{probe}}$  and  $\Delta r_{\text{debris}}$ , are not necessarily equal.** During the collision, the center of the probe will move a different distance than the center of the debris<sup>[1]</sup>. Since these two distances are different, the works will *not* cancel as the impulses did, and the equations are *not* solvable!

In fact, since we know  $v_2 = 4032$  m/s from our momentum analysis,

Probe	Debris
$3.13 \times 10^9 - F_{\text{debris}} (\Delta r_{\text{probe}}) = 125(4032)^2$	$F_{\text{probe}} (\Delta r_{\text{debris}}) = 30(4032)^2$
$3.13 \times 10^9 - F_{\text{debris}} (\Delta r_{\text{probe}}) = 2.03 \times 10^9$	$W_{\text{on debris}} = F_{\text{probe}} (\Delta r_{\text{debris}}) = 0.49 \times 10^9 \text{ J}$
$W_{\text{on probe}} = -F_{\text{debris}} (\Delta r_{\text{probe}}) = -1.1 \times 10^9 \text{ J}$	

Obviously, the two works do not cancel. In fact, the *internal work*, or work done by the objects on each other, totals  $-0.61 \times 10^9$  J. This means that there is  $0.61 \times 10^9$  J *less* kinetic energy in the system of the probe and the debris after the collision than before the collision. This is sometimes referred to as the energy lost in the collision, although the energy is not lost but rather converted into other forms of energy (i.e., other hiding places for the wooden blocks that have yet to be discussed), such as thermal energy.

In short, the work-energy relation (as it now stands) can not be used to effectively analyze collisions unless additional information regarding the internal energy is available. Occasionally, an approximation is made in which the total internal work is zero. When this approximation is made, the collision is referred to as an *elastic* collision. Realistic collisions, in which the total internal energy is not zero and kinetic energy is "lost", are referred to as *inelastic* collisions. In other words, during an elastic collision the two objects "bounce apart" whereas during an inelastic collision they "stick together."

[1] If the two objects were *actually* particles, rather than being *approximated* as particles, then the two distances would have to be the same and the two works would cancel when the equation were added together.

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## 05. Activities

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## 01. Selected Answers For Conservation

### Selected Answers for ModelSpiral Physics: Mechanics

#### Kinematics

- 41.  $r_2 = 1.4 \text{ m}$
- 42.  $t_2 = 1.45 \text{ s}$
- 43.  $t_3 = 3.4 \text{ s}$
- 44.  $t_3 = 4.55 \text{ s}$
- 45.  $r_3 = 36.6 \text{ m}$
- 46.  $r_4 = 2.0 \text{ km}$
- 47.  $t_4 = 370 \text{ ks}$
- 49.  $t_2 = 14.9 \text{ s}$
- 50.  $t_2 = 7.8 \text{ s}$
- 51.  $t_3 = 2.87 \text{ s}$

#### Dynamics

- 80.  $F_{\text{rope}} = 420 \text{ N}$
- 81.  $F_{\text{cushion}} = 1.91 \text{ kN}$
- 83.  $F_{\text{bottom cable}} = 11.4 \text{ N}$
- 84.  $F_{\text{cushion}} = 2.83 \text{ kN}$
- 85.  $F_{\text{ground}} = 43.7 \text{ kN}$
- 86.  $r_3 = 63.5 \text{ m}$
- 87.  $r_3 = 1.55 \text{ km}$
- 89.  $m_{\text{block}} = 240 \text{ kg}$
- 90.  $m_{\text{block}} = 26 \text{ kg}$
- 91.  $F_{\text{rope}} = 500 \text{ N}$
- 92.  $F_{\text{rope}} = 490 \text{ N}$

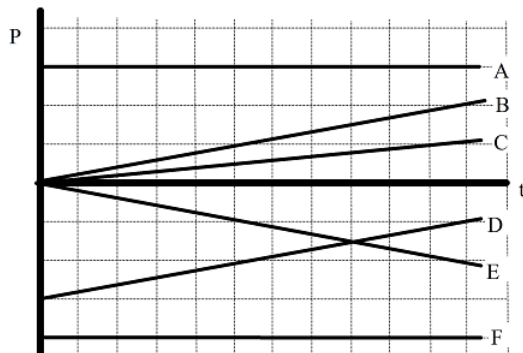
#### Conservation Laws

- 122. a.  $F_{\text{scale}} = 755 \text{ N}$       b.  $F_{\text{scale}} = 780 \text{ N}$
- 123. a.  $v = 116 \text{ m/s}$       b.  $t = 26.8 \text{ s}$
- 124. a.  $v = 12.5 \text{ m/s}$       b.  $F_{\text{ground}} = 43.7 \text{ kN}$
- 125.  $t_2 = 0.89 \text{ s}$        $r_2 = 1.78 \text{ m}$
- 127.  $m_{\text{student}} = 94 \text{ kg}$
- 128.  $v_2 = 6.0 \text{ m/s}$
- 129.  $v_2 = 1.5 \text{ m/s}$
- 130.  $v_{2 \text{ ship}} = 24.3 \text{ m/s}$
- 131.  $v_{\text{probe}} = 667 \text{ m/s}$
- 132.  $v_{2 \text{ platform}} = 0.28 \text{ m/s}$
- 133.  $v_{2 \text{ platform}} = 0.06 \text{ m/s}$
- 134.  $v_{2 \text{ balloon}} = 6.3 \text{ m/s}$

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## I-109

Below are momentum vs. time graphs for six different objects.



a. Rank these graphs on the basis of the change in momentum of the object over the time interval indicated. Rank positive changes in momentum as larger than negative changes in momentum.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. Rank these graphs on the basis of the total impulse on the object over the time interval indicated. Rank a positive total impulse as larger than a negative total impulse.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

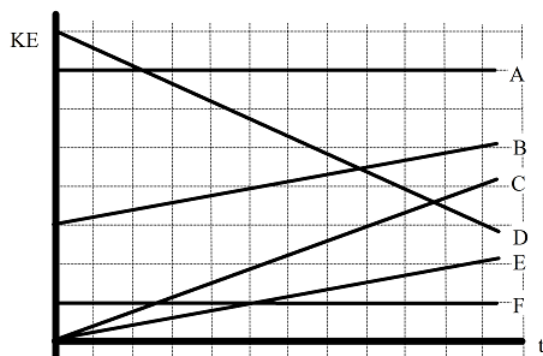
Explain the reason for your ranking:

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## I-110

Below are kinetic energy vs. time graphs for six different objects. All of the objects move horizontally.



a. Rank these graphs on the basis of the change in kinetic energy of the object over the time interval indicated. Rank positive changes in kinetic energy as larger than negative changes in kinetic energy.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. Rank these graphs on the basis of the total work on the object over the time interval indicated. Rank a positive total work as larger than a negative total work.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

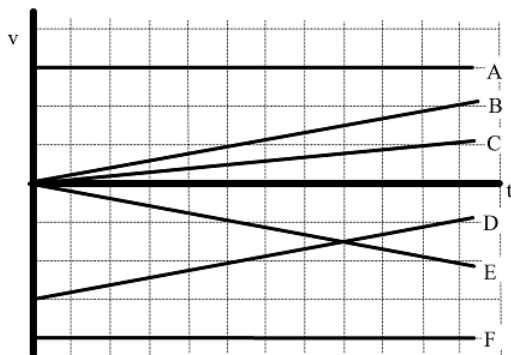
\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

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## I-111

Below are velocity vs. time graphs for six equal-mass objects.



a. Rank these graphs on the basis of the change in momentum of the object over the time interval indicated. Rank positive changes in momentum as larger than negative changes in momentum.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest  
 \_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. Rank these graphs on the basis of the change in kinetic energy of the object over the time interval indicated. Rank positive changes in kinetic energy as larger than negative changes in kinetic energy.

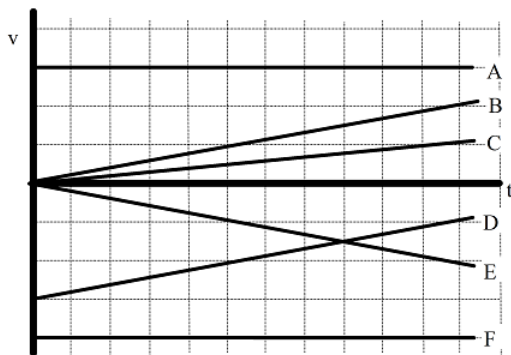
Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest  
 \_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

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## I-112

Below are velocity vs. time graphs for six equal-mass objects. All of the objects move horizontally.



a. Rank these graphs on the basis of the total impulse on the object over the time interval indicated. Rank a positive total impulse as larger than a negative total impulse.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. Rank these graphs on the basis of the total work on the object over the time interval indicated. Rank a positive total work as larger than a negative total work.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

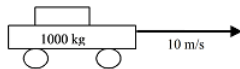
Explain the reason for your ranking:

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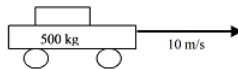
## I-113

Below are six automobiles initially traveling at the indicated velocity. The automobiles have different masses and velocities.

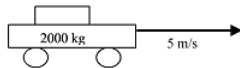
**A**



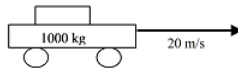
**B**



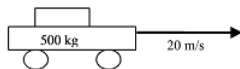
**C**



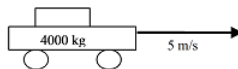
**D**



**E**



**F**



a. All automobiles will be stopped in the same amount of time. Rank these automobiles on the basis of the magnitude of the force needed to stop them.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. All automobiles will be stopped in the same amount of distance. Rank these automobiles on the basis of the magnitude of the force needed to stop them.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

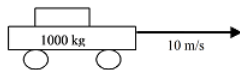
Explain the reason for your ranking:

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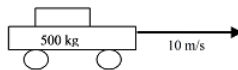
## I-114

Below are six automobiles initially traveling at the indicated velocity. The automobiles have different masses and velocities.

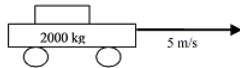
**A**



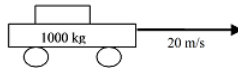
**B**



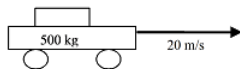
**C**



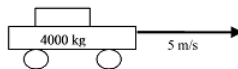
**D**



**E**



**F**



a. Rank these automobiles on the basis of the magnitude of the force needed to stop them.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

b. Rank these automobiles on the basis of the magnitude of the work needed to stop them.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

Explain the reason for your ranking:

c. Rank these automobiles on the basis of the magnitude of the impulse needed to stop them.

Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest

\_\_\_\_ The ranking can not be determined based on the information provided.

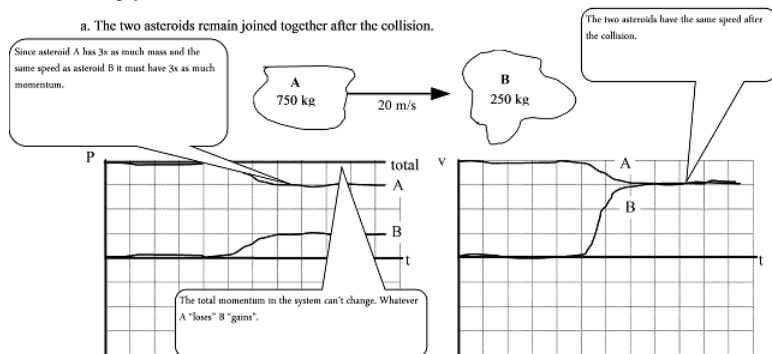
Explain the reason for your ranking:

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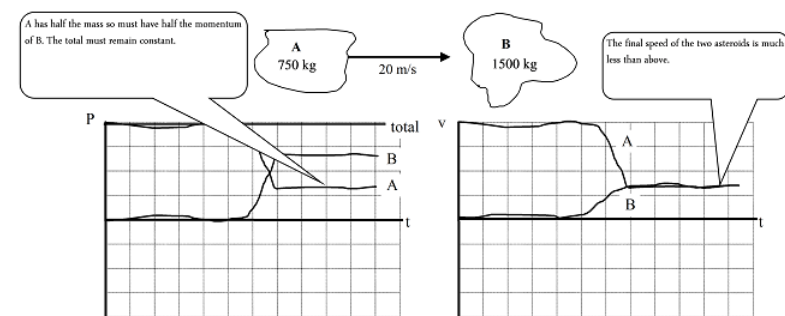
## I-115

For each of the collisions illustrated below, sketch a graph of the momentum and velocity of asteroid A, the momentum and velocity of asteroid B, and the total momentum in the system of the two asteroids. Begin your graph before the collision takes place and continue it after the collision is over. Use a consistent scale on all graphs.

a. The two asteroids remain joined together after the collision.



b. The two asteroids remain joined together after the collision.

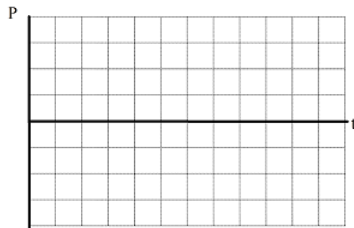
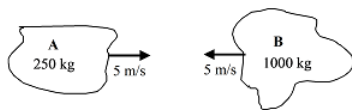


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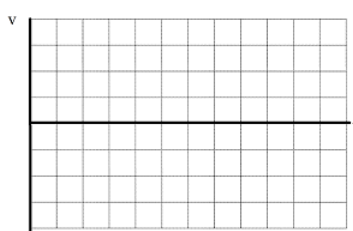
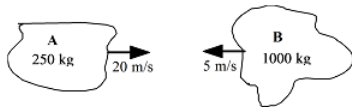
## I-116

For each of the collisions illustrated below, sketch a graph of the momentum and velocity of asteroid A, the momentum and velocity of asteroid B, and the total momentum in the system of the two asteroids. Begin your graph before the collision takes place and continue it after the collision is over. Use a consistent scale on all graphs.

a. The two asteroids remain joined together after the collision.



b. The two asteroids remain joined together after the collision.

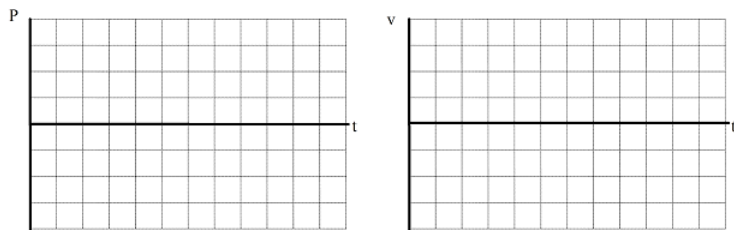
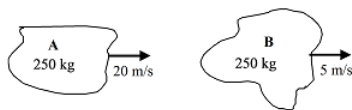


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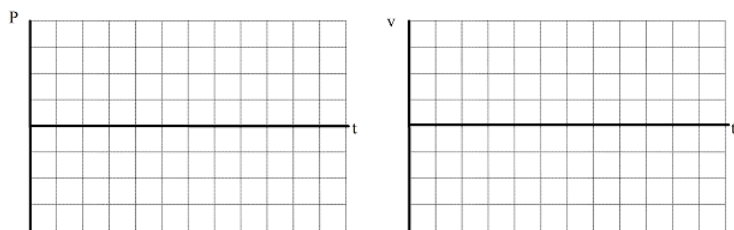
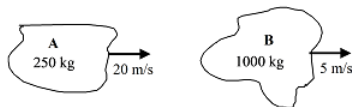
## I-117

For each of the collisions illustrated below, sketch a graph of the momentum and velocity of asteroid A, the momentum and velocity of asteroid B, and the total momentum in the system of the two asteroids. Begin your graph before the collision takes place and continue it after the collision is over. Use a consistent scale on all graphs.

a. The two asteroids remain joined together after the collision.



b. The two asteroids remain joined together after the collision.



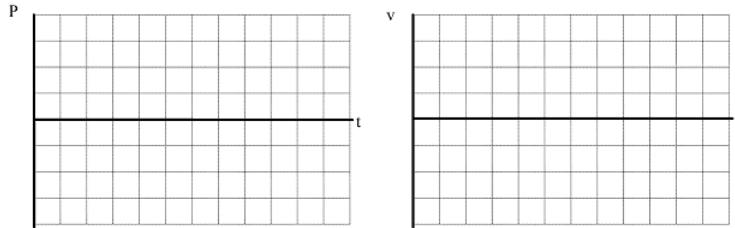
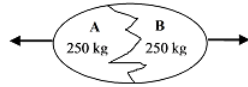
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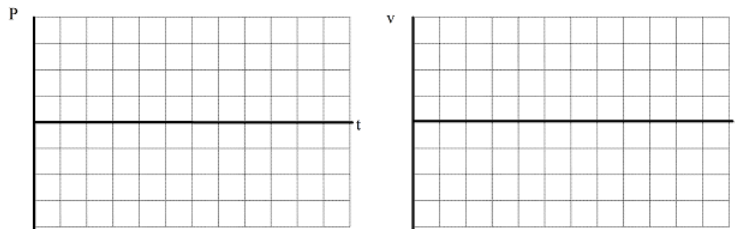
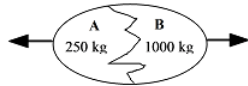
## I-118

For each of the explosions illustrated below, sketch a graph of the momentum and velocity of fragment A, the momentum and velocity of fragment B, and the total momentum in the system of the two fragments. Begin your graph before the explosion takes place and continue it while the fragments travel away from the sight of the explosion. Use a consistent coordinate system and scale on all graphs.

a. The exploding egg is initially at rest.



b. The exploding egg is initially at rest.

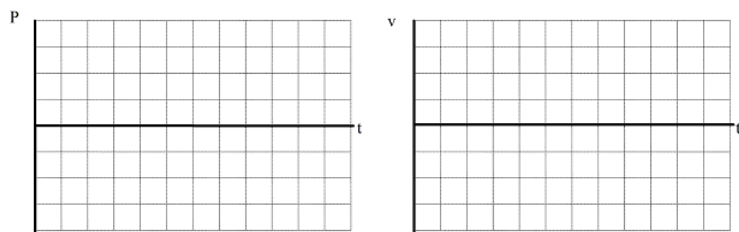


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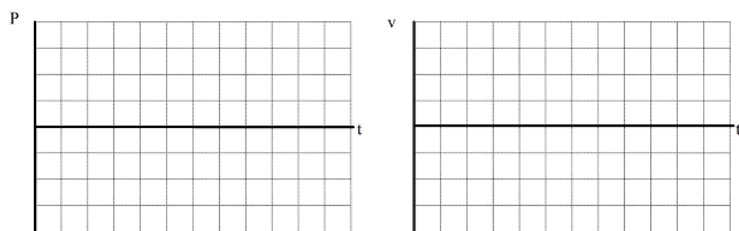
## I-119

A 200 kg astronaut is initially at rest on the extreme edge of a 1000 kg space platform. She wears special magnetic shoes that allow her to walk along the metal platform. For each of the situations illustrated below, sketch a graph of the momentum and velocity of the astronaut, the momentum and velocity of the platform, and the total momentum in the system of the two objects. Begin your graph before the astronaut begins to walk and continue it while she walks along the platform. Use a consistent coordinate system and scale on all graphs.

- a. The astronaut and platform are initially at rest.



- b. The astronaut and platform are initially drifting to the right.

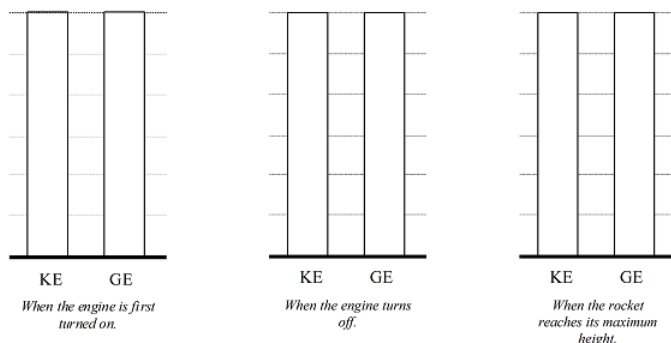


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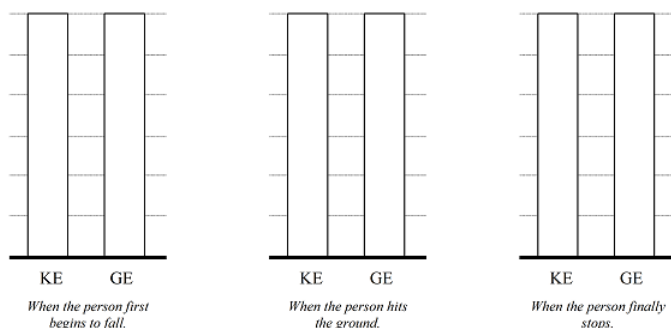
## I-120

For each of the scenarios described below, indicate the amount of kinetic energy and gravitational potential energy of the object at each of the events listed. Use a consistent scale throughout each motion. Set the lowest point of the motion as the zero-point of gravitational potential energy

a. A 4000 kg rocket's engine produces a thrust of 70,000 N for 15 s. The rocket is fired vertically upward.



b. To practice falling, a 55 kg pole-vaulter falls off of a wall 6.0 m above a 2.0 m thick foam cushion resting on the ground. However, he misses the cushion. The pole-vaulter sinks about 0.10 m into the ground before stopping.

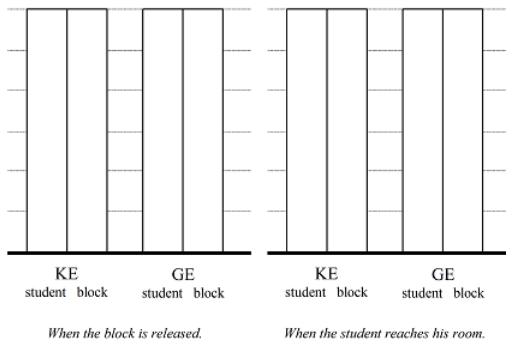


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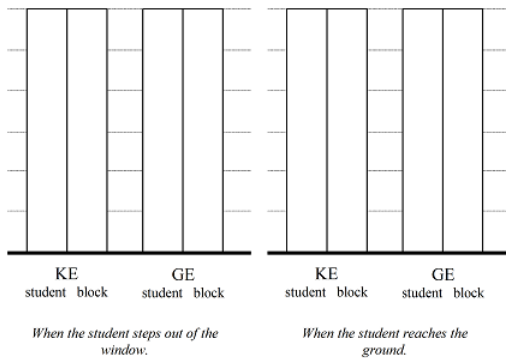
## I-121

For each of the scenarios described below, indicate the amount of kinetic energy and gravitational potential energy of each object at each of the events listed. Use a consistent scale throughout each motion. Set the lowest point of the motion as the zero-point of gravitational potential energy

a. Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching his third floor dorm room. An 84 kg block is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the 84 kg block is released, the student is pulled up to his dorm room, 8.0 m off the ground.



b. Tired of walking down the stairs, a 75 kg engineering student designs an ingenious device for reaching the ground from his third floor dorm room. A 60 kg block at rest on the ground is attached to a rope that passes over a pulley. The student grabs the other end of the rope and steps out of his window.



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## I-122

A 100 kg man concerned about his weight decides to weigh himself in an elevator. He stands on a bathroom scale in an elevator that is moving upward at 3.0 m/s. As the elevator reaches his floor, it slows to a stop.

a. If the elevator slows to a stop over a distance of 2.0 m, what is the reading on the bathroom scale?

**Free-Body Diagram**



**Mathematical Analysis**

Event 1:

Event 2:

b. If the elevator slows to a stop in 1.5 s, what is the reading on the bathroom scale?

**Free-Body Diagram**



**Mathematical Analysis**

Event 1:

Event 2:

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## I-123

*A 4000 kg rocket's engine produces a thrust of 70,000 N for 15 s. The rocket is fired vertically upward.*

a. What is the speed of the rocket when its engine turns off?

**Free-Body Diagram**



**Mathematical Analysis**

Event 1:

Event 2:

b. How long does it take the rocket to reach its maximum height?

**Free-Body Diagram**



**Mathematical Analysis**

Event 1:

Event 2:

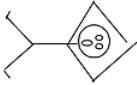
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## I-124

*To practice falling, a 55 kg pole-vaulter falls off of a wall 6.0 m above a 2.0 m thick foam cushion resting on the ground. However, he misses the cushion. The pole-vaulter sinks about 0.10 m into the ground before stopping.*

- a. What is the speed of the pole-vaulter when he hits the ground?

**Free-Body Diagram**



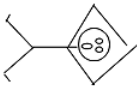
**Mathematical Analysis**

Event 1:

Event 2:

- b. What is the force exerted on the pole-vaulter by the ground as he comes to rest?

**Free-Body Diagram**



**Mathematical Analysis**

Event 1:

Event 2:

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## I-125

*A decorative light fixture in an elevator consists of a 2.0 kg light suspended by a cable from the ceiling of the elevator. From this light, a separate cable suspends a second 0.80 kg light. The elevator is moving downward at 4.0 m/s when someone presses the emergency stop button. During the stop, the upper cable snaps. The elevator engineer says that the cable could withstand a force of 40 N without breaking.*

### Free-Body Diagrams

#### Mathematical Analysis

*the top light*

*the bottom light*

Event 1:

Event 2:



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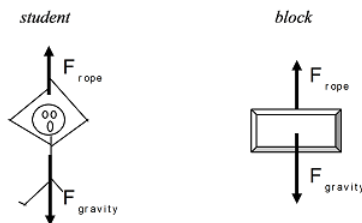


## I-126

Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching his third floor dorm room. An 84 kg block is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the 84 kg block is released, the student is pulled up to his dorm room, 8.0 m off the ground.

### Free-Body Diagrams

#### Mathematical Analysis



Event 1: The instant the block is released.  
Event 2: The instant the student reaches his room.

Since the distance the student and block travel is known, applying work-energy should allow us to solve the problem. I'll apply it separately to each object.

$$\begin{array}{lcl}
 \text{student} & & \text{block} \\
 \frac{1}{2}mv_i^2 + mgh_i + \Sigma|F||\Delta r|\cos\phi = \frac{1}{2}mv_f^2 + mgh_f & & \frac{1}{2}mv_i^2 + mgh_i + \Sigma|F||\Delta r|\cos\phi = \frac{1}{2}mv_f^2 + mgh_f \\
 0 + 0 + F_{\text{rope}}(8)\cos 0 = \frac{1}{2}(80)v_f^2 + 80(9.8)(8) & & 0 + 84(9.8)(8) + F_{\text{rope}}(8)\cos 180 = \frac{1}{2}(84)v_f^2 + 0 \\
 8F_{\text{rope}} = 40v_f^2 + 6272 & & 6586 - 8F_{\text{rope}} = 42v_f^2
 \end{array}$$

$F_{\text{rope}}$  is the same in both equations, as is the final velocity. Thus the two equations can be added together to yield:

$$\begin{aligned}
 6586 &= 40v_f^2 + 42v_f^2 + 6272 \\
 314 &= 82v_f^2 \\
 v_f &= 1.96 \text{ m/s}
 \end{aligned}$$

Notice that if you applied work-energy to the entire system you would have generated this same equation. Initially, the only form of energy present is the gravitational energy of the block ( $mgh = 6586 \text{ J}$ ). At the second event, both objects have kinetic energy plus the student has gravitational potential energy ( $mgh = 6272 \text{ J}$ ).

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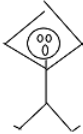
## I-127

Tired of walking up the stairs, an engineering student designs an ingenious device for reaching his third floor dorm room. A 100 kg block is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the block is released, the student is pulled up to his dorm room, 8.0 m off the ground. He is traveling at 2.2 m/s when he reaches his room.

### Free-Body Diagrams

#### Mathematical Analysis

student



block



Event 1:

Event 2:

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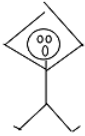
## I-128

Tired of walking down the stairs, a 75 kg engineering student designs an ingenious device for reaching the ground from his third floor dorm room. A 60 kg block at rest on the ground is attached to a rope that passes over a pulley. The student grabs the other end of the rope and steps out of his window. He falls for 5.5 s before reaching the ground.

### Free-Body Diagrams

#### Mathematical Analysis

student



block



Event 1:

Event 2:

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## I-129

*Far from any other masses, a 2000 kg asteroid traveling at 12 m/s collides with a 1200 kg asteroid traveling in the other direction at 16 m/s. After the collision they remain joined together and move with a common velocity.*

### Free-Body Diagrams

#### Mathematical Analysis

2000 kg asteroid



1200 kg asteroid



Event 1:

Event 2:

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## I-130

In the farthest reaches of deep space, an 8000 kg spaceship, including contents, is at rest relative to a space station. The spaceship recoils after it launches a 600 kg scientific probe with a speed of 300 m/s relative to the space station.

### Free-Body Diagrams

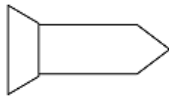
#### Mathematical Analysis

*spaceship*

*probe*

Event 1:

Event 2:



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## I-131

*In the farthest reaches of deep space, an 8000 kg spaceship, including contents, is drifting at 50 m/s relative to a space station. The spaceship is brought to rest, relative to the space station, by the recoil from launching a 600 kg scientific probe.*

### Free-Body Diagrams

#### Mathematical Analysis

*spaceship*

*probe*

Event 1:

Event 2:



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## I-132

A 140 kg astronaut is standing on the extreme edge of a 1000 kg space platform, at rest relative to the mother ship. She walks at a constant speed of 2.0 m/s, relative to the mother ship, toward the other edge of the platform. (She wears special magnetic shoes that allow her to walk along the metal platform.) As she walks, the platform moves through space.

### Free-Body Diagrams

#### Mathematical Analysis

astronaut

platform

Event 1:

Event 2:



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## I-133

Two astronauts, one 140 kg and the other 170 kg, are standing on opposite edges of a 1000 kg space platform, at rest relative to the mother ship. They walk at constant speeds of 2.0 m/s, relative to the mother ship, toward the other edge of the platform. (They wear special magnetic shoes that allow them to walk along the metal platform.) As they walk, the platform moves through space.

### Free-Body Diagrams

140 kg astronaut



platform



170 kg astronaut



### Mathematical Analysis

Event 1:

Event 2:

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## I-134

A 70 kg student is hanging from a 280 kg helium balloon. The balloon is rising at a constant speed of 8.0 m/s relative to the ground. The lift on the balloon due to the buoyant force is constant. The student begins to climb up the rope at a speed of 15 m/s relative to the ground. The balloon's upward speed is decreased as the student climbs.

### Free-Body Diagrams

#### Mathematical Analysis

student



balloon



Event 1:

Event 2:

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## I-135

A man of mass  $m$ , concerned about his weight, decides to weigh himself in an elevator. He stands on a bathroom scale in an elevator which is moving upward at  $v$ . As the elevator reaches his floor, it slows to a stop over a time interval,  $T$ . Determine the reading on the bathroom scale ( $F_{\text{scale}}$ ) as a function of  $m$ ,  $v$ ,  $T$ , and  $g$ .

### Free-Body Diagram



### Mathematical Analysis

Event 1:

Event 2:

### Questions

If  $T = \infty$ , what should  $F_{\text{scale}}$  equal? Does your function agree with this observation?

For what combination of  $v$  and  $T$  would the bathroom scale read 0 N?

If the elevator were initially going down, would the above combination of  $v$  and  $T$  also lead to a scale reading of 0 N?

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## I-136

A rocket of mass  $m$  is fired vertically upward from rest. The rocket's engine produces a thrust of constant magnitude  $F$  for  $t_{\text{thrust}}$  seconds. Determine the time it takes the rocket to reach its apex ( $t_{\text{apex}}$ ) as a function of  $F$ ,  $t_{\text{thrust}}$ ,  $m$ , and  $g$ .

### Free-Body Diagram



### Mathematical Analysis

Event 1:

Event 2:

### Questions

If  $g = 0 \text{ m/s}^2$ , what should  $t_{\text{apex}}$  equal? Does your function agree with this observation?

If  $F = mg$ , what should  $t_{\text{apex}}$  equal? Does your function agree with this observation?

For what value of  $F$  would  $t_{\text{apex}} = 2t_{\text{thrust}}$ ?

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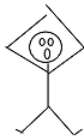
## I-137

Tired of walking up the stairs, an engineering student of mass  $m$  designs an ingenious device for reaching his third floor dorm room. A block of mass  $M$  is attached to a rope that passes over a pulley. The student holds the other end of the rope. When the block is released, the student is pulled up to his dorm room in a time  $T$ . Determine the velocity of the student ( $v$ ) when he reaches his room as a function of  $m$ ,  $M$ ,  $T$  and  $g$ .

### Free-Body Diagrams

#### Mathematical Analysis

student



block



Event 1:

Event 2:

### Questions

If  $g = 0 \text{ m/s}^2$ , what should  $v$  equal? Does your function agree with this observation?

If  $m = M$ , what should  $v$  equal? Does your function agree with this observation?

If  $M = \infty$ , what should  $v$  equal? Does your function agree with this observation?

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## I-138

In the farthest reaches of deep space, a spaceship of mass  $M$ , including contents, is at rest relative to a space station. The spaceship recoils after it launches a scientific probe of mass  $m$  at a speed  $v$  relative to the space station. Determine the recoil speed of the spaceship ( $V$ ) as a function of  $M$ ,  $m$ , and  $v$ .

### Free-Body Diagrams

#### Mathematical Analysis

spaceship

probe

Event 1:

Event 2:



### Questions

If  $M = 2m$ , what should  $V$  equal? Does your function agree with this observation?

If  $M = \infty$ , what should  $V$  equal? Does your function agree with this observation?

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## CHAPTER OVERVIEW

### Model 2: The constant-force particle model

#### Topic hierarchy

- 01. Model Specifics
- 02. Kinematics
  - 01. Concepts and Principles
  - 02. Drawing Motion Diagrams
  - 03. Drawing Motion Graphs
  - 04. Tabulating Motion Information & Doing The Math
  - 05. Hints and Suggestions
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### 03. Dynamics

01. Concepts and Principles

02. Drawing Free-Body Diagrams

03. Applying Newton's Second Law

04. Choosing a Coordinate System

05. Activities

01. Selected Answers For Dynamics

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#### 04. Conservation

01. Concepts and Principles

02. Analysis Tools

03. Analysis Tools 2

04. Activities

01. Selected Answers For Conservation

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## 01. Model Specifics

---

For our second pass through the study of mechanics, we will eliminate one of the major restrictions in our original model. We will now allow objects to move through three-dimensional space. We will still, however, restrict our model with the following approximations.

**The object is acted on by constant forces.**

**The object's size and shape are unimportant.**

**The object is classical.**

---

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## 02. Kinematics

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[01. Concepts and Principles](#)

[02. Drawing Motion Diagrams](#)

[03. Drawing Motion Graphs](#)

[04. Tabulating Motion Information & Doing The Math](#)

[05. Hints and Suggestions](#)

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## 01. Concepts and Principles

1. Position ( $r$ )
2. Velocity ( $v$ )
3. Acceleration ( $a$ )

An empirical fact about nature is that motion in one direction (for example, the horizontal) does not appear to influence aspects of the motion in a perpendicular direction (the vertical). Imagine a coin dropped from shoulder height. The elapsed time for the coin to hit the ground, the rate at which its vertical position is changing, and its vertical acceleration are the same whether you do this experiment in a stationary bus or one traveling down a smooth, level highway at 65 mph. The horizontal motion of the coin does not affect these aspects of its vertical motion.<sup>[1]</sup>

Thus, to completely describe the motion of an object moving both horizontally and vertically, you can first ignore the horizontal motion, and describe only the vertical *component* of the motion, and then ignore the vertical motion, and describe the horizontal component. Putting these kinematic components together gives you a complete description of the motion. This experimental fact about nature will make analyzing multi-dimensional motion no more conceptually difficult than analyzing one-dimensional motion.

Given this independence between motions in perpendicular directions, the same kinematic concepts and relationships utilized in one-dimensional motion will be utilized for multi-dimensional motion.

### Position ( $r$ )

The position of an object is its location relative to a well-defined coordinate system. In multi-dimensional situations, however, you must designate coordinate systems for all perpendicular directions of interest. The zero and positive direction for one dimension is completely independent of the zero and positive direction for another direction. The location at which all coordinate system zeros intersect is referred to as the origin of the coordinate system.

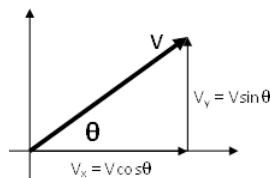
### Velocity ( $v$ )

The velocity is the rate at which the position is changing. Thus, we will define the velocity component in the vertical direction, for example, as the rate at which the vertical position is changing. The velocity component in the vertical direction is *completely* independent of the horizontal position or the rate at which the horizontal position changes.

As long as the coordinate directions are perpendicular, the speed, or magnitude of the object's velocity, can be determined by:

$$v = \sqrt{v_x^2 + v_y^2 + v_z^2}$$

The direction of the object's velocity can be determined via right-angle trigonometry.



### Acceleration ( $a$ )

The acceleration is the rate at which the velocity is changing. Thus, we will define the acceleration component in the vertical direction, for example, as the rate at which the velocity component in the vertical direction is changing. The acceleration component in the vertical direction is *completely* independent of the velocity component in the horizontal direction or the rate at which the velocity component in the horizontal direction changes.

As long as the coordinate directions are perpendicular, the magnitude of the object's acceleration can always be determined by:

$$a = \sqrt{a_x^2 + a_y^2 + a_z^2}$$

The direction of the object's acceleration can be determined via right-angle trigonometry

Doing kinematics in multiple dimensions involves a concerted effort on your part to disregard motion in one direction when considering motion in a perpendicular direction. The ability to mentally break down a complicated motion into its component motions requires considerable practice.

---

[1] Actually, at extremely high speeds the horizontal and vertical motions are not independent. At speeds comparable to the speed of light, the interdependence between horizontal and vertical motion (because of [time dilation](#)) becomes noticeable.

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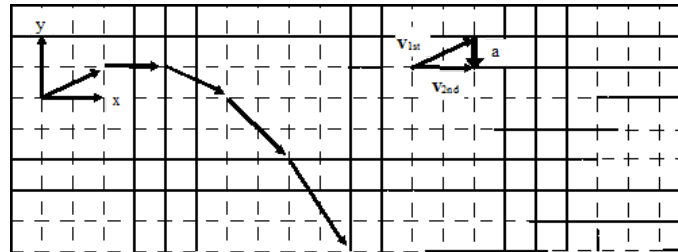
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## 02. Drawing Motion Diagrams

Motion diagrams are of crucial importance in investigating scenarios involving multi-dimensional motion. For example, analyze the following scenario:

*In the shot put, a large mass is thrown at an angle of  $22^\circ$  above horizontal, from a position of 2 m above the ground, a horizontal distance of 25 m.*

A motion diagram for this scenario is sketched below.



- Horizontal (x) and vertical (y) coordinate systems are clearly indicated.
- The acceleration is determined by the same method as in one-dimensional motion. In this case, the acceleration was determined near the beginning of the motion. Determining the acceleration at any other time will also indicate that its direction is straight downward
- In constructing the motion diagram, only a portion of the entire motion of the shot put is illustrated. For this motion diagram, analysis will begin **the instant after the shot put leaves the putter's hand**, and analysis will end **the instant before the shot put hits the ground**. It is of extreme importance to clearly understand the beginning and the end of the motion that you will describe. The acceleration of the shot put while in the putter's hand, and the acceleration upon contact with the ground, has been conveniently left out of this analysis. Unless explicit information is either provided or desired about these accelerations, it is best to focus analysis on the simplest essentials of the motion.

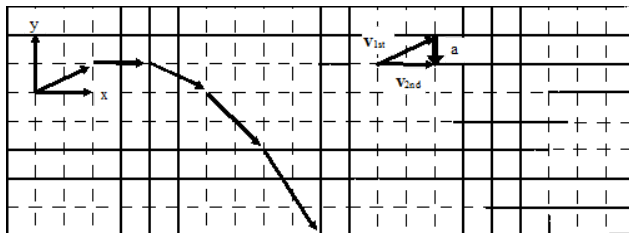
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## 03. Drawing Motion Graphs

1. Drawing the position vs. time graph
  1. Horizontal Position
  2. Vertical Position
2. Drawing the velocity vs. time graph
3. Drawing the acceleration vs. time graph

Let's look at the situation again:

*In the shot put, a large mass is thrown at an angle of  $22^\circ$  above horizontal, from a position of 2 m above the ground, a horizontal distance of 25 m.*



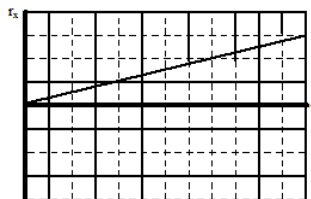
The verbal representation of the situation has already been translated into a pictorial representation, the motion diagram. A careful reading of the motion diagram allows the construction of the motion graphs.

### Drawing the position vs. time graph

First, examine the position of the shot put as it moves through the air. Remember, the analysis of the horizontal position must be independent of the analysis of the vertical position.

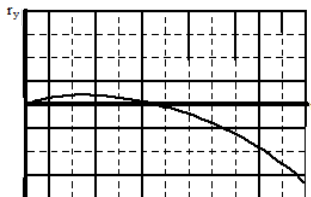
#### Horizontal Position

From the motion diagram, the shot put starts at position zero, and then has positive, increasing positions throughout the remainder of its motion. The horizontal position increases by even amounts in even time intervals.

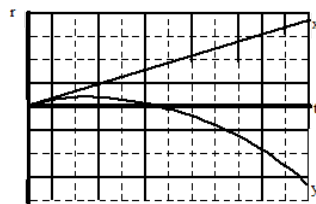


#### Vertical Position

The shot put starts at position zero, increases its vertical position at a rate that is decreasing, then begins to decrease its vertical position at a rate that is increasing, even as it drops to negative positions.

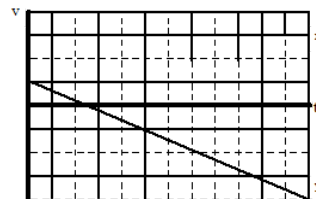


Typically, both the horizontal and vertical positions are displayed on the same axis.



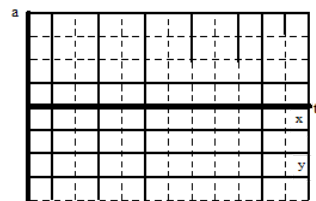
## Drawing the velocity vs. time graph

In the horizontal direction, the rate at which the position changes is constant. Hence, the horizontal component of velocity is constant, and positive. In the vertical direction, the velocity component begins positive, decreases to zero, and then increases in the negative direction.



## Drawing the acceleration vs. time graph

From the motion diagram, the acceleration of the shot-put can be determined to be directed downward at every point. Thus, the horizontal component of acceleration is zero and the vertical component is negative, and approximately constant due to our model's approximations.



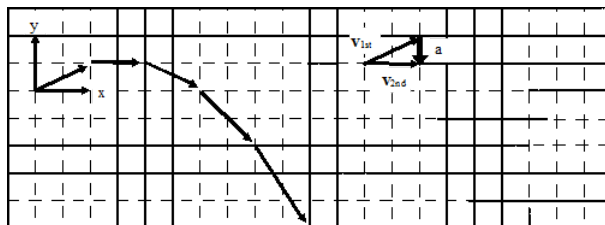
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## 04. Tabulating Motion Information & Doing The Math

1. Tabulating Motion Information
2. Doing the Math

### Tabulating Motion Information

In the shot put, a large mass is thrown at an angle of  $22^\circ$  above horizontal, from a position of 2 m above the ground, a horizontal distance of 25 m.

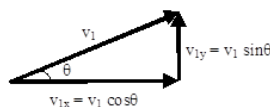


Now that you have constructed a motion diagram and motion graphs, you should be able to assign numerical values to several of the kinematic variables. A glance at the situation description should indicate that information is presented about the shot put at two distinct events: when the shot put leaves the putter's hand and when the shot put strikes the ground. Other information can also be determined about these events by referencing the motion diagram. To organize this information, you should construct a motion table.

Event 1: The instant after the shot put leaves the hand.	Event 2: The instant before the shot put hits the ground.
$t_1 = 0 \text{ s}$	$t_2 =$
$r_{1x} = 0 \text{ m}$	$r_{2x} = +25 \text{ m}$
$r_{1y} = 0 \text{ m}$	$r_{2y} = -2 \text{ m}$
$v_{1x} = v_1 \cos 22^\circ$	$v_{2x} =$
$v_{1y} = v_1 \sin 22^\circ$	$v_{2y} =$
$a_{12x} = 0 \text{ m/s}^2$	
$a_{12y} = -9.8 \text{ m/s}^2$	

In addition to the information explicitly given (the initial and final positions), information is available about both the initial velocity and the acceleration.

- **Initial velocity:** Although the magnitude of the initial velocity ( $v_1$ ) is unknown, its orientation in space is known. Thus, via the right-angle trigonometry shown below, the components of this unknown magnitude velocity in the horizontal and vertical directions can be determined. Since we will analyze the x- and y-motion separately, we *must* break the initial velocity into its x- and y-components.



- **Acceleration:** The only force acting on the shot-put *during the time interval of interest* is the force of gravity, which acts directly downward. This is because the analysis of the motion is restricted to the time interval *after* leaving the thrower's hand and *before* striking the ground. Thus, there is no horizontal acceleration of the shot-put and the vertical acceleration has a magnitude of  $9.8 \text{ m/s}^2$ .

### Doing the Math



In Model 1, you were presented with two kinematic relationships. These relationships are valid in both the horizontal and vertical directions. Thus, you have a total of four relationships with which to analyze the scenario given. In the example above, there are four unknown kinematic variables. You should remember from algebra that four equations are sufficient to calculate four unknowns. Thus, by applying the kinematic relations in both the horizontal and vertical directions, you should be able to determine the initial velocity of the shot-put, the time in the air, and the final horizontal and vertical velocity components.

First, let's examine the horizontal component of the motion. Note that the positions, velocities, and accelerations in the following equations are all horizontal components.

#### x-direction

$$\begin{aligned}
 r_2 &= r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2 \\
 25 &= 0 + (v_1 \cos 22)(t_2 - 0) + \frac{1}{2}(0)(t_2 - 0)^2 \\
 v_2 &= v_1 + a_{12}(t_2 - t_1) & 25 &= 0.927v_1t_2 \\
 v_{2x} &= (v_1 \cos 22) + 0(t_2 - 0) & v_1 &= \frac{27.0}{t_2} \\
 v_{2x} &= 0.927v_1
 \end{aligned}$$

Now let's examine the vertical component of the motion. All the positions, velocities, and accelerations in the following equations are now vertical components.

#### y-direction

$$\begin{aligned}
 r_2 &= r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2 \\
 v_2 &= v_1 + a_{12}(t_2 - t_1) & -2 &= 0 + (v_1 \sin 22)(t_2 - 0) + \frac{1}{2}(-9.8)(t_2 - 0)^2 \\
 v_{2y} &= (v_1 \sin 22) + (-9.8)(t_2 - 0) & -2 &= 0.375v_1t_2 - 4.9t_2^2 \\
 v_{2y} &= 0.375v_1 - 9.8t_2
 \end{aligned}$$

Substituting the value of  $v_1$  from above yields:

$$\begin{aligned}
 -2 &= 0.375\left(\frac{27.0}{t_2}\right)t_2 - 4.9t_2^2 \\
 -2 &= 10.1 - 4.9t_2^2 \\
 -12.1 &= -4.9t_2^2 \\
 2.47 &= t_2^2 \\
 t_2 &= 1.57s
 \end{aligned}$$

Plugging  $t_2 = 1.57$  s into all of the remaining equations gives:

$$v_1 = 17.2 \text{ m/s } v_{2x} = 15.9 \text{ m/s } v_{2y} = -8.94 \text{ m/s}$$

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## 05. Hints and Suggestions

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### 1. Selecting Events

### Selecting Events

Let's look again at the shot-putter.

*In the shot put, a large mass is thrown at an angle of  $22^\circ$  above horizontal, from a position of 2 m above the ground, a horizontal distance of 25 m.*

Imagine a videotape of the shot put event. Fast-forward over the frames showing the shot putter picking up the shot and stepping into the ring. Begin to watch the imaginary videotape frame-by-frame as the shot putter begins to push the shot off of her shoulder and forward. Stop the videotape on the frame when the shot first leaves the putter's hand.

Why is it so important that we begin the analysis at this frame and **explicitly disregard** all the motion that has taken place before this frame? The reason is that in every frame preceding this frame, the shot put was in contact with the putter. Thus, the putter was exerting a force on the shot. Since no information is presented concerning this force, we have no way to determine the acceleration during these frames and hence no way to determine any other kinematic variables. Thus, we disregard all motion preceding the instant the shot leaves the putter's hand because that portion of the motion is simply impossible to analyze with the information provided. Once the shot leaves her hand, the only force acting on the shot is the force of gravity, which greatly simplifies the analysis.

Continue playing the imaginary videotape forward. Begin playing the tape frame-by-frame as the shot approaches the ground. Stop the videotape the frame before the shot hits the ground. We will stop our analysis at this frame. Why? Because starting with the next frame, the shot is in contact with the ground. Once in contact with the ground, an additional, unknown magnitude force begins to act on the shot. Once an unknown magnitude force begins to act, the acceleration of the shot becomes unknown and we are stuck. Thus, we conveniently stop our analysis before things get too complicated!

Since our analysis stops the instant before contact, note that the shot is still moving at this instant. (If it wasn't, how could it ever reach the ground?) Thus, resist the temptation to think that the velocity of the shot is zero at the end of analysis. The velocity of the shot is ultimately equal to zero (after it makes a big divot into the ground) but that happens long after it strikes the ground and hence long after our analysis is finished.

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## 01. Selected Answers For Kinematics

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## 01. Selected Answers For Kinematics

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## 03. Dynamics

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- [01. Concepts and Principles](#)
  - [02. Drawing Free-Body Diagrams](#)
  - [03. Applying Newton's Second Law](#)
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## 01. Concepts and Principles

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### Concepts and Principles

Just like in kinematics, it's an empirical fact about nature that when a force acts on an object in one direction (for example, the horizontal) this action does not appear to cause changes in the motion in a perpendicular direction (the vertical). Therefore, to investigate the effects of forces on the motion of an object in the vertical direction, you can ignore all forces acting in the horizontal direction. Of course, many forces will simultaneously act in both the horizontal and vertical directions. As in kinematics, the effect of these forces can be examined by concentrating on the *components* of the forces in the various directions. Again, as long as the directions of interest are perpendicular, the force components can be determined through right-angle trigonometry, and the magnitude of the force can always be determined by:

$$F = \sqrt{F_x^2 + F_y^2 + F_z^2}$$

Thus, Newton's second law,

$$\Sigma F = ma$$

is independently valid in any member of a set of perpendicular directions. The total force in the horizontal direction, for example, is equal to the mass times the acceleration in that direction. Note that the mass has been verified to be independent of direction, meaning that objects possess the same inertia in all directions.

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## 02. Drawing Free-Body Diagrams

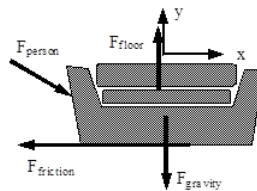
1. Drawing Free-Body Diagrams
2. The Force of Friction
  1. Kinetic friction
  2. Static friction

### Drawing Free-Body Diagrams

The free-body diagram is still the most important analysis tool for determining the forces that act on a particular object. As an example, start with a verbal description of a situation:

*While rearranging furniture, a 600 N force is applied at an angle of 25° below horizontal to a 100 kg sofa at rest.*

A free-body diagram for the sofa is sketched below:



This is a complete free-body diagram for the couch.

The only non-contact interaction is the force of gravity, directed vertically downward.

The couch is in contact with two external objects, the person, pushing the couch across the floor, and the floor, exerting a force directed upward to prevent the couch from sinking into the floor. In addition, experience tells us that the interaction between the couch and the floor also hinders the motion of the couch in the direction of the person's push. This portion of the couch-floor interaction is commonly referred to as *friction*.

### The Force of Friction

The interaction between objects in direct contact typically consists of two parts. One part of the interaction is directed perpendicular to the surface of contact. [1] The other part of the interaction is the portion commonly called friction. The frictional portion of the interaction depends on many variables.

For most situations, a *model* of friction limiting the number of variables effecting the interaction to two is adequate. These two variables are the magnitude of the perpendicular portion of the interaction, generically called the *contact force*, and a unit-less constant that reflects the relative roughness of the surface-to-surface contact, termed the *coefficient of friction*. This linear model of sliding friction further differentiates between the frictional interaction when the two surfaces are moving with respect to each other, termed *kinetic friction*, and when they are not, termed *static friction*.

#### Kinetic friction

The kinetic friction model states that the frictional interaction between the surfaces is approximately equal to the product of the contact force,  $F_{\text{contact}}$  and the coefficient of friction for kinetic situations,  $\mu_k$ ;

$$F_{\text{friction}} = \mu_k F_{\text{contact}}$$

The direction of this force on a particular object is in opposition to the relative motion of the two surfaces in contact.

#### Static friction

The static friction model states that the frictional interaction between the surfaces must be less than, or at most equal to, the product of the contact force,  $F_{\text{contact}}$ , and the coefficient of friction for static situations,  $\mu_s$ ;

$$F_{\text{friction}} \leq \mu_s F_{\text{contact}}$$

The direction of this force on a particular object is in opposition to the motion that *would* result if the frictional interaction were not present.

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[1]The portion of the interaction directed perpendicular to the surface of contact is sometimes referred to as the *normal* force, where normal has its mathematical definition of perpendicular.

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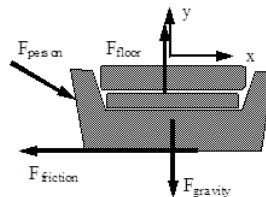
## 03. Applying Newton's Second Law

### Applying Newton's Second Law

Using this model for friction, we can now quantitatively analyze the original situation. Note that the two coefficients of friction will typically be given as an ordered pair,  $(\mu_s, \mu_k)$ .

*While rearranging furniture, a 600 N force is applied at an angle of 25° below horizontal to a 100 kg sofa at rest. The coefficient of friction between the sofa and the floor is (0.5, 0.4).*

A free-body diagram for the sofa is sketched below:



Hopefully you realize that two quite different outcomes can result from this push. Either the person pushes too weakly to move the couch or the push is sufficient to make the couch move. Since the frictional forces acting in these two cases are quite different, we can't really numerically analyze the situation until we make an *assumption* as to the outcome of the push. Of course, we will then have to check the validity of our assumption once we have completed our analysis. If our assumption turns out to be incorrect, we will then have to re-analyze the situation using the other possible outcome.

I will assume the couch doesn't move for the analysis below, and then later check the assumption. Assuming the couch doesn't move is equivalent to assuming  $a_x = 0 \text{ m/s}^2$  and that the relevant type of friction to use is static friction.

Applying Newton's Second Law independently in the horizontal (x) and vertical (y) directions yields:

#### x-direction   y-direction

$$\begin{array}{ll} \Sigma F = ma & \Sigma F = ma \\ F_{\text{person}} \cos(25^\circ) - F_{\text{static friction}} = ma_x & -F_{\text{person}} \sin(25^\circ) - F_{\text{gravity}} + F_{\text{floor}} = ma_y \\ 600 \cos(25^\circ) - F_{\text{static friction}} = 100(0) & -600 \sin(25^\circ) - (100)(9.8) + F_{\text{floor}} = 100(0) \\ F_{\text{static friction}} = 544 \text{ N} & F_{\text{floor}} = 1,234 \text{ N} \end{array}$$

Notice that the acceleration of the couch in the vertical direction must be zero regardless of my assumption, unless the couch begins to levitate or crash through the floor.

Assuming the couch doesn't move leads to a calculated value of static friction equal to 544 N. Can static friction create a force of this magnitude to prevent the couch's motion? I can check this calculated value against the allowed values for static friction:

$$\begin{array}{l} F_{\text{static friction}} \leq \mu_s F_{\text{floor}} \\ F_{\text{static friction}} \leq (0.5)(1,234) \\ F_{\text{static friction}} \leq 617 \text{ N} \end{array}$$

Since the calculated value of the static frictional force is below the maximum possible value of the static frictional force, my analysis and assumption are valid, the couch does not budge. The person is not pushing hard enough to overcome the static frictional force that acts to prevent the couch's motion relative to the floor.

Therefore, in this scenario the actual value of the static frictional force is 544 N (remember, it can be any value less than or equal to 617 N) and the acceleration of the couch is equal to zero.

How would the analysis change if the couch *was* initially in motion? Assume you enlisted a friend to help get the couch moving, but as soon as it began to move your friend stopped pushing. Would the couch stop immediately, gradually slow down to a stop, or could you keep the couch in motion across the room?

If the couch was initially moving, two things must change in our analysis. First, the horizontal acceleration of the couch is no longer necessarily zero. Second, the frictional force acting on the couch is kinetic.

Applying Newton's Second Law independently in the horizontal (x) and vertical (y) directions now yields:

**x-direction y-direction**

$$\begin{array}{ll} \Sigma F = ma & \Sigma F = ma \\ F_{\text{person}} \cos(25) - F_{\text{kinetic friction}} = 100a_x & -F_{\text{person}} \sin(25) - F_{\text{gravity}} + F_{\text{floor}} = ma_y \\ 600 \cos(25) - F_{\text{kinetic friction}} = 100a_x & -600 \sin(25) - (100)(9.8) + F_{\text{floor}} = 100(0) \\ 544 - F_{\text{kinetic friction}} = 100a_x & F_{\text{floor}} = 1,234 \text{ N} \end{array}$$

To finish the analysis, we need to calculate the kinetic frictional force.

$$\begin{array}{l} F_{\text{kinetic friction}} = \mu_k F_{\text{floor}} \\ F_{\text{kinetic friction}} = (0.4)(1,234) \\ F_{\text{kinetic friction}} = 494 \text{ N} \end{array}$$

Substituting this into the x-equation above yields:

$$\begin{array}{l} 544 - 494 = 100a_x \\ 50 = 100a_x \quad 50 = 100a_x \\ a_x = 0.50 \text{ m/s}^2 \end{array}$$

Thus, if the couch is already moving the kinetic frictional force is 494 N and the couch accelerates toward the right at  $0.50 \text{ m/s}^2$ . In summary, if the couch is initially moving it will continue to move and accelerate at  $0.50 \text{ m/s}^2$  to the right. If it is initially at rest, the person pushing on it will not be able to get it to move.

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## 04. Choosing a Coordinate System

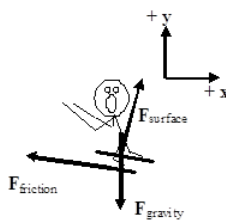
### 1. Choosing a Coordinate System

#### Choosing a Coordinate System

In analyzing a scenario, you are always free to choose whatever coordinate system you like. If you make up negative, or left positive, this will not make you get the wrong answer. However, certain coordinate systems may make the mathematical analysis simpler than other coordinate systems. For example;

*A 75 kg skier starts from rest at the top of a  $20^\circ$  slope. He's a show-off, so he skies down the hill backward. The frictional coefficient between his skis and the snow is (0.10,0.05).*

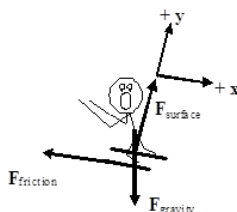
In attempting to analyze this situation, first draw a free-body diagram.



Notice that I have chosen the traditional horizontal and vertical coordinate system. I could analyze the situation using this coordinate system, but there are two difficulties with this choice.

1. Neither the force of the surface nor the force of friction is oriented in the x- or y-direction. (The force of gravity is oriented in the negative y-direction.) Therefore, I will have to use trigonometry to determine the x- and y-components of both of these forces.
2. The skier is accelerating down the inclined slope. Thus, I will also need trigonometry to determine the x- and y-components of the acceleration.

Although these difficulties are by no means insurmountable, why make the task more difficult than it has to be?



Contrast the above choice of coordinate system with a coordinate system in which the x-direction is tilted parallel to the surface on which the skier slides and the y-direction, remaining perpendicular to the x, is perpendicular to the surface.

1. Using the tilted coordinate system, the only force not oriented in the x- or y-direction is the force of gravity. Therefore, I will only need to use trigonometry to determine the x- and y-components of one force rather than two.
2. The skier is accelerating down the inclined slope. Since the x-direction is oriented parallel to the slope, the skier has an acceleration in the x-direction and zero acceleration in the y-direction.

This simple rotation of the coordinate system has made the mathematical analysis of this situation much easier. Applying Newton's second law in the x- and y-direction leads to:

### x-direction

$$\Sigma F = m a$$

$$F_{\text{gravity}} (\sin 20^\circ) - F_{\text{friction}} = m a_x$$

$$(75)(9.8)(\sin 20^\circ) - F_{\text{friction}} = 75 a_x$$

$$251 - F_{\text{friction}} = 75 a_x$$

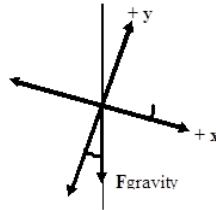
### y-direction

$$\Sigma F = m a$$

$$-F_{\text{gravity}} (\cos 20^\circ) + F_{\text{surface}} = m (0)$$

$$-(75)(9.8)(\cos 20^\circ) + F_{\text{surface}} = m (0)$$

$$F_{\text{surface}} = 691 \text{ N}$$



Notice that if the x-axis is rotated by  $20^\circ$  from horizontal to become parallel to the slope, the y-axis is rotated by  $20^\circ$  from vertical. Since the force of gravity is always oriented vertically downward, it's now  $20^\circ$  from the y-axis.

Thus, the force of gravity has a component in the positive x-direction of  $F_{\text{gravity}} (\sin 20^\circ)$  and a component in the negative y-direction of  $F_{\text{gravity}} (\cos 20^\circ)$ .

Now that the contact force between the skier and the slope is known, the static friction force can be determined.

$$F_{\text{friction}} \leq \mu_s F_{\text{surface}}$$

$$F_{\text{friction}} \leq (0.10)(691)$$

$$F_{\text{friction}} \leq 69 \text{ N}$$

Since the x-component of the force of gravity on the skier (251 N) is larger than the force of static friction (69 N), the skier will accelerate down the hill. Once he begins to move, the frictional force must be calculated using the kinetic friction model.

$$F_{\text{friction}} = \mu_k F_{\text{surface}}$$

$$F_{\text{friction}} = (0.05)(691)$$

$$F_{\text{friction}} = 35 \text{ N}$$

Examining the x-component of Newton's second law:

$$251 - F_{\text{friction}} = 75 a_x$$

$$251 - 35 = 75 a_x$$

$$a_x = 2.9 \text{ m/s}^2$$

The skier accelerates down the slope with an acceleration of  $2.9 \text{ m/s}^2$ .

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## 05. Activities

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## 04. Conservation

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- [02. Analysis Tools](#)
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- [04. Activities](#)

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## 01. Concepts and Principles

1. [The Impulse-Momentum Relation](#)
2. [The Work-Energy Relation](#)

### The Impulse-Momentum Relation

Just like the kinematic relations and Newton's second law, the impulse-momentum relation is independently valid in any member of a set of perpendicular directions. Thus, we will typically apply the impulse-momentum relation in its component forms:

$$\begin{aligned}P_i + J &= P_f \\mv_{xi} + \Sigma(F_x(\Delta t)) &= mv_{xf} \\mv_{yi} + \Sigma(F_y(\Delta t)) &= mv_{yf} \\mv_{zi} + \Sigma(F_z(\Delta t)) &= mv_{zf}\end{aligned}$$

### The Work-Energy Relation

From Model 1, our expression for the Work-Energy Relation, with gravitational potential energy terms, is:

$$\begin{aligned}E_i + Work &= E_f \\ \frac{1}{2}mv_i^2 + mgh_i + \Sigma(F\|\Delta r\|\cos\phi) &= \frac{1}{2}mv_f^2 + mgh_f\end{aligned}$$

It's very important to remember that the work-energy relation is a *scalar* equation, meaning it can not be broken into components and "solved" separately in the x-, y-, and z-directions. This is even more important to remember now that we are working in multiple dimensions. This observation results in two important points:

- The work-energy relation involves the actual initial and final velocities, *not* their components. The kinetic energy of an object does not depend on the direction of travel of the object.
- In the expression for work,  $F\|\Delta r\|\cos\phi$ , the product of the *magnitude* of the force and the *magnitude* of the displacement is multiplied by  $\cos\phi$ , where  $\phi$  is defined to be the angle between the applied force and the displacement of the object. If the force and displacement are in the same direction  $\phi = 0^\circ$ , and the work is positive (the object gains energy). If the force and displacement are in the opposite direction  $\phi = 180^\circ$ , and the work is negative (the object loses energy). If the force and displacement are perpendicular, no work is done. Note that the actual directions of the force and the displacement are unimportant, only their directions *relative to each other* affect the work.

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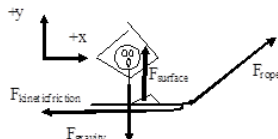
## 02. Analysis Tools

1. Applying the Impulse-Momentum Relation to a Single Object
2. Applying the Work-Energy Relation to a Single Object

### Applying the Impulse-Momentum Relation to a Single Object

Let's investigate the following scenario:

*A boy pulls a 30 kg sled, including the mass of his kid brother, along ice. The boy pulls on the tow rope, oriented at  $60^\circ$  above horizontal, with a force of 110 N until his kid brother begins to cry. Like clockwork, his brother always cries upon reaching a speed of 2.0 m/s. The frictional coefficient is (0.20,0.15).*



Event 1: The instant before the sled begins to move.

Event 2: The instant the sled reaches 2.0 m/s.

Let's apply the impulse-momentum relations to see what they reveal about the situation.

Applying impulse-momentum separately in the x- and y-directions yields:

**x-direction y-direction**

$$\begin{aligned}
 mv_{yi} + \Sigma(F_y(\Delta t)) &= mv_{yf} \\
 30(0) + F_{\text{rope}} \sin 60(\Delta t) - F_{\text{gravity}}(\Delta t) + F_{\text{surface}}(\Delta t) &= 30(0) \\
 0 + 110 \sin 60(\Delta t) - (30)(9.8)(\Delta t) + F_{\text{surface}}(\Delta t) &= 0 \\
 95(\Delta t) - 294(\Delta t) + F_{\text{surface}}(\Delta t) &= 0 \\
 95 - 294 + F_{\text{surface}} &= 0 \\
 F_{\text{surface}} &= 199 \text{ N}
 \end{aligned}$$

$$\begin{aligned}
 mv_{xi} + \Sigma(F_x(\Delta t)) &= mv_{xf} \\
 30(0) + F_{\text{rope}} \cos 60(\Delta t) - F_{\text{fr}}(\Delta t) &= 30(2) \\
 0 + 110 \cos 60(\Delta t) - 0.15 F_{\text{surface}}(\Delta t) &= 60 \\
 55(\Delta t) - 0.15 F_{\text{surface}}(\Delta t) &= 60
 \end{aligned}$$

Substituting the value for the force of the surface into the x-equation,

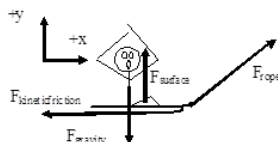
$$\begin{aligned}
 55(\Delta t) - 0.15(199)(\Delta t) &= 60 \\
 55(\Delta t) - 30(\Delta t) &= 60 \\
 25(\Delta t) &= 60 \\
 \Delta t &= 2.4 \text{ s}
 \end{aligned}$$

The kid brother begins to cry after only 2.4 s.

### Applying the Work-Energy Relation to a Single Object

What will the work-energy relation tell us about the same scenario?

*A boy pulls a 30 kg sled, including the mass of his kid brother, along ice. The boy pulls on the tow rope, oriented at  $60^\circ$  above horizontal, with a force of 110 N until his kid brother begins to cry. Like clockwork, his brother always cries upon reaching a speed of 2.0 m/s. The frictional coefficient is (0.20,0.15).*



Event 1: The instant before the sled begins to move.

Event 2: The instant the sled reaches 2.0 m/s.

$$\begin{aligned}\frac{1}{2}mv_i^2 + mgh_i + \Sigma(F\|\Delta r\|\cos\phi) &= \frac{1}{2}mv_f^2 + mgh_f \\ 0 + 0 + F_{\text{rope}}(\Delta r)\cos 60 + F_{\text{surface}}(\Delta r)\cos 90 + F_{\text{kinetic friction}}(\Delta r)\cos 180 &= \frac{1}{2}(30)(2)^2 + 0 \\ (110)(\Delta r)\cos 60 - F_{\text{kinetic friction}}(\Delta r) &= 60 \\ 55(\Delta r) - 0.15F_{\text{surface}}(\Delta r) &= 60\end{aligned}$$

Notice that the force of the surface does no work, the force of the rope does positive work, and the force of friction does negative work. Each of these terms should make sense if you remember that work is the transfer of energy into (positive) or out of (negative) the system of interest. Also recall that in this form of the work-energy relation we conceptualize gravity as a source of potential energy, not as a force that does work.

Using the result for the force of the surface determined in the first example,  $F_{\text{surface}} = 199 \text{ N}$ , gives:

$$\begin{aligned}55(\Delta r) - 0.15F_{\text{surface}}(\Delta r) &= 60 \\ 55(\Delta r) - 0.15(199)(\Delta r) &= 60 \\ 55(\Delta r) - 30(\Delta r) &= 60 \\ 25(\Delta r) &= 60 \\ \Delta r &= 2.4 \text{ m}\end{aligned}$$

The kid brother begins to cry after traveling 2.4 m.

---

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## 03. Analysis Tools 2

### Applying Work-Energy with Gravitational Potential Energy

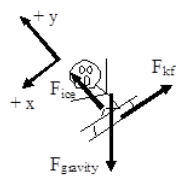
Let's use the work-energy relation, with gravitational potential energy terms, to analyze the following scenario:

*A 30 kg child on his 15 kg sled slides down his parents' 10 m long, 15° above horizontal driveway after an ice storm. The coefficient of friction between the sled and the driveway is (0.10, 0.08).*

Event 1: The instant before the sled begins to move.

Event 2: The instant the sled reaches the bottom of the driveway.

To calculate the gravitational energy terms, let the bottom of the driveway be zero and up positive. The coordinate system used to calculate gravitational energy does not in general have to be the same as the system you use for the rest of the problem. In fact, since the work-energy relation is a scalar equation, the other portions of the equation should not depend on your choice of coordinate system at all!



$$\frac{1}{2}mv_i^2 + mgh_i + \Sigma(F\|\Delta r\|\cos\phi) = \frac{1}{2}mv_f^2 + mgh_f$$

$$0 + 45(9.8)(10\sin 15) + F_{ice}(10)\cos 90 + F_{kf}(10)\cos 180 = \frac{1}{2}45v_2^2 + 0$$

$$1140 - 10F_{kf} = 22.5v_2^2$$

Note:

- The only forces that *could* do work are the force of the ice and the force of friction, since the action of the force of gravity is already incorporated into the gravitational potential energy terms.
- The heights in the gravitational potential energy function were measured from the bottom of the driveway, with the positive direction as upward, as required. Notice that the initial height is not the same as the length of the driveway. Since the driveway is 10 m long, at an angle of 15°, the height of the top of the driveway relative to the bottom is (10 m) sin 15°. The height at the bottom of the driveway is defined to be 0 m.

To finish the analysis we need to determine the kinetic frictional force. Since this depends on the force of the ice, apply Newton's Second Law in the y-direction and find:

$$\Sigma F = ma$$

$$+ F_{ice} - F_{gravity}\cos 15 = 45(0)$$

$$F_{ice} - (45)(9.8)\cos 15 = 0$$

$$F_{ice} = 426\text{ N}$$

$$F_{kf} = \mu_k F_{ice}$$

$$F_{kf} = (0.08)(426)$$

$$F_{kf} = 34\text{ N}$$

Plugging this value into the work-energy relation yields:

$$1140 - 10(34) = 22.5v_2^2$$

$$1140 - 340 = 22.5v_2^2$$

$$800 = 22.5v_2^2$$

$$v_2 = 5.96\text{ m/s}$$

### A Two-Dimensional Collision

Let's try a two-dimensional collision.

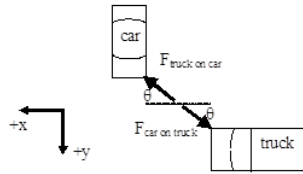
*At a busy intersection, an impatient driver heading south runs a red-light and collides with a delivery truck originally moving at 15 m/s west. The vehicles become entangled and the skid marks from the wreckage are at 22° south of west. The auto mass is 755 kg and the truck mass is 1250 kg.*

I'll choose:

Event 1: The instant before the car and truck collide.

Event 2: The instant the car and truck reach a common velocity

Partial free-body diagrams (top view) for both the car and the truck during this time interval are shown below.



These are only partial free-body diagrams because:

- Forces perpendicular to the earth's surface (the force of gravity and the force of the road) are not shown.
- During a collision, the force *between* the colliding objects is normally much greater in magnitude than any other forces acting on the objects. Therefore we will often ignore the other forces acting on colliding objects for the duration of a collision. This approximation is termed the *impulse approximation*. Under the impulse approximation, the frictional forces between the car and truck and the road are ignored.

Also note that the direction of the force acting between the car and truck is unknown. The angle  $\theta$  is not determined from the situation description.

Applying the impulse-momentum relation to the car and truck yields:

#### Car

##### x-direction y-direction

$$mv_i + \Sigma F(\Delta t) = mv_f \quad mv_i + \Sigma F(\Delta t) = mv_f$$

$$755(0) + F_{\text{truck on car}} \cos \theta(\Delta t) = 755(v_2 \cos 22) \quad 755(v_{\text{car}}) - F_{\text{truck on car}} \sin \theta(\Delta t) = 755(v_2 \sin 22)$$

#### Truck

##### x-direction y-direction

$$mv_i + \Sigma F(\Delta t) = mv_f \quad mv_i + \Sigma F(\Delta t) = mv_f$$

$$1250(15) - F_{\text{car on truck}} \cos \theta(\Delta t) = 1250(v_2 \cos 22) \quad 1250(15) - F_{\text{car on truck}} \sin \theta(\Delta t) = 1250(v_2 \sin 22)$$

Since the magnitude of the force on the car due to the truck and the force on the truck due to the car are equal, when the x-equations for the car and truck are added, the impulses cancel!

##### x-direction

$$755(0) + 1250(15) = 755(v_2 \cos 22) + 1250(v_2 \cos 22)$$

$$18750 = 2005(v_2 \cos 22)$$

$$v_2 = 10.1 \text{ m/s}$$

This is the speed of the wreckage immediately after the collision. Note that this is exactly the same equation we would have written if we had considered the *system* of the car and truck right from the start.

Adding the two y-equations yields:

##### y-direction

$$755(v_{\text{car}}) + 1250(0) = 755(v_2 \sin 22) + 1250(v_2 \sin 22)$$

$$755(v_{\text{car}}) = 2005(10.1 \sin 22)$$

$$v_{\text{car}} = 10.0 \text{ m/s}$$

This is the speed of the car immediately before colliding with the truck.

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## 04. Activities

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## 01. Selected Answers For Conservation

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## CHAPTER OVERVIEW

### Model 3: The Particle Model

#### Topic hierarchy

- 01. Model Specifies
- 02. Kinematics
  - 01. Concepts and Principles
  - 02. Using the Calculus
  - 03. Using the Calculus Another Example
  - 04. Circular Motion
  - 05. Circular Motion 2
  - 06. The Kinematics of Circular Motion
  - 07. Activities
- 04. Conservation
  - 01. Concepts and Principles
  - 02. Applying the Impulse-Momentum Relation
  - 03. Elastic Potential Energy
  - 04. Applying Newton's 2nd Law to a Spring-Mass System
  - 05. Potential Energy Functions
  - 06. Activities
- 3.3: Dynamics
  - 01. Concepts and Principles
  - 02. Time-Dependent Forces
  - 03. Applying Newton's Second Law to Circular Motion
  - 04. Activities

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## 01. Model Specifies

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### Model Specifics

For our third pass through the study of mechanics, we will finally allow forces to vary in both magnitude and direction. Peeling away this approximation requires us to utilize a new mathematical tool, calculus. We will still, however, require the following two approximations.

**The object's size and shape are unimportant.**

**The object is classical.**

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## 02. Kinematics

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- [01. Concepts and Principles](#)
  - [02. Using the Calculus](#)
  - [03. Using the Calculus Another Example](#)
  - [04. Circular Motion](#)
  - [05. Circular Motion 2](#)
  - [06. The Kinematics of Circular Motion](#)
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- 

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## 01. Concepts and Principles

1. Position ( $r$ )
2. Velocity ( $v$ )
3. Acceleration ( $a$ )

If the forces acting on an object are not constant, then the acceleration of the object is not constant. To analyze the kinematics of an object undergoing non-constant acceleration requires the use of calculus. By re-examining our original definitions, valid in the limit of very small time intervals during which the acceleration is *approximately* constant, the relationships between position, velocity, and acceleration can be constructed in terms of the derivative.

### Position ( $r$ )

Let  $r(t)$  be the location of the object at every time,  $t$ , in the time interval of interest.

### Velocity ( $v$ )

Our original definition of velocity,

$$v = \frac{r_{\text{final}} - r_{\text{initial}}}{t_{\text{final}} - t_{\text{initial}}}$$

remains valid if the acceleration is constant. The acceleration will always be constant in the limit of infinitesimally small time intervals. By the fundamental definition of calculus, the above expression, in the limit of infinitesimally small time intervals, becomes the derivative of the position function. Thus,  $v(t)$ , the velocity of the object at every time,  $t$ , is defined to be the derivative of the position function,  $r(t)$ .

$$v(t) = \frac{dr(t)}{dt}$$

### Acceleration ( $a$ )

Our original definition of acceleration,

$$a = \frac{v_{\text{final}} - v_{\text{initial}}}{t_{\text{final}} - t_{\text{initial}}}$$

also remains valid only in the limit of infinitesimally small time intervals. By the fundamental definition of calculus, the above expression, in the limit of infinitesimally small time intervals, becomes the derivative of the velocity function. Thus,  $a(t)$ , the acceleration of the object at every time,  $t$ , is defined to be the derivative of the velocity function,  $v(t)$ .

$$a(t) = \frac{dv(t)}{dt}$$

Thus, if the position of the object is known as a function of time, the velocity and acceleration functions can be constructed through differentiation of  $r(t)$ .

$$\begin{aligned} v(t) &= \frac{dr(t)}{dt} \\ a(t) &= \frac{dv(t)}{dt} = \frac{d^2r(t)}{dt^2} \end{aligned}$$

On the other hand, if the acceleration of the object is known as a function of time, the velocity and position functions can be constructed through anti-differentiation, or integration, of  $a(t)$ .

$$\begin{aligned} v(t) &= \int a(t) dt \\ r(t) &= \int v(t) dt = \iint a(t) dt \end{aligned}$$

An important distinction, however, is that when integrating  $a(t)$  to form  $v(t)$ , an arbitrary constant will be introduced into the expression for  $v(t)$ . This constant can often be determined from knowledge of the object's velocity at some specific instant in time. Another integration, to form  $r(t)$ , will introduce an additional arbitrary constant that can often be determined from knowledge of the object's position at some specific instant in time.

In closing, please remember that the kinematic relations that have been used throughout this course were derived assuming a constant acceleration. If the acceleration is not constant, those relations are *incorrect*, and the correct kinematic relationships *must* be determined through direct integration and differentiation.

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## 02. Using the Calculus

### 1. Using The Calculus

### Using The Calculus

Investigate the scenario described below.

*In a 100 m dash, detailed video analysis indicates that a particular sprinter's speed can be modeled as a quadratic function of time at the beginning of the race, reaching 10.6 m/s in 2.70 s, and as a decreasing linear function of time for the remainder of the race. She finished the race in 12.6 seconds.*

To analyze this situation, we should first carefully determine and define the sequence of events that take place. At each of these instants, let's tabulate what we know about the motion.

Event 1: The race begins	Event 2: She reaches 10.6 m/s	Event 3: She crosses the finish line
$t_1 = 0 \text{ s}$	$t_2 = 2.7 \text{ s}$	$t_3 = 12.6 \text{ s}$
$r_1 = 0 \text{ m}$	$r_2 =$	$r_3 = 100 \text{ m}$
$v_1 = 0 \text{ m/s}$	$v_2 = 10.6 \text{ m/s}$	$v_3 =$
$a_1 =$	$a_2 =$	$a_3 =$

Since the acceleration is no longer necessarily constant between instants of interest, it is no longer useful to speak of  $a_{12}$  or  $a_{23}$ . The acceleration, like the position and the velocity, is a *function*. What the table represents is the *value* of that function at specific instants of time.

Between event 1 and 2, the sprinter's velocity can be modeled by a generic quadratic function of time, or

$$v(t) = At^2 + B$$

Our job is to first determine (if possible) the arbitrary constants A and B, and then use this velocity function to find the position and acceleration functions.

Since the sprinter starts from rest, we can evaluate the function at  $t = 0 \text{ s}$  and set the result equal to 0 m/s:

$$v(0) = A(0)^2 + B = 0$$

$$B = 0$$

Since we also know the sprinter reaches a speed of 10.6 m/s in 2.7 s, we can evaluate the function at  $t = 2.7 \text{ s}$  and set the result equal to 10.6 m/s:

$$v(2.7) = A(2.7)^2 + 0 = 10.6$$

$$A = 1.45$$

Now that we know the two constants in the velocity function, we have a complete description of the sprinter's speed during this time interval:

$$v(t) = 1.45t^2$$

Once the velocity function is determined, we differentiate to determine her acceleration function.

$$a(t) = \frac{dv(t)}{dt}$$

$$a(t) = \frac{d}{dt}1.45t^2$$

$$a(t) = 2.90t$$

Evaluating this function at  $t = 0 \text{ s}$  and  $t = 2.7 \text{ s}$  yields  $a_1 = 0 \text{ m/s}^2$  and  $a_2 = 7.83 \text{ m/s}^2$ .

We can also integrate to determine her position function.

$$r(t) = \int v(t) dt$$

$$r(t) = \int 1.45t^2 dt$$

$$r(t) = \frac{1.45}{3}t^3 + C$$

Since we know  $r = 0$  m when  $t = 0$  s, we can determine the integration constant:

$$r(0) = \frac{1.45}{3}(0)^3 + C = 0$$

$$C = 0$$

Therefore, the position of the sprinter is given by the function:

$$r(t) = \frac{1.45}{3}t^3$$

Evaluating this function at  $t = 0$  s and  $t = 2.7$  s yields  $r_1 = 0$  m and  $r_2 = 9.51$  m.

During the second portion of the race when her speed is decreasing linearly, her acceleration is constant. Therefore, we can use the kinematic relations developed for constant acceleration, and her acceleration is simply a constant value, denoted  $a_{23}$ .

$$\begin{aligned} r_3 &= r_2 + v_2(t_3 - t_2) + \frac{1}{2}a_{23}(t_3 - t_2)^2 \\ 100 &= 9.51 + 10.6(12.6 - 2.7) + \frac{1}{2}a_{23}(12.6 - 2.7)^2 & v_3 &= v_2 + a_{23}(t_3 - t_2) \\ a_{23} &= -0.295 \text{ m/s}^2 & v_3 &= 10.6 - 0.295(12.6 - 2.7) \\ & & v_3 &= 7.68 \text{ m/s} \end{aligned}$$

She crosses the finish line running at 7.68 m/s.

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## 03. Using the Calculus Another Example

### 1. Another Example

### Another Example

Investigate the scenario described below.

*A sports car can accelerate from rest to a speed of 40 m/s while traveling a distance of 200 m. Assume the acceleration of the car can be modeled as a decreasing linear function of time, with a maximum acceleration of 10.4 m/s<sup>2</sup>.*

Event 1: The car begins from rest	Event 2: The car reaches 40 m/s
$t_1 = 0$ s	$t_2 =$
$r_1 = 0$ m	$r_2 = 200$ m
$v_1 = 0$ m/s	$v_2 = 40$ m/s
$a_1 = 10.4$ m/s <sup>2</sup>	$a_2 =$

Between event 1 and 2, the car's acceleration can be modeled by a generic linear function of time, or

$$a(t) = At + B$$

Since the acceleration is decreasing, the maximum value occurs at  $t = 0$  s,

$$\begin{aligned} a(0) &= A(0) + B = 10.4 \\ B &= 10.4 \end{aligned}$$

Since we don't know the value of the acceleration at  $t_2$ , or even the value of  $t_2$ , we can't determine A, and all we can currently say about the acceleration function is that it is given by:

$$a(t) = At + 10.4$$

Nonetheless, we can still integrate the acceleration to determine the velocity,

$$\begin{aligned} v(t) &= \int a(t) dt \\ v(t) &= \int (At + 10.4) dt \\ v(t) &= \frac{1}{2} At^2 + 10.4t + C \end{aligned}$$

Since we know  $v = 0$  m/s when  $t = 0$  s, we can determine the integration constant:

$$\begin{aligned} v(0) &= \frac{1}{2} A(0)^2 + 10.4(0) + C = 0 \\ C &= 0 \end{aligned}$$

We also know that  $v = 40$  m/s at  $t_2$ , so:

$$v(t_2) = \frac{1}{2} At_2^2 + 10.4t_2 = 40$$

This equation can't be solved, since it involves two unknowns. However, if we can generate a second equation involving the same two unknowns, we can solve the two equations simultaneously. This second equation must involve the position function of the car:

$$\begin{aligned} r(t) &= \int v(t) dt \\ r(t) &= \int \left( \frac{1}{2} At^2 + 10.4t \right) dt \\ r(t) &= \frac{1}{6} At^3 + 5.2t^2 + D \end{aligned}$$

Since we know  $r = 0$  when  $t = 0$  s, we can determine the integration constant:

$$r(0) = \frac{1}{6}A(0)^3 + 5.2(0)^2 + D = 0$$
$$D = 0$$

We also know that  $r = 200$  m at  $t_2$ , so:

$$r(t_2) = \frac{1}{6}At_2^3 + 5.2t_2^2 = 200$$

These two equations,

$$\frac{1}{2}At_2^2 + 10.4t_2 = 40$$

$$\frac{1}{6}At_2^3 + 5.2t_2^2 = 200$$

can be solved by substitution (or by using a solver). Solve the first equation for  $A$ , and substitute this expression into the second equation. This will result in a quadratic equation for  $t_2$ . The solution is  $t_2 = 7.57$  s, the time for the car to reach 40 m/s. Plugging this value back into the original equations allows you to complete the description of the car's motion.

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## 04. Circular Motion

### 1. Circular Motion

#### 1. Defining angular position, angular velocity and angular acceleration

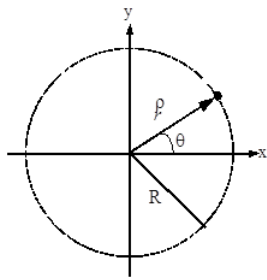
1. Angular Position
2. Angular Velocity
3. Angular Acceleration

#### 2. Deriving Relationships for velocity and acceleration

1. Position
2. Velocity
3. Acceleration

## Circular Motion

If the acceleration of an object is not constant, in either magnitude or direction, the development of a kinematic description necessitates the use of calculus. A very common class of motion, in which the acceleration is guaranteed to change in at least direction, is the motion of an object on a circular path. Let's examine general circular motion in more detail before we attempt to describe a specific situation.



The x- and y-position of the object moving along a circular path of radius R can always be described by the functions:

$$r_x(t) = R \cos \theta(t)$$

$$r_y(t) = R \sin \theta(t)$$

assuming the origin of the coordinate system is placed at the center of the circle.

### Defining angular position, angular velocity and angular acceleration

#### Angular Position

The function  $q(t)$  specifies the *angular position* of the object, and is typically measured in radians. It specifies where, *along the circle*, the object is at every instant of time. For example, if  $q(t)$  is a constant, the object doesn't move. if  $q(t)$  is a linear function of time, the object moves with constant velocity around the circle. If  $q(t)$  is a more complex function, the object speeds up or slows down as it moves around the circle.

#### Angular Velocity

The rate at which  $q(t)$  is changing,  $\frac{d\theta(t)}{dt}$ , is termed the *angular velocity* of the object, and denoted  $w(t)$ . ( $w$  is the lower-case Greek letter "omega".) Since  $w(t)$  is the rate at which the angular position is changing, it has units of rad/s.

#### Angular Acceleration

The rate at which  $w(t)$  is changing,  $\frac{d\omega(t)}{dt}$ , is termed the *angular acceleration* of the object, and denoted  $a(t)$ . ( $a$  is the lower-case Greek letter "alpha".) Since  $a(t)$  is the rate at which the angular velocity is changing, it has units  $\text{rad/s}^2$ .

## Deriving Relationships for velocity and acceleration

Now that we have the definitions of the angular quantities out of the way, let's determine the velocity and acceleration of an object undergoing circular motion. I'll begin by writing the position function in  $\hat{i}\hat{j}\hat{k}$  notation, common "short-hand" method of writing the  $\mathbf{x}$ -,

$y$ -, and  $z$ - components of a vector all together. In this notation, the  $\hat{i}$  simply stands for the  $x$ -component, the  $\hat{j}$  for the  $y$ -component, and the  $\hat{k}$  for any  $z$ -component. Hold onto your hat and try not to get lost in the calculus.

### Position

$$\mathbf{p}(t) = R(\cos \theta(t)\hat{i} + \sin \theta(t)\hat{j})$$

### Velocity

$$\mathbf{v}(t) = \frac{d}{dt} \mathbf{p}(t)$$

$$\mathbf{v}(t) = \frac{d}{dt} [R(\cos \theta(t)\hat{i} + \sin \theta(t)\hat{j})]$$

$$\mathbf{v}(t) = R\left(\frac{d}{dt} \cos \theta(t)\hat{i} + \frac{d}{dt} \sin \theta(t)\hat{j}\right)$$

$$\mathbf{v}(t) = R\left(-\sin \theta(t) \frac{d}{dt} \theta(t)\hat{i} + \cos \theta(t) \frac{d}{dt} \theta(t)\hat{j}\right)$$

$$\mathbf{v}(t) = R(-\sin \theta(t)\omega(t)\hat{i} + \cos \theta(t)\omega(t)\hat{j})$$

$$\mathbf{v}(t) = R\omega(t)(-\sin \theta(t)\hat{i} + \cos \theta(t)\hat{j})$$

### Acceleration

$$\mathbf{a}(t) = \frac{d}{dt} \mathbf{v}(t)$$

$$\mathbf{a}(t) = \frac{d}{dt} [R\omega(t)(-\sin \theta(t)\hat{i} + \cos \theta(t)\hat{j})]$$

Remembering to use the chain rule for differentiation,

$$\mathbf{a}(t) = R\left[\frac{d}{dt} \omega(t)(-\sin \theta(t)\hat{i} + \cos \theta(t)\hat{j}) + \omega(t)\left[\frac{d}{dt} (-\sin \theta(t)\hat{i} + \cos \theta(t)\hat{j})\right]\right]$$

$$\mathbf{a}(t) = R\alpha(t)(-\sin \theta(t)\hat{i} + \cos \theta(t)\hat{j}) + R\omega(t)\left(-\cos \theta(t) \frac{d}{dt} \theta(t)\hat{i} + -\sin \theta(t) \frac{d}{dt} \theta(t)\hat{j}\right)$$

$$\mathbf{a}(t) = R\alpha(t)(-\sin \theta(t)\hat{i} + \cos \theta(t)\hat{j}) + R\omega(t)(-\cos \theta(t)\omega(t)\hat{i} + -\sin \theta(t)\omega(t)\hat{j})$$

$$\mathbf{a}(t) = R\alpha(t)(-\sin \theta(t)\hat{i} + \cos \theta(t)\hat{j}) - R\omega^2(t)(\cos \theta(t)\hat{i} + \sin \theta(t)\hat{j})$$

These relationships for velocity and acceleration *look* intimidating, but are actually rather simple. (You don't have to believe me just yet...) The problem is that they are written using an awkward choice of coordinate system. In a previous Model, we used inclined coordinates for situations involving objects moving on an inclined surface. For an object moving in a circle, it's almost as if the surface upon which the object moves is continually changing its angle of incline! Perhaps we should use a coordinate system in which the orientation of the system continually changes, always keeping one axis parallel and one axis perpendicular to the motion. This is exactly what we will do. This coordinate system is referred to as the *polar coordinate system*.

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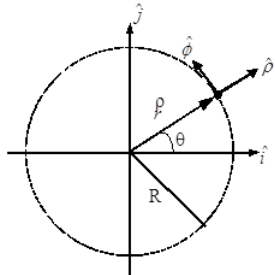
## 05. Circular Motion 2

### 1. Polar Coordinates

1. Position
2. Velocity
3. Acceleration

### Polar Coordinates

In the polar coordinate system, one axis (the radial axis, or  $\hat{\rho}$  is perpendicular to the surface of the circular path pointing radially away from the center, and the other axis (the tangential, or  $\hat{\phi}$  is parallel to the surface of the circular path pointing in the counterclockwise direction. ( $r$  is the lower-case Greek letter "rho" and  $\phi$  is "phi".)



Notice that  $\hat{\rho}$  is inclined by an angle  $\theta$  from the positive x-axis. Therefore, in terms of  $\hat{i}$  and  $\hat{j}$ ;

$$\hat{\rho} = \cos \theta(t) \hat{i} + \sin \theta(t) \hat{j}$$

$\hat{\phi}$ , on the other hand, is inclined by an angle  $\theta$  to the left of the positive y-axis. Therefore, in terms of  $\hat{i}$  and  $\hat{j}$ ;

$$\hat{\phi} = -\sin \theta(t) \hat{i} + \cos \theta(t) \hat{j}$$

Now re-examine the relationships for position, velocity and acceleration.

### Position

$$\vec{r}(t) = R(\cos \theta(t) \hat{i} + \sin \theta(t) \hat{j})$$

becomes

$$\vec{r}(t) = R \hat{\rho}$$

In component form this is:

$$\begin{aligned} r_{\rho} &= R \\ r_{\phi} &= 0 \end{aligned}$$

This means that the position of an object undergoing circular motion is *only* in the radial direction, and has a constant magnitude equal to the radius of the circle. Basically, the coordinate system is constructed so that the location of the object *defines* the radial direction.

### Velocity

$$\vec{v}(t) = R \omega(t) (-\sin \theta(t) \hat{i} + \cos \theta(t) \hat{j})$$

becomes

$$\vec{v}(t) = R \omega(t) \hat{\phi}$$

In component form this is:

$$\begin{aligned} v_{\rho} &= 0 \\ v_{\phi} &= R \omega \end{aligned}$$

This means that the velocity of an object undergoing circular motion is *only* in the tangential direction, and has a magnitude equal to the product of the radius and angular velocity. The only way an object can have a radial velocity is if the radius of its path changes, but that can't happen for an object moving along a circular path. If the object moved along an elliptical path, for example, then it would have both tangential and radial velocities.

## Acceleration

$$\vec{a}(t) = R\alpha(t)(-\sin\theta(t)\hat{i} + \cos\theta(t)\hat{j}) - R\omega^2(t)(\cos\theta(t)\hat{i} + \sin\theta(t)\hat{j})$$

becomes

$$\vec{a}(t) = R\alpha(t)\hat{\phi} - R\omega^2(t)\hat{\rho}$$

In component form this is:

$$a_{\phi} = -R\omega^2$$

$$a_{\rho} = R\alpha$$

The acceleration of an object undergoing circular motion has two components. If the object is speeding up or slowing down, the angular acceleration does not equal zero and there is an acceleration component in the tangential direction. The magnitude of the tangential acceleration is equal to the product of the radius and angular acceleration.

However, *even if the object is moving at constant speed*, there is an acceleration component in the negative radial direction, i.e., pointing *toward* the center of the circle. By virtue of traveling in a circle, the velocity vector of an object continually changes its orientation. This change in orientation is directed toward the center of the circle. Draw a motion diagram and convince yourself of this fact!

The magnitude of the radial acceleration is equal to the product of the radius and the *square* of the angular *velocity*.

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## 06. The Kinematics of Circular Motion

### 1. The Kinematics of Circular Motion

#### The Kinematics of Circular Motion

Let's try out our new tools by examining the following scenario.

*An automobile enters a U-turn of constant radius of curvature 95 m. The car enters the U-turn traveling at 33 m/s north and exits at 22 m/s south. Assume the speed of the car can be modeled as a quadratic function of time.*

To analyze this situation, we should first carefully determine and define the sequence of events that take place. At each of these events we will tabulate the position, velocity and acceleration in polar coordinates as well as the angular position, angular velocity, and angular acceleration.

Event 1: The auto enters the turn.	Event 2: The auto exits the turn.
$t_1 = 0 \text{ s}$	$t_2 =$
$r_{1p} = 95 \text{ m}$	$r_{2p} = 95 \text{ m}$
$r_{1\phi} = 0 \text{ m}$	$r_{2\phi} = 0 \text{ m}$
$\theta_1 = 0 \text{ rad}$	$\theta_2 = \pi \text{ rad}$
$v_{1p} = 0 \text{ m/s}$	$v_{2p} = 0 \text{ m/s}$
$v_{1\phi} = 33 \text{ m/s}$	$v_{2\phi} = 22 \text{ m/s}$
$\omega_1 =$	$\omega_2 =$
$a_{1p} =$	$a_{2p} =$
$a_{1\phi} =$	$a_{2\phi} =$
$\alpha_1 =$	$\alpha_2 =$

Notice that the position of the car is strictly in the radial direction and the velocity of the car is strictly in the tangential direction. It is *impossible*, for an object moving along a circular path, to have a non-zero tangential position or radial velocity.

Also notice that since the car completely reverses its direction of travel, it must have traveled halfway around a circular path. Thus,  $q_2 = \pi \text{ rad}$ .

We can quickly determine the angular velocity of the car at the two events by using the relationship:

$$v_{\phi} = R\omega$$

With  $R = 95 \text{ m}$ , the angular velocity of the car as it enters the turn is  $0.35 \text{ rad/s}$ , and as it exits the turn,  $0.23 \text{ rad/s}$ .

Between event 1 and 2, the car's speed can be modeled by a generic quadratic function of time, or

$$v_{\phi}(t) = At^2 + B$$

Since the car enters the turn at  $33 \text{ m/s}$ , we can evaluate the function at  $t = 0 \text{ s}$  and set the result equal to  $33 \text{ m/s}$ :

$$v_{\phi}(0) = A(0)^2 + B = 33$$

$$B = 33$$

The car exits the turn at  $22 \text{ m/s}$ , but since we don't know the value of  $t_2$  we can't determine  $A$ . We do know, however, that:

$$v_{\phi}(t_2) = At_2^2 + 33 = 22$$

If we can determine another equation involving  $A$  and  $t_2$  we can solve these two equations simultaneously.

The only other important piece of information regarding the motion is that  $q_2 = \pi$ . To make use of that information, we must "convert" our velocity equation into an angular velocity ( $\omega$ ) equation and then integrate the resulting equation into an equation for  $q$ . To determine the angular velocity function,

$$v_{\phi} = R\omega$$

$$\omega = \frac{v_{\phi}}{R}$$

$$\omega(t) = \frac{At^2 + 33}{95}$$

$$\omega(t) = \frac{A}{95}t^2 + 0.35$$

The angular position ( $\theta$ ) is the integral of the angular velocity,

$$\theta(t) = \int \omega(t) dt$$

$$\theta(t) = \int \left( \frac{A}{95}t^2 + 0.35 \right) dt$$

$$\theta(t) = \frac{A}{285}t^3 + 0.35t + D$$

Since we know  $\theta = 0$  when  $t = 0$  s, we can determine the integration constant:

$$\theta(0) = \frac{A}{285}(0)^3 + 0.35(0) + D = 0$$

$$D = 0$$

We also know that  $\theta = \pi$  at  $t_2$ , so:

$$\theta(t_2) = \frac{A}{285}t_2^3 + 0.35t_2 = \pi$$

These two equations,

$$\frac{A}{285}t_2^3 + 0.35t_2 = \pi$$

$$At_2^2 + 33 = 22$$

can be solved by substitution (or by using a solver). Solve the second equation for  $A$ , and substitute this expression into the first equation. You can solve the resulting equation for  $t_2$ . The solution is  $t_2 = 10.0$  s, the time for the car to complete the turn. Plugging this value back into the original equation allows you to determine  $A = -0.11$ .

Once you know  $A$ , you can complete the rest of the motion table. For example, since

$$\omega(t) = \frac{A}{95}t^2 + 0.35$$

$$\omega(t) = -0.00116t^2 + 0.35$$

We can now determine angular acceleration,

$$\alpha(t) = \frac{d\omega(t)}{dt}$$

$$\alpha(t) = \frac{d}{dt}(-0.00116t^2 + 0.35)$$

$$\alpha(t) = -0.00232t$$

radial acceleration,

$$a_r(t) = -R\omega(t)^2$$

$$a_r(t) = -95(-0.00116t^2 + 0.35)^2$$

and tangential acceleration.

$$a_t(t) = R\alpha(t)$$

$$a_t(t) = 95(-0.00232t)$$

$$a_t(t) = -0.22t$$

Each of these functions could be evaluated at  $t_1 = 0$  s and  $t_2 = 10.0$  s to complete the table.

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## 07. Activities

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## 04. Conservation

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  - [02. Applying the Impulse-Momentum Relation](#)
  - [03. Elastic Potential Energy](#)
  - [04. Applying Newton's 2nd Law to a Spring-Mass System](#)
  - [05. Potential Energy Functions](#)
  - [06. Activities](#)
- 

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## 01. Concepts and Principles

1. [The Impulse-Momentum Relation](#)
2. [The Work-Energy Relation](#)

### The Impulse-Momentum Relation

We've already used the impulse-momentum relation to analyze situations involving constant forces. The relation is typically applied in its component form:

$$\begin{aligned}mv_{x_i} + \Sigma(F_x(\Delta t)) &= mv_{x_f} \\mv_{y_i} + \Sigma(F_y(\Delta t)) &= mv_{y_f} \\mv_{z_i} + \Sigma(F_z(\Delta t)) &= mv_{z_f}\end{aligned}$$

Hopefully it's not too much of a stretch to argue that for forces that vary in magnitude or direction that simple summation over a time interval ( $\Delta t$ ) must be replaced by an integral over an infinitesimal time ( $dt$ ):

$$\begin{aligned}mv_{x_i} + \int_{t_i}^{t_f} F_x dt &= mv_{x_f} \\mv_{y_i} + \int_{t_i}^{t_f} F_y dt &= mv_{y_f} \\mv_{z_i} + \int_{t_i}^{t_f} F_z dt &= mv_{z_f}\end{aligned}$$

Regardless of whether that forces acting on an object are constant or not, the impulse they exert on the object is precisely equal to the change in the object's momentum.

### The Work-Energy Relation

Our previous encounter with the work-energy relation resulted in:

$$\frac{1}{2}mv_i^2 + mgh_i + \Sigma(|F||\Delta r|\cos\phi) = \frac{1}{2}mv_f^2 + mgh_f$$

where the forces acting on the object of interest were constant in both magnitude and direction. Again, for forces that vary, I will generalize this result to:

$$\frac{1}{2}mv_i^2 + mgh_i + \int_{r_i}^{r_f} (F \cos\phi) dr = \frac{1}{2}mv_f^2 + mgh_f$$

where the angle  $\phi$  is the angle between the instantaneous force acting on the object and the instantaneous displacement of the object ( $dr$ ). Thus, this angle can change as the object moves along its path. Technically, this integral is termed a *line integral* and its evaluation can be rather complicated.

Also recall from our previous discussion of work-energy that this is *not* a vector equation, meaning it is not applied independently in each of the coordinate directions.

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## 02. Applying the Impulse-Momentum Relation

1. Applying the Impulse-Momentum Relation
2. Applying the Work-Energy Relation
  1. The Force Exerted by a Spring

### Applying the Impulse-Momentum Relation

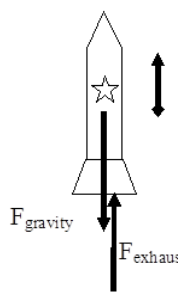
Let's re-examine the same situation we examined at the beginning of the previous chapter, a rocket launched directly upward with a time-dependent thrust.

*A 2.0 kg toy rocket is fitted with an engine that provides a thrust roughly modeled by the function  $F(t) = (60 \text{ N/s})t - (15 \text{ N/s}^2)t^2$ , for  $0 < t < 4.0 \text{ s}$ , and zero thereafter. The rocket is launched directly upward.*

For analysis, we'll apply the impulse-momentum relation between:

Event 1: The instant the rocket leaves the launch-pad

Event 2: The instant the thrust drops to zero.



$$mv_{y_i} + \int_{t_i}^{t_f} F_y dt = mv_{y_f}$$

$$2(0) + \int_{0.36}^{4.0} (F_{\text{thrust}} - F_{\text{gravity}}) dt = 2v_{y_f}$$

$$\int_{0.36}^4 (60t - 15t^2 - 2(9.8)) dt = 2v_{y_f}$$

Remember from last chapter that the rocket does not leave the launch pad until 0.36 s after the engine is ignited.

$$\int_{0.36}^4 (60t - 15t^2 - 19.6) dt = 2v_{y_f}$$

$$[30t^2 - 5t^3 - 19.6t]_{0.36}^4 = 2v_{y_f}$$

$$(30(4)^2 - 5(4)^3 - 19.6(4)) - (30(0.36)^2 - 5(0.36)^3 - 19.6(0.36)) = 2v_{y_f}$$

$$v_{y_f} = 42.5 \text{ m/s}$$

When the engine shuts off, the rocket is traveling at 42.5 m/s upward.

We could also apply the impulse-momentum relation between:

Event 1: The instant the thrust drops to zero.

Event 2: The instant the rocket reaches its maximum height.

During this interval, the only force acting on the rocket is the force of gravity, and the impulse-momentum relation is:

$$mv_{y_i} + \int_{t_i}^{t_f} F_y dt = mv_{y_f}$$

$$2(42.5) + \int_4^{8.34} (-F_{\text{gravity}}) dt = 2(0)$$

$$85 + \int_4^{8.34} (-19.6) dt = 0$$

$$85 - 19.6(t_f) + 19.6(4) = 0$$

$$t_f = 8.34 \text{ s}$$

Thus, the rocket reaches its highest altitude 8.34 s after launch.

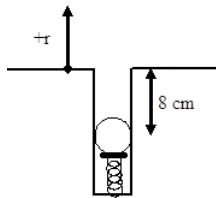
It's important to note that when a force is a function of time, it's relatively easy to integrate the function and determine the impulse. However, it should be clear that it would *not* be easy to determine the work done by a force of this type. Since work is expressed as

an integral of a force with respect to a displacement ( $dr$ ), the force function has to be expressed in terms of position,  $r$ . In general, it's not an easy (or sometimes possible) task to "convert" a function of time into a function of position, so work-energy is not a particularly useful way to analyze systems when the forces acting are time dependent. However, if the forces depend upon the *position* of the object, work-energy is a powerful analysis tool.

## Applying the Work-Energy Relation

*A 0.15 kg ball is launched vertically upward by means of a spring-loaded plunger, pulled back 8.0 cm and released. It requires a force of about 10 N to push the plunger back 8.0 cm.*

### The Force Exerted by a Spring



The force that the spring exerts on the ball depends on the amount by which the spring is compressed. The more the spring is compressed, the larger the force it exerts on the ball. A common *model* is that the force exerted by a spring is directly proportional to the amount of deformation of the spring, in this case compression. Deformation ( $s$ ) is defined to be the difference between the current length of the spring ( $L$ ) and the equilibrium length ( $L_0$ ):

$$s = L - L_0$$

If we define the positive coordinate direction to point in the same direction as positive deformation (stretch), then

$$F_{\text{spring}} = -ks$$

with the proportionality constant,  $k$ , referred to as the *spring constant*.

For the plunger, since it takes 10 N to compress the spring by 8.0 cm,

$$\begin{aligned} F_{\text{spring}} &= -ks \\ 10 &= -k(-0.08) \\ k &= 125 \frac{\text{N}}{\text{m}} \end{aligned}$$

For analysis, we'll apply the work-energy relation between:

Event 1: The instant the plunger is released.

Event 2: The instant the ball reaches its maximum height.

For these two events, work-energy looks like this:

$$\begin{aligned} \frac{1}{2}mv_i^2 + mgh_i + \int_{r_i}^{r_f} (F \cos \phi) dr &= \frac{1}{2}mv_f^2 + mgh_f \\ 0 + 0.15(9.8)(-0.08) + \int_{r_i}^{r_f} (F_{\text{spring}} \cos \phi) dr &= 0 + 0.15(9.8)h_f \end{aligned}$$

To do the integral, we must express  $F_{\text{spring}}$  in terms of the variable of integration,  $r$ . For the coordinate system chosen,  $s$  and  $r$  are identical. Also note that the force of the spring and the direction of motion of the ball point in the same direction. Thus,  $\phi = 0$ .

$$\begin{aligned} -0.118 + \int_{-0.08}^0 (-125s)(\cos 0) ds &= 1.47 h_f \\ -0.118 + \left[ -62.5s^2 \right]_{-0.08}^0 &= 1.47 h_f \\ -0.118 + \left[ 0 + 62.5(-0.08)^2 \right] &= 1.47 h_f \\ -0.118 + 0.40 &= 1.47 h_f \\ h_f &= 0.19 \text{ m} \end{aligned}$$

The ball reaches a maximum height of 19 cm above the top of the plunger.

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## 03. Elastic Potential Energy

1. Elastic Potential Energy
2. Applying the Work-Energy Relation with Elastic Potential Energy

### Elastic Potential Energy

When an object interacts with a spring, or other elastic material, a common model is that the material reacts linearly, i.e., with a force directly proportional to the deformation of the material. It is possible to calculate the work done by the linear material in general, and to rewrite the work-energy relation in such a way as to incorporate the effects of this work from the start. This is referred to as constructing a *potential energy function* for the work done by the elastic material. We did exactly the same thing earlier for the force of gravity. Imagine a spring of spring constant  $k$ , initially deformed by a distance  $s_i$ . It changes its deformation, ultimately resulting in deformation  $s_f$ . To calculate the work done by the spring on the object causing the deformation:

$$Work = \int_{r_i}^{r_f} F \cos \phi \, dr \quad (1)$$

$$Work = \int_{r_i}^{r_f} (-ks) \cos \phi \, dr \quad (2)$$

Choosing a coordinate system in which  $s$  and  $r$  are interchangeable (the origin is located at the point where the spring is at its natural length and the direction of elongation is positive) and allowing the force and the displacement to be in the same direction ( $\phi = 0$ ) results in,

$$Work = \int_{r_i}^{r_f} (-ks) \cos 0 \, dr \quad (3)$$

$$Work = -\frac{1}{2}ks_f^2 + \frac{1}{2}S_i^2 \quad (4)$$

The terms,  $\frac{1}{2}ks^2$ , are referred to as *elastic potential energy*. Inserting this result into the work-energy relation results in

$$\frac{1}{2}mv_i^2 + \frac{1}{2}ks_i^2 + mgh_i + \int_i^{r_f} (F \cos \phi) dr = \frac{1}{2}mv_f^2 + \frac{1}{2}ks_f^2 + mgh_f$$

with the understanding that the forces remaining in the equation, which may do work on the system, do not include the force of gravity or the force of the spring. The work done by the force of gravity and the force of the spring are already included in the relation via the inclusion of the potential energy terms.

### Applying the Work-Energy Relation with Elastic Potential Energy

*A 0.15 kg ball is launched vertically upward by means of a spring-loaded plunger, pulled back 8.0 cm and released. It requires a force of about 10 N to push the plunger back 8.0 cm.*

The spring constant of the plunger is known to be 125 N/m from above. Applying the work-energy relation with potential energy terms between the instant the plunger is released and the instant the ball reaches its maximum height results in

$$\begin{aligned} \frac{1}{2}mv_i^2 + \frac{1}{2}ks_i^2 + mgh_i + \int_i^{r_f} (F \cos \phi) dr &= \frac{1}{2}mv_f^2 + \frac{1}{2}ks_f^2 + mgh_f \\ 0 + \frac{1}{2}(125)(-0.08)^2 + 0.15(9.8)(-0.08) + 0 &= 0 + 0 + 0.15(9.8)h_f \\ h_f &= 0.19 \text{ m} \end{aligned}$$

Of course, this results in the same answer as before.

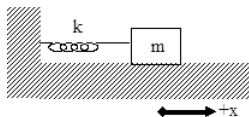
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## 04. Applying Newton's 2nd Law to a Spring-Mass System

### 1. Applying Newton's Second Law to a Spring-Mass System

### Applying Newton's Second Law to a Spring-Mass System

Since we did not analyze systems involving springs in the previous dynamics section, we should make up for that omission now.



The system above consists of a spring with spring constant  $k$  attached to a block of mass  $m$  resting on a frictionless surface. The origin of the coordinate system is located at the position in which the spring is unstretched.

Now imagine the block is pulled to the right and let go. Hopefully you can convince yourself that the block will oscillate back and forth. Let's apply *Newton's Second Law* at the instant the mass is at an arbitrary position,  $x$ . The only force acting on the mass in the  $x$ -direction is the force of the spring.

$$\sum \vec{F} = m\vec{a} \quad (1)$$

$$\vec{F}_{spring} = m\vec{a} \quad (2)$$

$$-ks = ma \quad (3)$$

Because of our choice of coordinate system, the stretch of the spring ( $s$ ) is exactly equal to the location of the block ( $x$ ). Therefore,

$$-kx = ma \quad (4)$$

Note that when the block is at a positive position, the force of the spring is in the negative direction and when the block is at a negative position, the force of the spring is in the positive direction. Thus, the force of the spring always acts to return the block to equilibrium.

Rearranging gives

$$\begin{aligned} -\frac{k}{m}x &= a \\ -\frac{k}{m}x &= \frac{d^2x}{dt^2} \end{aligned}$$

and defining a constant,  $\omega^2$ , as

$$\omega^2 = \frac{k}{m}$$

(Granted, it seems pretty silly to define  $k/m$  as the square of a constant, but just play along. You may also find it frustrating to learn that this "omega" is *not* an angular velocity. The block does not even have an angular velocity!)

yields,

$$-\omega^2 x = \frac{d^2x}{dt^2} \quad (5)$$

Therefore, the position function for the block must have a second time derivative equal to the product of  $(-\omega^2)$  and itself. The only functions whose second time derivative is equal to the product of a negative constant and itself are the sine and cosine functions. Therefore, a [solution to this differential equation](#)<sup>1</sup> can be written:

$$x(t) = A \cos(\omega t + \phi) \quad (6)$$

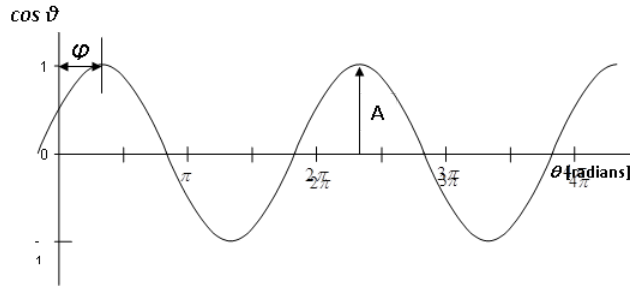
or equivalently with the sine function, where  $A$  and  $\phi$  are arbitrary constants.<sup>2</sup>

- $A$  is the *amplitude* of the oscillation. The amplitude is the maximum displacement of the object from equilibrium.
- $\phi$  is the *phase angle*. The phase angle is used to adjust the function forward or backward in time. For example, if the particle is at the origin at  $t = 0$  s,  $\phi$  must equal  $+\pi/2$  or  $-\pi/2$  to ensure that the cosine function evaluates to zero at  $t = 0$  s. If the particle



is at its maximum position at  $t = 0$  s, then the phase angle must be zero or  $\pi$  to ensure that the cosine function evaluates to +1 or -1 at  $t = 0$  s.

- $\omega$  is the *angular frequency* of the oscillation.<sup>3</sup>



Note that the cosine function repeats itself when its argument increases by  $2\pi$ . Thus, when

$$\Delta(\omega t + \phi) = 2\pi$$

the function repeats. Since  $\omega$  and  $\phi$  are constant,

$$\Delta(\omega t + \phi) = \omega \Delta t$$

Therefore, the time interval when

$$\omega \Delta t = 2\pi$$

is the time interval for one complete cycle of the oscillatory motion. The time for one complete cycle of the motion is termed the *period*,  $T$ . Thus,

$$T = \frac{2\pi}{\omega}$$

Therefore, the physical significance of the angular frequency is that it is inversely proportional to the period. Substituting in the definition of  $\omega$ :

$$\omega = \sqrt{\frac{k}{m}}$$

yields

$$T = 2\pi \sqrt{\frac{m}{k}}$$

In summary, a mass attached to a spring will oscillate about its equilibrium position with a position function given by:

$$x(t) = A \cos(\omega t + \phi)$$

This function repeats with a period of

$$T = 2\pi \sqrt{\frac{m}{k}}$$

1. A differential equation is an equation involving a function and its derivative(s).
2. To prove to yourself that this is indeed the solution to the equation, you should substitute the function,  $x(t)$ , into the left side of the equation and the second derivative of  $x(t)$  into the right side. This will verify that the two sides of the equation are equal. In addition to mathematically verifying this solution, you should verify the solution physically by sketching a graph of the motion that you *know* would result if the block were displaced to the right and comparing that sketch to a sketch of the function.
3. Again, note that  $\omega$  is *not* the angular velocity. The block is *not* rotating; it does not have an angular velocity.

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## 05.Potential Energy Functions

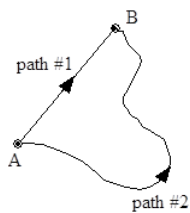
### 1. Potential Energy Functions

### Potential Energy Functions

The potential energy functions for the work done by gravity and by springs make analyzing many situations much easier. In light of this, why don't we construct potential energy functions for the work done by every different type of force that could possibly act on an object? Of course, one reason is that there are too many different types of forces. Having a potential energy function for the work done by every one of them would lead to so many potential energy functions that it would be hard to keep them all straight.

Another, more subtle reason is that it is impossible to construct potential energy functions for certain forces. A potential energy function must, by definition, be a *function*. Mathematically speaking, a potential energy *function* of position must assign a single, specific value of potential energy to every position. Functions must be single-valued. If you specify a height off of the ground, or the deformation of the spring, the potential energy function tells you exactly how much energy the system possesses at that position, regardless of the path the object took to reach that position.

In light of this observation, let's try to create a potential energy function to replace the work done by friction.



Imagine sliding an object along a rough surface from point A to point B. If you slide the object along path #1, the force of friction will do a certain amount of work on the object. (This work will be negative because the direction of the force of friction is always in opposition to the change in position of the object.)

If you slide the object along path #2, you should see that the magnitude of the work done by friction will be greater. (Although the frictional force will be the same in magnitude, the distance over which the frictional force acts will be larger. Thus, a larger amount of negative work will be done by friction.)

Now let's try to create a potential energy function for the work done by friction. Since we are always free to choose a coordinate system, we can choose a system in which the potential energy at A is zero. What is the value for the potential energy at B? Since the work done by friction depends on the path taken from A to B, so must the potential energy. However, this leaves us with a potential energy at B that can be either one of two values, depending on the path taken! Since a function must be single-valued, the work done by friction cannot be represented by a function. You can not create a potential energy function for the work done by friction!

There are also other forces whose work can not be represented by a potential energy function. In general, forces whose work can be represented by a potential energy function are termed *conservative* forces, while those for which potential energy functions can not be constructed are termed *non-conservative*.

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## 06. Activities

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## SECTION OVERVIEW

### 3.3: Dynamics

#### Topic hierarchy

01. Concepts and Principles

02. Time-Dependent Forces

03. Applying Newton's Second Law to Circular Motion

04. Activities

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## 01. Concepts and Principles

This project was supported, in part, by the National Science Foundation. Opinions expressed are those of the author and not necessarily those of the Foundation.

**Model Specifics** We are now ready to consider how the actual size and shape of an object affects its motion. The most immediate ramification of this Model is that the objects we study can now rotate. However, there are still three important restrictions to the types of motions we will investigate. The object is rigid. The size and shape of the object under investigation do not change during the motion. The rotation is planar. This restriction may take some explaining. Imagine an object spinning in place. The axis (real or imaginary) about which the object spins is referred to as the rotation axis. Note that all points that lie on the rotation axis do not move. All other points on the object, however, exhibit circular motion around the rotation axis. If the object, in addition to spinning, is also moving, this model will restrict us to motions where the center-of-mass moves exclusively in the plane perpendicular to the rotation axis. I will refer to this as planar rotation. For example, a yo-yo is an object that typically moves in a plane perpendicular to its rotation axis. Thus, we can study the motion of most yo-yos. The motion of a wheel is typically in a plane perpendicular to its rotation axis. Thus, we can study the motion of most wheels. The motion of the earth around the sun, however, is beyond this model's capabilities since the rotation axis of the earth is tilted relative to the plane of the earth's motion. A wobbling top is also beyond this model's capabilities (although a steadily spinning top is not). The object is classical.

**Kinematics Concepts and Principles**

**Rotation about a Fixed Axis (Spinning)** Imagine a rigid body constrained to rotate about a fixed axis. Non-physicists would say that the object is spinning. Place the origin of a coordinate system at the location of the rotation axis. Examine an arbitrary point on the object, denoted by the position  $\vec{r}$ . Notice that as the object spins, this point undergoes circular motion (denoted by the dashed line). Although the actual object may be of irregular shape, as it spins every point on the object undergoes circular motion. Therefore, to understand the kinematics of spinning objects we first must understand the kinematics of circular motion, so let's examine circular motion in detail.

**Defining the Angular Kinematic Variables** The position of an object moving along a circular path of radius  $R$  (whether part of a spinning object or a separate point particle) can always be described by the following relations, and assuming the origin of the coordinate system is placed at the center of the circle.  $\theta$  specifies the angular position of the object, and is typically measured in radians. It specifies where along the circle the object is located. For example, if  $\theta$  is a constant, the object's position along the circle is constant (the object doesn't move). If  $\theta$  varies, the object's position along the circle varies. The rate at which  $\theta$  is changing,  $\omega$ , is termed the angular velocity of the object and denoted by  $\omega$ , the Greek letter omega. Thus, Since  $\omega$  is the rate at which the angular position is changing, it will have units of radians/second, or rad/s. The rate at which  $\omega$  is changing,  $\alpha$ , is termed the angular acceleration of the object and denoted by  $\alpha$ , the Greek letter alpha. Thus, Since  $\alpha$  is the rate at which the angular velocity is changing, it will have units of rad/s<sup>2</sup>. Notice that since angular position, angular velocity, and angular acceleration are defined with exactly the same inter-relationships as "regular" position, velocity, and acceleration, the kinematic equations involving  $\theta$ ,  $\omega$ , and  $\alpha$  must follow exactly the same form as the kinematic equations involving  $r$ ,  $v$ , and  $a$ . In other words, and must be valid relationships between the angular variables. Deriving the Relationships between Angular and Linear Kinematics Now that we have the inter-relationships between the angular quantities out of the way, let's determine the relationships between the angular variables and the linear variables. At left is illustrated the position of the object ( $r_1$  and  $r_2$ ) at two instants separated by  $\Delta t$ , as well as the change in the position ( $\Delta r$ ) between these times. The length of  $\Delta r$  is approximately equal to the arc length between  $r_1$  and  $r_2$ , especially when  $\Delta \theta$  is very small. From geometry, the arc length is equal to the product of the radius of the circle and the enclosed angle. Thus, Since the change in position takes place over a time interval  $\Delta t$ , dividing both sides by  $\Delta t$  yields  $v = R\omega$ . Thus, the magnitude of the velocity of an object undergoing circular motion is the product of the radius and the angular velocity. Also note that the direction of the velocity is tangent to the circular path of the object. We will refer to this direction as the ' $\hat{\theta}$ ' direction, since it is in the direction of increasing angular position,  $\theta$ . In short, Now let's determine a relationship for the acceleration of the particle. At left is illustrated the velocity of the object ( $v_1$  and  $v_2$ ) at two instants separated by  $\Delta t$ . Examining the change in the velocity allows us to construct the acceleration. Notice that the acceleration can be thought of as having two components: one directed radially inward toward the center of the circle ( $a_R$ ) and one directed tangent to the circular path ( $a_\theta$ ). The ' $R$ ', or radial, direction will refer to the direction of increasing radius, i.e., away from the center of the circle. Thus  $a_R$  will turn out to be negative since it points toward the center of the circle. Let's look at the radial component of acceleration first. The radial component of the acceleration is due to the changing direction of the velocity. As the velocity changes direction, the arrow representing the velocity rotates. If the speed of the object is constant, the tip of this arrow traces a circular path. The magnitude of  $\Delta v$  is approximately equal to the arc length between  $v_1$  and  $v_2$  (ignoring the change in length of  $v_2$ ), especially when  $\Delta \theta$  is very small. From geometry, the arc length is equal to the product of the "radius" of the circle,  $v$ , and the enclosed angle. Thus, Since the change in velocity takes place over a time interval  $\Delta t$ , dividing both sides by  $\Delta t$  yields  $a_R = -v\omega$ . Substituting, results in: As noted above, this acceleration points toward the center of the circular path, in the direction of decreasing radius, so our final result for the radial acceleration is:  $a_R = -v\omega$ . Now let's examine the tangential component of the acceleration. The

tangential component of the acceleration is due to the changing magnitude of the velocity. Ignoring the change in direction of the velocity (since we've taken care of that in computing  $a_R$ ), Since the change in velocity takes place over a time interval  $\Delta t$ , dividing both sides by  $\Delta t$  yields  $a_t = \frac{\Delta v}{\Delta t}$ . Thus, the acceleration of an object undergoing circular motion has two components. If the object is speeding up or slowing down,  $a_t$  does not equal zero and there is an acceleration component in the tangential direction. However, even if the object is moving at constant speed there is an acceleration component in the radial direction, pointing toward the center of the circle. By virtue of traveling in a circle, the velocity of an object continually changes its orientation and this change in orientation results in an acceleration toward the center of the circle. Everything described above is true for both a point on a spinning rigid body as well as for a particle undergoing circular motion. However, for a rigid body, since each and every point on the object has to complete an entire cycle around the rotation axis in the same amount of time, every point must undergo circular motion with the same angular speed,  $\omega$ , and the same angular acceleration,  $\alpha$ . Since every point on a rigid body must have the same angular speed and the same angular acceleration, we will speak of the angular speed and angular acceleration of the object, rather than of each point on the object.

**Rotation and Translation** To describe the motion of an object undergoing pure rotation (spinning), we have to describe the circular motion each point on the object undergoes. What must we do if the object is simultaneously rotating and translating (moving in a plane perpendicular to the rotation axis), like a wheel rolling down an incline? The answer lies in the independence of these two types of motion. In much the same manner as we attacked kinematics in two dimensions by independently analyzing the horizontal and vertical motions, we will attack rigid-body kinematics by independently analyzing the rotational and translational motions. In short, we will model any motion of a rigid-body as a superposition of a translation of the object's center-of-mass (CM) (which we will analyze by particle kinematics) and a rotation about an axis passing through the CM (which we will analyze by the kinematics of spinning, detailed above). For example, examine the motion of the thin rod between  $t_1$  and  $t_2$ . Although the rod may have been spinning crazily through space between these two times, we can model its motion as a superposition of a simple translation of its CM without rotation (denoted by the vector  $\mathbf{r}$ ), which leaves the rod in the orientation denoted by the dashed lines, and a simple rotation about an axis through its CM without translation (denoted by  $\theta$ ), which leaves the rod in its proper, final orientation. If we imagine the time difference ( $t_2 - t_1$ ) shrinking toward zero, hopefully it becomes plausible that we can model any motion through this method. In summary, to describe the motion of an arbitrary rigid body we will break the motion down into a pure translation of the CM and a pure rotation about the CM. We will use particle kinematics to describe the translational portion of the motion and the kinematics of circular motion to describe the rotational portion. The velocity (or acceleration) of any point on the object is then determined by the sum of the velocity (or acceleration) due to the translation and the velocity (or acceleration) due to the rotation.

**Kinematics Analysis Tools Pure Rotation** After the off button is pressed, a ceiling fan takes 22 s to come to rest. During this time, it completes 18 complete revolutions. To analyze this situation, we should first carefully determine and define the sequence of events that take place. At each of these instants, let's tabulate what we know about the motion. Since we are dealing with pure rotation, the relevant kinematic variables are the angular position, velocity, and acceleration of the fan. Also, let's take the direction that the fan is initially rotating to be the positive direction. Event 1: The 'off' button is pressed  $t_1 = 0 \text{ s}$   $\theta_1 = 0 \text{ rad}$   $\omega_1 = 12 \text{ rad/s}$   $\alpha_1 = 0 \text{ rad/s}^2$   
 = Event 2: The ceiling fan stops  $t_2 = 22 \text{ s}$   $\theta_2 = 18(2\pi) = 113 \text{ rad}$   $\omega_2 = 0 \text{ rad/s}$   $\alpha_2 = 0 \text{ rad/s}^2$  Using the angular versions of the kinematic equations yields: Now substitute this expression into the other equation: Substitute this result back into the original equation: Notice that the sign of the angular acceleration is negative. This indicates that the angular acceleration is in the opposite direction of the angular velocity, as it should be since the fan is slowing down. Also note that every point on the ceiling fan has the angular velocity and angular acceleration calculated above.

**Circular Motion** An automobile enters a U-turn of constant radius of curvature 95 m. The car enters the U-turn traveling at 33 m/s north and exits at 22 m/s south. Since the car moves along a circular path, we will not break the position, velocity and acceleration vectors into x- and y-components. Rather, we will use polar coordinates, where we view these vectors as having either radial (R) or tangential (t) components. The position is solely in the radial direction (which is written as  $rR$ , read as "position in the radial direction", or simply the radius of the circular path), the velocity is solely in the tangential direction ( $v_t$ ), and the acceleration has components in both directions ( $a_R$  and  $a_t$ ). Additionally, we can tabulate the angular position, angular velocity, and angular acceleration of the car.

Event 1: The auto enters the turn.  $t_1 = 0 \text{ s}$   $rR = 95 \text{ m}$   $\theta_1 = 0 \text{ rad}$   $v_1 = 33 \text{ m/s}$   $\omega_1 = 12 \text{ rad/s}$   $a_R = 0 \text{ m/s}^2$   $a_t = 0 \text{ m/s}^2$   
 = Event 2: The auto exits the turn.  $t_2 = 22 \text{ s}$   $rR = 95 \text{ m}$   $\theta_2 = 113 \text{ rad}$   $v_2 = 22 \text{ m/s}$   $\omega_2 = 0 \text{ rad/s}$   $a_R = 0 \text{ m/s}^2$   $a_t = 0 \text{ m/s}^2$  Notice that since the car completely reverses its direction of travel, it must have traveled halfway around a circular path. Thus,  $\theta_2 - \theta_1 = \pi \text{ rad}$ .

The unknown variables listed above are not all independent. In fact, we know we can go back and forth between position, velocity, and acceleration in polar coordinates and angular position, angular velocity, and angular acceleration using the relationships we derived above. Using these relationships, we can quickly determine the angular velocity of the car at the two events: and the radial acceleration of the car: All that's left to determine is the time it takes for the car to complete the u-turn and the tangential

acceleration and angular acceleration it has as it performs this turn. Since there are only two variables in the set of angular quantities, we can use the two kinematic equations, written in terms of angular quantities, to finish our analysis.

Now substitute this expression into the other equation: Substitute this result back into the original equation: The tangential acceleration can be determined from: Thus, the car has acceleration in both the radial and tangential directions as it goes around the u-turn. The tangential acceleration corresponds to the rate at which the speed of the car decreases, and is due to the driver stepping on the brakes as she goes around the turn. The radial acceleration is due to the change in direction of the car as it goes around the turn, and is ultimately due to the frictional forces acting on the tires of the car which prevent it from sliding off the road. When the car initially enters the turn, this acceleration is quite large and the required frictional force to stay on the road is large. As the car slows, this acceleration decreases as the required frictional force to stay on the road decreases.

Connecting Pure Rotation to Pure Translation The device at right is used to lift a heavy load. The free rope is attached to a truck which accelerates from rest at a rate of  $1.5 \text{ m/s}^2$ . The inner radius of the pulley is  $20 \text{ cm}$  and the outer radius is  $40 \text{ cm}$ . The load must be raised  $15 \text{ m}$ . The coordinate system chosen indicates that the block moving upward, the pulley rotating clockwise, and the truck moving to the right are all positive. There are three different objects that we should be able to describe kinematically; the truck, the pulley, and the block. Let's tabulate everything we know about each object:

Object	Event	Time ( $t$ )	Position ( $r$ )	Velocity ( $v$ )	Acceleration ( $a$ )
Truck	Event 1: Truck begins to move	$t_1 = 0 \text{ s}$	$r_1 = 0 \text{ m}$	$v_1 = 0 \text{ m/s}$	$a_1 = 1.5 \text{ m/s}^2$
	Event 2: Block raised $15 \text{ m}$	$t_2 = 2 \text{ s}$	$r_2 = 30 \text{ m}$	$v_2 = 6 \text{ m/s}$	$a_2 = 1.5 \text{ m/s}^2$
Pulley	Event 1: Truck begins to move	$t_1 = 0 \text{ s}$	$\theta_1 = 0 \text{ rad}$	$\omega_1 = 0 \text{ rad/s}$	$\alpha_1 = 3.75 \text{ rad/s}^2$
	Event 2: Block raised $15 \text{ m}$	$t_2 = 2 \text{ s}$	$\theta_2 = 75 \text{ rad}$	$\omega_2 = 23.7 \text{ rad/s}$	$\alpha_2 = 3.75 \text{ rad/s}^2$
Block	Event 1: Truck begins to move	$t_1 = 0 \text{ s}$	$r_1 = 0 \text{ m}$	$v_1 = 0 \text{ m/s}$	$a_1 = 0.75 \text{ m/s}^2$
	Event 2: Block raised $15 \text{ m}$	$t_2 = 2 \text{ s}$	$r_2 = 15 \text{ m}$	$v_2 = 3 \text{ m/s}$	$a_2 = 0.75 \text{ m/s}^2$

The truck and block exhibit translational motion, so we only have to tabulate linear variables. The pulley spins, so the relevant variables are the angular variables. Notice that just because the block moves  $15 \text{ m}$  doesn't mean the truck moves  $15 \text{ m}$ . Also, the acceleration of the block is not equal to the acceleration of the truck. However, these variables are related to each other since both objects are attached to the same pulley. Also notice that we can't currently solve this problem. Each of the objects has three unknown quantities. However, since the kinematics of the three objects are related, we will be able to solve the problem once we've worked out the exact relationship between each object's kinematics. Let's start with the acceleration. Assuming the rope from the truck does not slip on the pulley, the point on the pulley in contact with the rope must be accelerating at the same rate as the truck. Notice that this acceleration is tangent to the pulley, and this rope is located  $0.4 \text{ m}$  from the center of the pulley. Therefore, from The pulley must have an angular acceleration of  $3.75 \text{ rad/s}^2$  since it is attached to the truck. In addition, assuming the rope from the block does not slip on the pulley, the point on the pulley in contact with this rope must be accelerating at: Since the point on the pulley attached to the block is accelerating at  $0.75 \text{ m/s}^2$ , the block itself must be accelerating at  $0.75 \text{ m/s}^2$ . In a nutshell, what we've done is used the acceleration of the truck to find the angular acceleration of the pulley, and then used the angular acceleration of the pulley to find the acceleration of the block. This chain of reasoning is very common. Now that we know the block's acceleration, we can solve the problem. Applying the two kinematic equations to the block yields: Now that we know the final speed of the block, we can find the final angular speed of the pulley and the final speed of the truck: Again assuming the rope from the block does not slip on the pulley, the point on the pulley in contact with the rope must be moving at the same rate as the block. Notice that this velocity is tangent to the pulley, and this rope is located  $0.2 \text{ m}$  from the center of the pulley. Therefore, from The pulley must have an angular velocity of  $23.7 \text{ rad/s}$  since it is attached to the block. In addition, assuming the rope from the truck does not slip on the pulley, the point on the pulley in contact with this rope must be moving at: The truck is moving at  $9.48 \text{ m/s}$  when the block reaches  $15 \text{ m}$ . We can relate the displacement of the block to the angular displacement of the pulley and the displacement of the truck using where  $r$  can be thought of as either the displacement of the object attached to the rope or the arc length over which the rope is wrapped around the pulley (since the amount of rope wrapped around a pulley is exactly equal to the displacement of the object attached to the rope), and  $\theta$  is the angular displacement of the pulley. Therefore, relating the block to the pulley yields The pulley must have turned through  $75 \text{ rad}$  since it is attached to the block. Relating the pulley to the truck results in: The truck moved  $30 \text{ m}$  while the block moved  $15 \text{ m}$ . Notice that in all cases, the values of the kinematic variables for the truck are exactly twice the values for the block. This is not a coincidence. Since the truck is attached to the pulley at twice the distance from the axle that the block is attached, all of its kinematic variables will have twice the value. Once you understand why this is the case, you can use this insight to simplify your analysis. Rotating and Translating Accelerating from rest, a Cadillac Sedan de Ville can reach a speed of  $25 \text{ m/s}$  in a time of  $6.2 \text{ s}$ . During this acceleration, the Cadillac's tires do not slip in their contact with the road. The diameter of the Cadillac's tires is  $0.80 \text{ m}$ . Let's examine the motion of one of the Cadillac's tires. Obviously, the tire both translates and rotates. We will imagine the motion to be a superposition of a pure translation of the CM of the tire and a pure rotation about the CM of the tire. To analyze the situation, we should define the sequence of events that take place and tabulate what we know about the motion at each event. Since we are dealing with translation as well as rotation, we will need to keep track of both the linear kinematic variables and the angular kinematic variables. Let's take the positive  $x$  direction to be the direction that the Cadillac translates and the positive  $\theta$  direction to be the direction in which the tires rotate. Event 1: The Cadillac begins to accelerate  $t_1 = 0 \text{ s}$   $r_1 = 0 \text{ m}$   $\theta_1 = 0 \text{ rad}$   $v_1 = 0 \text{ m/s}$   $\omega_1 = 0 \text{ rad/s}$

rad/s  $a_{12} =$

12 = Event 2: The Cadillac reaches 25 m/s  $t_2 = 6.2$  s  $r_2 =$  2 =

$v_2 = 25$  m/s 2 = First, let's examine the translational portion of the tire's motion. What about the rotational kinematics of the tire? The motion of the tire is the superposition of the pure translational motion of the CM and the pure rotational motion about the CM. Thus, the velocity of any point on the tire due to rotation is given by where  $R$  is the distance from the rotation axis, i.e., the distance from the CM. The key insight into studying the rotation of the tire is to realize that at any instant the velocity of the point on the tire in contact with the road is zero because the tire never slips in its contact with the road. Thus, if the CM of the tire is moving forward at 25 m/s, the point on the bottom of the tire must be moving backward relative to the CM at the exact same speed in order for its velocity relative to the ground to be zero! Thus, only a very particular value for will allow the tire to roll without slipping. From above, Therefore, at the end of the acceleration the tire has an angular speed of 62.5 rad/s. We can find the angular acceleration and angular position of the tire by using the linear variables determined above and "translating" those values into angular variables, or by using the angular versions of the kinematic equations. Thus, the wheel rotates through 194 rad, or 30.8 revolutions, while accelerating. Kinematics Activities Construct motion diagrams for the motions described below. a. A satellite has been programmed to circle a stationary space station at a radius of 10 km and a constant angular speed of 0.02 rad/s. b. A rider on a merry-go-round, 3 m from the axis, is traveling at 4 m/s. The merry-go-round slows, and the rider reaches a speed of 0.5 m/s in 11 seconds. c. In a device built to acclimate astronauts to large accelerations, astronauts are strapped into a pod that is swung in a 6 m radius circle at high speed. The linear speed of the pod is increased from rest to a speed of 17 m/s in a time interval of 25 seconds. A car drives around a semi-circular U-turn at constant speed. Set the origin of the coordinate system at the center-of-curvature of the U-turn. a. Using polar coordinates, sketch position, velocity, and acceleration vs. time graphs for this motion. b. Sketch angular position, angular velocity, and angular acceleration vs. time graphs for this motion. A car drives around a semi-circular U-turn at decreasing speed. Set the origin of the coordinate system at the center-of-curvature of the U-turn. a. Using polar coordinates, sketch position, velocity, and acceleration vs. time graphs for this motion. b. Sketch angular position, angular velocity, and angular acceleration vs. time graphs for this motion. A girl pedals her bicycle at constant speed along a level path. Examine the motion of a point on the rim of her tire and a point midway along a spoke for one complete cycle. Set the origin of the coordinate system at the center of the tire. a. Using polar coordinates, sketch position, velocity, and acceleration vs. time graphs for each point. b. Sketch angular position, angular velocity, and angular acceleration vs. time graphs for each point. A girl pedals her bicycle, from rest, at a constantly increasing speed along a level path. Examine the motion of a point on the rim of her tire and a point midway along a spoke for one complete cycle, starting from rest. Set the origin of the coordinate system at the center of the tire. a. Using polar coordinates, sketch position, velocity, and acceleration vs. time graphs for each point. b. Sketch angular position, angular velocity, and angular acceleration vs. time graphs for each point. A block is attached to a rope that is wound around a pulley. Sketch velocity vs. time and acceleration vs. time graphs for a point on the rim of the pulley and for the block. Set the origin of the coordinate system at the center of the pulley. a. Sketch the polar components of the point on the rim's motion and the Cartesian components of the block's motion when the block is lowered at constant speed. b. Sketch the polar components of the point on the rim's motion and the Cartesian components of the block's motion when the block is released from rest. The two wheels at right are coupled together by an elastic conveyor belt that does not slip on either wheel. Examine the motion of a point on the rim of the large wheel and a point on the rim of the small wheel when the conveyor belt's speed increases from rest. Set the origin of a coordinate system at the center of each pulley. a. Using polar components, sketch velocity and acceleration vs. time graphs for each point. b. Sketch angular velocity and angular acceleration vs. time graphs for each point. The two gears at right have equal teeth spacing and are well-meshed. Examine the motion of a point on the rim of the large gear and a point on the rim of the small gear when the small gear's speed increases from rest in the counterclockwise direction. Set the origin of a coordinate system at the center of each gear and let counterclockwise be the positive direction. a. Using polar components, sketch velocity and acceleration vs. time graphs for each point. b. Sketch angular velocity and angular acceleration vs. time graphs for each point. Below are angular position vs. time graphs for six different objects. a. Rank these graphs on the basis of the angular velocity of the object. Rank positive angular velocities as larger than negative angular velocities. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank these graphs on the basis of the angular acceleration of the object. Rank positive angular accelerations as larger than negative angular accelerations. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Below are angular velocity vs. time graphs for six different objects. a. Rank these graphs on the basis of the angular acceleration of the object. Rank positive angular accelerations as larger than negative angular accelerations. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank these graphs on the basis of the angular displacement of the object over the time interval shown. Rank positive angular displacements as larger than negative angular displacements. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_



Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: An artificial satellite circles a space station at constant speed. The satellite passes through the six labeled points. For all questions below, use the indicated coordinate system. a. Rank the x-velocity of the satellite at each of the labeled points. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided.

b. Rank the y-velocity of the satellite at each of the labeled points. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. c. Rank the tangential velocity of the satellite at each of the labeled points. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. d. Rank the angular velocity of the satellite at each of the labeled points. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided.

Six artificial satellites of identical mass circle a space station with constant period  $T$ . The satellites are located a distance  $R$  from the space station.  $R$   $T$  A 5000 m 160 hrs B 2500 m 40 hrs C 2500 m 80 hrs D 10000 m 160 hrs E 5000 m 120 hrs F 10000 m 80 hrs a. Rank these satellites on the basis of their angular velocity. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. b. Rank these satellites on the basis of their tangential velocity. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your rankings: Six artificial satellites of identical mass circle a space station at constant speed  $v$ . The satellites are located a distance  $R$  from the space station.  $R$   $v$  A 5000 m 160 m/s B 2500 m 40 m/s C 2500 m 80 m/s D 10000 m 160 m/s E 5000 m 120 m/s F 10000 m 80 m/s a. Rank these satellites on the basis of their angular acceleration. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. b. Rank these satellites on the basis of their radial acceleration. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. c. Rank these satellites on the basis of their tangential acceleration. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your rankings: The pulley below has the outer radius and inner radius indicated. In all cases, the horizontal rope is pulled to the right at the same, constant speed. Router Rinner A 0.4 m 0.2 m B 0.4 m 0.3 m C 0.8 m 0.4 m D 0.6 m 0.5 m E 0.2 m 0.1 m F 0.6 m 0.2 m a. Rank these scenarios on the basis of the angular speed of the pulley. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank these scenarios on the basis of the speed of the block. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: The two wheels, with radii indicated below, are linked together by an elastic conveyor belt that does not slip on either wheel. In all cases, the small wheel is turning at the same, constant angular speed. Rlarge Rsmall A 0.4 m 0.2 m B 0.4 m 0.3 m C 0.8 m 0.4 m D 0.6 m 0.5 m E 0.2 m 0.1 m F 0.6 m 0.2 m a. Rank the scenarios on the basis of the speed of the belt. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank the scenarios on the basis of the angular speed of the large wheel. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: The two gears, with radii indicated below, have equal teeth spacing and are well-meshed. In all cases, the large gear is turning at the same, constant angular speed. Rlarge Rsmall A 0.4 m 0.2 m B 0.4 m 0.3 m C 0.8 m 0.4 m D 0.6 m 0.5 m E 0.2 m 0.1 m F 0.6 m 0.2 m a. Rank the scenarios on the basis of the speed of the teeth on the small gear. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank the scenarios on the basis of the angular speed of the small gear. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A pulley system is illustrated below. The horizontal rope is pulled to the right at constant speed. Each letter designates a point on either the pulley or on one of the two ropes. Neither rope slips in its contact with the pulley. a. Rank the designated points on the basis of their speed. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank the designated points on the basis of the magnitude of their acceleration. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: At the instant shown, the pulley below has the outer radius, inner radius, angular velocity, and angular acceleration indicated. Positive angular quantities are counterclockwise. Router Rinner A 0.4 m 0.2 m 10 rad/s 0 rad/s<sup>2</sup> B 0.4 m 0.3 m 10 rad/s 1 rad/s<sup>2</sup> C 0.8 m 0.4 m 10 rad/s 0 rad/s<sup>2</sup> D 0.6 m 0.5 m 5 rad/s -1 rad/s<sup>2</sup> E 0.2 m 0.1 m 20 rad/s -4 rad/s<sup>2</sup> F 0.6 m 0.2 m 15 rad/s 2 rad/s<sup>2</sup> a. Rank these scenarios on the basis of the speed of the block. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank these scenarios on the basis of the magnitude of the acceleration of the block. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_

\_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: At the instant shown, the pulley below has the outer radius, inner radius, angular velocity, and angular acceleration indicated. Positive angular quantities are counterclockwise. Router Rinner A 0.4 m 0.2 m 10 rad/s 0 rad/s<sup>2</sup> B 0.4 m 0.3 m 10 rad/s 0 rad/s<sup>2</sup> C 0.8 m 0.4 m 10 rad/s 0 rad/s<sup>2</sup> D 0.6 m 0.5 m 5 rad/s 0 rad/s<sup>2</sup> E 0.2 m 0.1 m 20 rad/s 0 rad/s<sup>2</sup> F 0.6 m 0.2 m 15 rad/s 0 rad/s<sup>2</sup> a. Rank these scenarios on the basis of the magnitude of the acceleration of a point on the inner rim of the pulley. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank these scenarios on the basis of the magnitude of the acceleration of the block. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: After being turned on, a record player (an antique device used to listen to music) reaches its rated angular speed of 45 rpm (4.71 rad/s) in 1.5 s. Motion Information Event 1:  $t_1 = 1 = 1 =$  12 = Event 2:  $t_2 = 2 = 2 =$  Mathematical Analysis A clothes dryer spins clothes at an angular speed of 6.8 rad/s. Starting from rest, the drier reaches its operating speed with an average angular acceleration of 7.0 rad/s<sup>2</sup>. Motion Information Event 1:  $t_1 = 1 = 1 =$  12 = Event 2:  $t_2 = 2 = 2 =$  Mathematical Analysis In a secret Las Vegas research laboratory, a roulette wheels is undergoing extensive testing. The wheel is spun at 25 rad/s. After 6.4 s, the wheel has rotated through 100 rad. Motion Information Event 1:  $t_1 = 1 = 1 =$  12 = Event 2:  $t_2 = 2 = 2 =$  Mathematical Analysis A rider on a merry-go-round, 3.0 m from the axis, is traveling at 4.0 m/s. The merry-go-round slows, and the rider reaches a speed of 0.50 m/s in 11 seconds. Motion Information Mathematical Analysis Event 1:  $t_1 = r_1R = 1 = v_1 = 1 = a_1R = a_{12} =$  12 = Event 2:  $t_2 = r_2R = 2 = v_2 = 2 = a_2R =$  A rider on a merry-go-round, 3.0 m from the axis, is traveling at 4.0 m/s. The merry-go-round slows to rest over three complete revolutions. Motion Information Mathematical Analysis Event 1:  $t_1 = r_1R = 1 = v_1 = 1 = a_1R = a_{12} =$  12 = Event 2:  $t_2 = r_2R = 2 = v_2 = 2 = a_2R =$  A rider on a 15 m diameter Ferris Wheel is initially at rest. The angular speed of the Ferris Wheel is increased to 1.5 rad/s over a time interval of 3.5 s. Motion Information Mathematical Analysis Event 1:  $t_1 = r_1R = 1 = v_1 = 1 = a_1R = a_{12} =$  12 = Event 2:  $t_2 = r_2R = 2 = v_2 = 2 = a_2R =$  An automobile enters a constant 90 m radius of curvature turn traveling at 25 m/s north and exits the curve traveling at 35 m/s east. Motion Information Mathematical Analysis Event 1:  $t_1 = r_1R = 1 = v_1 = 1 = a_1R = a_{12} =$  12 = Event 2:  $t_2 = r_2R = 2 = v_2 = 2 = a_2R =$  An automobile enters a constant 90 m radius of curvature turn traveling at 25 m/s north and exits the curve traveling east. The car completes the turn in 5.4 s. Motion Information Mathematical Analysis Event 1:  $t_1 = r_1R = 1 = v_1 = 1 = a_1R = a_{12} =$  12 = Event 2:  $t_2 = r_2R = 2 = v_2 = 2 = a_2R =$  In a device built to acclimate astronauts to large accelerations, astronauts are strapped into a pod that is swung in a 6.0 m radius circle at high speed. The angular speed of the pod is increased from rest to an angular speed of 2.8 rad/s in a time interval of 20 seconds. The device is then linearly slowed to rest over a time interval of 40 seconds. Motion Information Mathematical Analysis Event 1:  $t_1 = r_1R = 1 = v_1 = 1 = a_1R = a_{12} =$  12 = Event 2:  $t_2 = r_2R = 2 = v_2 = 2 = a_2R = a_{23} =$  23 = Event 3:  $t_3 = r_3R = 3 = v_3 = 3 = a_3R =$  In a device built to acclimate astronauts to large accelerations, astronauts are strapped into a pod that is swung in a 6.0 m radius circle at high speed. The linear speed of the pod is increased from rest to a speed of 17 m/s after three complete revolutions of the pod. The device is then linearly slowed to rest over a time interval of 35 seconds. Motion Information Mathematical Analysis Event 1:  $t_1 = r_1R = 1 = v_1 = 1 = a_1R = a_{12} =$  12 = Event 2:  $t_2 = r_2R = 2 = v_2 = 2 = a_2R = a_{23} =$  23 = Event 3:  $t_3 = r_3R = 3 = v_3 = 3 = a_3R =$  A diver rotating at approximately constant angular velocity completes four revolutions before hitting the water. She jumped vertically upward with an initial velocity of 5 m/s from a diving board 4 m above the water. Motion Information Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_{12} =$  12 = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 =$  Mathematical Analysis A baton twirler throws a spinning baton directly upward. As it goes up and returns to the twirler's hand, it turns through four complete revolutions. The constant angular speed of the baton while in the air is 10 rad/s. Motion Information Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_{12} =$  12 = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 =$  Mathematical Analysis A quarterback throws a pass that is a perfect spiral, spinning at 50 rad/s. The ball is thrown at 19 m/s at an angle of 35° above the horizontal. The ball leaves the quarterback's hand 2.0 m above the Astroturf and is caught just before it hits the turf. (The opposing coach thinks the ball was trapped. We are still waiting for the result of the challenge.) Motion Information Event 1:  $t_1 = r_1x = r_1y = 1 = v_1x = v_1y = 1 = a_{12}x = a_{12}y =$  12 = Event 2:  $t_2 = r_2x = r_2y = 2 = v_2x = v_2y = 2 =$  Mathematical Analysis A car, with 0.75 m diameter tires, slows from 35 m/s to 15 m/s over a distance of 70 m. The car's tires do not slip in their

contact with the road. Motion Information Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_{12} =$  12 = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 =$  Mathematical Analysis A bowling ball of diameter 21.6 cm is rolled down the alley at 4.7 m/s. The ball slows with an acceleration of 0.2 m/s<sup>2</sup> until it strikes the pins. The pins are located 18.3 m from the release point of the ball. Motion Information Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_{12} =$  12 = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 =$  Mathematical Analysis The man at right is trapped inside a section of large pipe. If that's not bad enough, the pipe begins to roll from rest down a 35 m long, 180 incline! The pipe has a diameter of 1.2 m. The pipe (and very dizzy man) reach the bottom of the incline after 6.32 s. Motion Information Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_{12} =$  12 = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 =$  Mathematical Analysis The unlucky man is falling at 20 m/s, 75 m above the crocodile-infested waters below! In an attempt to save him, the brake shoe is pressed against the spinning pulley. The action of the brake shoe gives the pulley an angular acceleration of 7.5 rad/s<sup>2</sup>. The man is saved! (Barely.) Motion Information Lucky Man Event 1:  $t_1 = r_1 = v_1 = a_{12} =$  Event 2:  $t_2 = r_2 = v_2 =$  Pulley Event 1:  $t_1 = 1 = 1 =$  12 = Event 2:  $t_2 = 2 = 2 =$  Mathematical Analysis Tired of walking up the stairs, an engineering student designs an ingenious device for reaching her third floor dorm room. An block is attached to a rope that passes over the outer diameter of a 0.7 m outer diameter, disk-shaped compound pulley. The student holds a second rope, wrapped around the inner 0.35 m diameter of the pulley. When the block is released, the student is pulled up to her dorm room, 8 m off the ground, in 11.2 s. Motion Information Student Event 1:  $t_1 = r_1 = v_1 = a_{12} =$  Event 2:  $t_2 = r_2 = v_2 =$  Pulley Event 1:  $t_1 = 1 = 1 =$  12 = Event 2:  $t_2 = 2 = 2 =$  Block Event 1:  $t_1 = r_1 = v_1 = a_{12} =$  Event 2:  $t_2 = r_2 = v_2 =$  Mathematical Analysis The device at right is used to lift a heavy load. The free rope is attached to a truck that, from rest, accelerates to the right with a constant acceleration. The block is lifted 25 m in 45 s. The inner and outer pulley diameters are 0.40 m and 0.90 m, respectively. Motion Information Truck Event 1:  $t_1 = r_1 = v_1 = a_{12} =$  Event 2:  $t_2 = r_2 = v_2 =$  Pulley Event 1:  $t_1 = 1 = 1 =$  12 = Event 2:  $t_2 = 2 = 2 =$  Block Event 1:  $t_1 = r_1 = v_1 = a_{12} =$  Event 2:  $t_2 = r_2 = v_2 =$  Mathematical Analysis The device at right is used to lift a load quickly. The free rope is attached to a truck that, from rest, accelerates to the right with a constant acceleration. The block is lifted 25 m in 15 s. The inner and outer pulley diameters are 0.40 m and 0.90 m, respectively. Motion Information Truck Event 1:  $t_1 = r_1 = v_1 = a_{12} =$  Event 2:  $t_2 = r_2 = v_2 =$  Pulley Event 1:  $t_1 = 1 = 1 =$  12 = Event 2:  $t_2 = 2 = 2 =$  Block Event 1:  $t_1 = r_1 = v_1 = a_{12} =$  Event 2:  $t_2 = r_2 = v_2 =$  Mathematical Analysis Dynamics Concepts and Principles To study the dynamics of an arbitrary rigid body we will break the motion down into a pure translation of the CM and a pure rotation about the CM. We will use particle dynamics, i.e., Newton's second law applied to the CM of the object, to study the translational portion of the motion. The study of the rotational portion of the motion requires a pair of new concepts. We will "invent" these concepts through the use of an analogy with linear dynamics. In linear dynamics, Newton's second law states that the linear acceleration of an object is proportional to the total force acting on the object and inversely proportional to the mass, or inertia, of the object. It would seem plausible that the angular acceleration of an object would depend on analogous concepts in the same manner. We will replace the concept of force, often thought of as the push or pull applied to an object, with a quantity measuring the twist applied to an object. We will call this new quantity torque, symbolized  $\tau$ . We will replace the concept of mass, the measure of the resistance of the object to changes in its linear velocity, with a quantity measuring the resistance of the object to changes in its angular velocity. We will call this new quantity rotational inertia, symbolized  $I$ . In summary, Before we go any further, however, let's define these new concepts more clearly. Torque In simple English, torque measures the twist applied to an object. The question remains, however, how do we quantify twist? Let's examine a common device used to generate twist, a wrench. The magnitude, location, and orientation of the force applied to the wrench by the person's hand are indicated. Each of these three parameters effects the amount of twist the person delivers to the wrench (and therefore to the bolt). If you've turned many bolts in your life, two things about this person's bo ). Remembering that a Newton is equivalent to kg.m.s<sup>-2</sup> leads to the units of rotational inertia being kg.m<sup>2</sup>. Thus, rotational inertia must be the product of a mass and a distance squared. It seems plausible (there's that word again!) that the distance that is squared in the relationship for rotational inertia is the distance from the rotation axis. The farther a piece of mass is from the rotation axis, the more difficult it is to give the object an angular acceleration. Enough with the "plausibilities", let's finally just define the rotational inertia to be: Imagine the object of interest divided into a large number of small chunks of mass, each with mass  $m$ . Each chunk of mass is a distance  $r$  from the rotation axis. If you take the product of the mass of each chunk and the distance of the chunk from the rotation axis, squared, and sum this quantity over all the chunks of mass, you have the rotational inertia of the object of interest. This summation has been done for a number of common objects and the results are tabulated in the next section. In summary, we now have quantitative relationships for measuring torque and rotational inertia, and the hypothesis that these two quantities are related in a manner analogous to Newton's second law, Rotational Inertia of Common, Uniform Solids Hoop about axis through CM Cylinder about axis through CM Solid sphere about axis through CM Hollow sphere about axis through CM Thin rod about axis through CM Rectangular plate about axis through CM Dynamics Analysis Tools Applying Newton's Second Law to Circular

Motion Investigate the situation below, in which an object travels on a circular path. (Note that this situation does not require the new concepts of torque and rotational inertia.) During a high speed auto chase in San Francisco, a 1745 kg police cruiser traveling at constant speed loses contact with the road as it goes over the crest of a hill with radius of curvature 80 m. A free-body diagram for the car at the instant it's at the crest of the hill is sketched below. Since the car is traveling along a circular path, it's best to analyze the situation using polar coordinates. Remember, in polar coordinates the radial direction points away from the center of the circular path. In general, you should apply Newton's second law independently in the radial (R) and tangential ( ) directions, although since the cruiser is moving at constant speed, there is no acceleration in the tangential direction and therefore no net force in the tangential direction. R direction Since the cruiser loses contact with the road,  $F_{road} = 0$  N. and therefore For the car to lose contact with the road, it must be traveling at a speed of 28 m/s or greater when it reaches the crest of the hill. At lower speeds, the car would remain in contact with the road. Since it is often useful to directly relate the radial acceleration of an object to its velocity, let's construct a direct relationship between these two variables: This relationship, although mathematically equivalent to , is in a more "useful" form. Another Circular Motion Scenario A 950 kg car, traveling at a constant 30 m/s, safely makes a left hand-turn with radius of curvature 75 m. First, let's draw a pair of free-body diagrams for the car, a side-view (on the left) and a rear-view (on the right) . The free-body diagram on the left is a side-view of the car. Notice that the upward direction is the y-direction, and the forward direction, tangent to the turn, is the direction. The free-body diagram on the right is a rear-view of the car. This is what you would see if you stood directly behind the car. Notice that the upward direction is still the y-direction, and the horizontal direction, perpendicular to the direction of travel and hence directed radially outward, is the R-direction. The frictional force indicated is perpendicular to the tread on the tire. This force causes the car to accelerate toward the center of the turn. Remember, if the car is going to travel along a circular path, it must have an acceleration directed toward the center of the circle. Something has to be supplying the force that creates this acceleration. This something is the static friction between the tire and the road that acts to prevent the car from sliding out of the turn. Since the car has no velocity in the radial direction, the frictional force that points in this direction must be static! Now that we have all that straightened out (maybe), let's apply Newton's Second Law. y direction R direction Thus, to safely make this turn requires at least 11400 N of static friction. Using this value, I should be able to compute the minimum coefficient of friction necessary for the car to safely round this turn at this speed. Although this is a large value for the coefficient of static friction, it is an attainable value for a sports car with performance tires. Applying Newton's Second Law in Translational and Rotational Form - I Investigate the scenario described below. The robotic arm at right consists of a pair of hydraulic jacks, each attached 10 cm from the pin "elbow". The elbow is 15 cm from the extreme edge of the 40 kg, 0.90 m long "forearm". The forearm and attached 90 kg load are held stationary at an angle of  $20^\circ$  below horizontal. Since the jacks primarily exert forces through extension, the front, or "biceps", jack is exerting negligible force. To study the dynamics of this situation, we will first need a free-body diagram of the forearm (and the attached load). The forces due to the jack, the rope, and gravity should not need much explanation. However, the forces due to the pin elbow may need some explanation. First, the free-body diagram drawn at left is of the forearm, not including the pin that serves as the elbow. Thus, the pin is external to the object of interest, and thus its interactions with the object are forces that must be indicated on the diagram. If you were to include the pin as part of the forearm, then the upper part of the arm would be external to the object of interest (pin plus forearm) and its interactions with the "pin plus forearm" would have to be indicated as forces on the diagram. Both of these approaches are completely valid. Second, the direction of the force of the pin on the forearm may not be obvious. Does the pin push straight down on the forearm, pull up on the elbow, or push at some unspecified angle? The only thing that's clear is that this force is directed somewhere in the xy-plane. If the force is in the xy-plane, then it must have components along both the x- and y-axis (although one of these components may turn out to have a magnitude of zero). Thus, to handle the generality of the situation you should include both an x- and y-component for the force of the pin on the forearm. For convenience, I'll draw these forces as pointing in the positive x and y directions. (If my guess is wrong, the mathematics will tell me.) Now that we have a free-body diagram, we can apply Newton's second law, both the linear form (in the x- and y-directions) and the rotational form (in the -direction). x-direction y-direction Since the forearm is held stationary, both  $a_x$  and  $a_y$  are equal to zero. Also,  $F_{rope}$  is equal to the force of gravity on the load, 882 N. Obviously, we need another equation in order to determine the two unknown forces. The obvious choice is Newton's second law in rotational form. To use the rotational form of Newton's second law, we must determine the torque due to each force. (Remember, the angle, , in the relation for torque is the angle between  $r$  (oriented along the forearm) and  $F$ .) Pin, Y Jack Gravity Rope Finally, let's apply Newton's second law in rotational form, paying careful attention to the algebraic sign of each torque. Note that our coordinate system indicates that all torques acting counterclockwise are positive. Therefore, all torques acting clockwise are negative. Since the forearm is held stationary, is equal to zero. The jack must exert a force of 7796 N to hold the forearm stationary. Plugging this value into our y-equation yields: The pin must exert a force of 9070 N upwards, since the mathematics determined that the algebraic sign of  $F_{Pin,y}$  was positive. Applying Newton's Second Law in Translational and Rotational Form - II A 75 kg man is attached to a rope wrapped around a 35 kg disk-shaped pulley, with inner and outer diameters 0.60 m and 0.90 m,

respectively. The man is initially at rest. The brake shoe is pressed against the pulley with a force of 500 N. The coefficient of friction between the brake shoe and the pulley is (0.9, 0.8). First, we have to determine whether the force applied to the brake shoe is sufficient to hold the man stationary. If it's not, we have to find the man's acceleration. To accomplish this, we need free-body diagrams for both the man and the pulley. We'll use counterclockwise as the positive  $\theta$ -direction and down as the positive  $y$ -direction. The pin, or axle, at the center of the pulley exerts forces in the  $x$ - and  $y$ -direction, although the exact directions may be unknown. We'll just label these forces as  $F_{pin,x}$  and  $F_{pin,y}$ . The frictional force that acts on the right edge of the pulley acts in a way to prevent the man from falling (or, if he does fall, to slow down the man's fall). To begin my analysis, I'll assume that the man is stationary, solve for the value of the frictional force required to keep him stationary, and then determine whether this frictional force is possible given the applied force of 500 N and the coefficient of static friction. Let's apply Newton's second law to the man in the  $y$ -direction. (Remember, we are assuming he is not falling.) Now look at Newton's second law applied to the pulley in the  $\theta$ -direction. Notice that all of the other forces acting on the pulley do not exert torques on the pulley. They are either located at  $r = 0$  m or have angular orientations of  $\theta = 180^\circ$ . Therefore, to hold the man stationary requires 490 N of friction. However, Therefore, the man must accelerate downward. Now that we know the man must fall, let's write the same two equations as before, although this time the accelerations (both angular for the pulley and linear for the man) are not zero. For the pulley: This assumes that the rotational inertia of a "compound" pulley, one with more than one location where a rope can be wrapped, is the same as a regular pulley, and that the outermost radius of the pulley determines the inertia. Since the man is falling, the frictional force is now kinetic, and Thus our "pulley" equation becomes: Examining Newton's second law for the man, We now have two equations with three variables. However, the angular acceleration of the pulley and the linear acceleration of the man are directly related. Since the rope that the man is attached to is wrapped 0.3 m from the center of the pulley, the man accelerates at the same rate as the tangential acceleration of a point on the pulley 0.3 m from the center. Thus, Using this relationship allows me to rewrite the two Newton's second law equations as: This pair of equations has the solution The man falls with this acceleration and the rope exerts this force (upward on the man and downward on the pulley).

Dynamics Activities

Six artificial satellites of mass,  $m$ , circle a space station at constant speed,  $v$ .  $m$   $v$  A 200 kg 160 m/s B 100 kg 160 m/s C 400 kg 80 m/s D 800 kg 40 m/s E 200 kg 120 m/s F 300 kg 80 m/s a. If the distance between the space station and the satellites is the same for all satellites, rank the magnitude of the net force acting on each satellite. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. If the period is the same for all satellites, rank the magnitude of the net force acting on each satellite. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Six artificial satellites of mass,  $m$ , circle a space station at distance,  $R$ .  $m$   $R$  A 200 kg 5000 m B 100 kg 10000 m C 400 kg 2500 m D 800 kg 5000 m E 100 kg 2500 m F 300 kg 7500 m a. If the tangential speed is the same for all satellites, rank the magnitude of the net force acting on each satellite. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. If the period is the same for all satellites, rank the magnitude of the net force acting on each satellite. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Six rollercoaster carts pass over semi-circular "bumps". The mass of each cart (including passenger),  $M$ , and the force of the track on the cart at the apex of each bump,  $F$ , are given below.  $M$   $F$  A 200 kg 400 N B 100 kg 400 N C 400 kg 200 N D 800 kg 800 N E 100 kg 800 N F 300 kg 300 N a. If the radius of each bump is the same, rank the speed of the cart as it passes over the apex of the bump. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. If the speed of each cart is the same at the apex, rank the radius of each bump. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A rollercoaster track has six semi-circular "dips" with different radii of curvature,  $R$ . The rollercoaster cart rides through each dip at a different speed,  $v$ .  $R$   $v$  A 30 m 16 m/s B 60 m 16 m/s C 15 m 8 m/s D 30 m 4 m/s E 15 m 12 m/s F 45 m 4m/s Rank the magnitude of the force of the rollercoaster track on the cart at the bottom of each dip. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A rollercoaster track has a semi-circular "bump" of radius of curvature  $R$ . A rollercoaster cart (including passenger) of total mass  $M$  rides over the bump.  $M$   $R$  A 200 kg 30 m B 100 kg 60 m C 400 kg 15 m D 800 kg 30 m E 100 kg 15 m F 300 kg 45 m Rank the minimum velocity necessary for the cart to momentarily lose contact with the track at the top of the bump. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: The wrench below has six equal magnitude forces acting on it. Rank these forces on the basis of the magnitude of the torque they apply to the wrench, measured about an axis centered on the bolt. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Two identical spheres are



attached together by a thin rod. The rod lies on a line connecting the centers-of-mass of the two spheres. The length of the rod is equal to the diameter of each sphere. The object has six equal magnitude forces acting on it at the locations shown. a. Rank these forces on the basis of the torque they apply to the object, measured about axis 1. Let the counterclockwise direction be positive and rank positive torques as larger than negative torques. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank these forces on the basis of the torque they apply to the object, measured about axis 2. Let the counterclockwise direction be positive and rank positive torques as larger than negative torques. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: An artist constructs the mobile shown below. The highest two crossbars are 3 units long, and are hung from a point  $\frac{1}{3}$  along their length. The lower three crossbars are 2 units long, and are hung from their midpoint. In this configuration, the mobile is perfectly balanced. Rank the masses of the six hanging shapes. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: An artist constructs the mobile shown below. The crossbars are either 3 units long, and hung from a point  $\frac{1}{3}$  along their length, or are 2 units long, and hung from their midpoint. In this configuration, the mobile is perfectly balanced. Rank the masses of the six hanging shapes. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: The six platforms below are initially at rest in deep space. The indicated forces act on the platforms at their endpoints, quarter-points, or midpoints. All forces have the same magnitude. Rank these platforms on the basis of the torque that acts on them, measured about their CM. Let the counterclockwise direction be positive and rank positive torques as larger than negative torques. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: The six platforms below are initially at rest in deep space. The indicated forces act on the platforms at their endpoints, quarter-points, or midpoints. All forces have the same magnitude. Rank these platforms on the basis of the torque that acts on them, measured about their CM. Let the counterclockwise direction be positive and rank positive torques as larger than negative torques. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: The six platforms below are initially at rest in deep space. The indicated forces act on the platforms at their endpoints, quarter-points, or midpoints. All forces have the same magnitude. Rank these platforms on the basis of the force that must be applied to them, at their left-edge, to keep them from rotating. Let the positive direction be upward and rank positive forces as larger than negative forces. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A roadside sign is to be hung from the end of a thin pole, and the pole supported by a single cable. Your design firm brainstorms the following six scenarios. In scenarios A, B, and D, the rope is attached to the pole  $\frac{3}{4}$  of the distance between the hinge and the sign. In C, the rope is attached to the mid-point of the pole. Rank the design scenarios on the basis of the tension in the supporting cable. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Two identical, uniform, solid spheres are attached together by a solid, uniform thin rod. The rod lies on a line connecting the centers-of-mass of the two spheres. The axes A, B, C, and D are in the plane of the page (which contains the centers-of-mass of the spheres and the rod), while axes E and F are perpendicular to the page. Rank the rotational inertia of this object about the axes indicated. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Below are top views of six identical turntables, upon which are fastened different masses. The distance from the center of the turntable to the center of the mass is either zero, one-half the radius of the turntable, or the radius of the turntable. Rank the rotational inertia of these turntable-mass systems. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Below are top views of six identical turntables, upon which are fastened different masses. The distance from the center of the turntable to the center of the mass is either zero, one-half the radius of the turntable, or the radius of the turntable. Rank the rotational inertia of these turntable-mass systems. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A block of mass M is attached to the inner radius (Rinner) of the pulley shown below. A second rope is attached to outer radius (Router) of the pulley. Assume friction in the pulley mount is very small. M Rinner Router A 10 kg 10 cm 20 cm B 5 kg 20 cm 40 cm C 10 kg 5 cm 20 cm D 10 kg 5 cm 10 cm E 20 kg 5 cm 10 cm F 20 kg 20 cm 30 cm Rank these scenarios on the basis of the magnitude of the force that must be applied to the free rope to hold the block stationary. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A 10 kg block is attached to the inner radius (Rinner) of the pulley shown below. A second rope, attached to outer radius (Router) of the pulley, is used to raise the block at constant speed

v. Assume friction in the pulley mount is very small. v Rinner Router A 10 m/s 10 cm 20 cm B 20 m/s 20 cm 40 cm C 5 m/s 5 cm 20 cm D 10 m/s 5 cm 10 cm E 5 m/s 5 cm 10 cm F 10 m/s 20 cm 30 cm Rank these scenarios on the basis of the magnitude of the force that must be applied to the free rope to raise the block at the speed indicated. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A rollercoaster cart travels through a semicircular “bump” of radius 20 m followed by a “dip” of radius 25 m. At the apex of the bump, the 50 kg passenger momentarily loses contact with her seat, and at the bottom of the dip the bathroom scale she’s sitting on reads 1020 N. a. How fast is the cart moving at the apex of the bump? Free-Body Diagram Mathematical Analysis b. How fast is the cart moving at the bottom of the dip? Free-Body Diagram Mathematical Analysis A 60 kg woman rides in a Ferris wheel of radius 16 m. In order to better understand physics, she takes along a bathroom scale and sits on it. When at the top, the scale reads 540 N. The Ferris wheel maintains the same speed throughout its motion. She checks the scale again at the bottom of the motion. a. What is the angular speed of the Ferris wheel? Free-Body Diagram Mathematical Analysis b. What does the bathroom scale read at the bottom of the motion? Free-Body Diagram Mathematical Analysis A 1.5 kg pendulum bob swings from the end of a 2 m long string. The maximum angle from vertical reached by the pendulum is  $15^\circ$ , and the bob reaches a speed of 1.16 m/s as it passes through its lowest point. a. What is the force exerted on the bob by the string at the maximum angle? Free-Body Diagram Mathematical Analysis b. What is the force exerted on the bob by the string at the lowest point? Free-Body Diagram Mathematical Analysis The 65 kg strange man at right is skateboarding in a half-pipe. The half-pipe is a half-circle with a radius of 5 m. At the instant shown he is 2 m from the bottom of the half-pipe, measured vertically, and is moving at 7.5 m/s. He reaches a speed of 9.5 m/s at the bottom. a. What is the force exerted by the skateboard on the man at the instant shown? Free-Body Diagram Mathematical Analysis b. What is the force exerted by the skateboard on the man at the bottom of the half-pipe? Free-Body Diagram Mathematical Analysis A 2.0 kg tether ball is attached to a vertical pole by a 1.5 m long rope. The ball is swung around the pole at an angle of  $300^\circ$  from vertical. Free-Body Diagram Mathematical Analysis y-direction -direction A 2.0 kg tether ball swings around a vertical pole attached to two 1.5 m long ropes, each at an angle of  $300^\circ$  from vertical. One rope is attached to the top of the ball and the top of the pole, the other rope is attached to the bottom of the ball and the bottom of the pole. The ball is traveling at 3.5 m/s. Free-Body Diagram Mathematical Analysis In the hammer throw, a 7.3 kg steel ball at the end of a 1.22 m wire is swung in an approximately circular path around a thrower’s head. Assume the ball is traveling at a constant speed of 7 m/s. Free-Body Diagram Mathematical Analysis A 1000 kg car rounds a 50 m radius curve at constant speed without slipping. The coefficient of friction between the car’s tires and the road is (0.6,0.5). Free-Body Diagram Mathematical Analysis rear view An 400 m radius, banked, highway curve is designed to allow cars to drive through the curve at a speed of 20 m/s and not slip, even when the road is extremely icy. Free-Body Diagram Mathematical Analysis rear view A 100 m radius, 80 banked, highway curve has just been built. You decide to find maximum velocity with which you can drive your 1200 kg pick-up truck through this curve without sliding up or down the incline. The coefficient of friction between your enormous tires and the road is (0.7, 0.6) . Free-Body Diagram Mathematical Analysis rear view You enter an 80 m radius, 60 banked, highway curve at 30 m/s. The coefficient of friction between your tires and the road is (0.9, 0.8) . Free-Body Diagram Mathematical Analysis rear view In an amusement park ride called the Rotor, a circular, 4.0 m radius room is spun around at high speed and then the floor is removed. The people riding the Rotor feel that they are being pressed against the wall with such a large force that they do not slide down the wall to the floor. (Obviously they do not understand physics.) The coefficient of friction between the riders and the wall is (0.6,0.5). Free-Body Diagram Mathematical Analysis The roadside sign at right has a mass of 22 kg. It hangs from the end of a 1.6 m long, 14 kg support pole hinged to the wall. A support cable is attached to the pole 1.1 m from the wall. Free-Body Diagram Mathematical Analysis The roadside sign at right has a mass of 22 kg. It hangs from a 1.6 m long, 14 kg support pole hinged to the wall. A support cable is attached to the pole at both 0.4 m and 1.2 m from the wall. The sign is hung directly below the outer support cable. Free-Body Diagram Mathematical Analysis The roadside sign at right has a mass of 22 kg. It hangs from the end of a 1.6 m long, 14 kg support pole hinged to the wall at an angle of  $250^\circ$  above horizontal. A support cable is attached to the pole at a point 1.1 m from the wall along the pole. Free-Body Diagram Mathematical Analysis The 70 kg strange man has built a simple scaffold by placing a 6 m long, 35 kg board on top of two piles of bricks. The brick supports are 2 m apart and centered on the CM of the board. The man’s 10 kg toolbox is 1 m from the right edge of the board. The man is 1.1 m from the left edge of the board when he stops and realizes he probably can’t walk all the way to the left edge of the board. Free-Body Diagram Mathematical Analysis The two 70 kg strange twins have built a simple scaffold by placing a 6 m long, 35 kg board on top of two piles of bricks. The brick supports are 2 m apart and centered on the CM of the board. One twin is at the extreme right edge of the board while the other is 1.2 m from the left edge. Free-Body Diagram Mathematical Analysis The 70 kg strange man has climbed three-quarters of the way up the 20 kg, 10 m long ladder when he stops and realizes that  $50^\circ$  may not be a very safe angle for a ladder. The coefficient of friction between the base of the ladder and the ground is (0.6, 0.5). The friction between the ladder and the wall is negligible. Free-Body Diagram Mathematical Analysis The 75 kg man is falling at 20 m/s, 75 m above the crocodile-infested waters below! In an

attempt to save him, the brake shoe is pressed against the spinning pulley. The coefficient of friction between the brake shoe and the pulley is (0.9, 0.8), and the 35 kg disk-shaped pulley has inner and outer diameters of 0.60 m and 0.90 m, respectively. Motion Information Free-Body Diagrams Object: Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_1 = 1$  = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 = a_2 = 2$  = Mathematical Analysis Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching his third floor dorm room. An 84 kg block is attached to a rope that passes over a 0.7 m diameter, 15 kg, disk-shaped pulley. The student holds the other end of the rope. When the 84 kg block is released, the student is pulled up to his dorm room, 8 m off the ground. Motion Information Free-Body Diagrams Object: Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_1 = 1$  = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 = a_2 = 2$  = Mathematical Analysis Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching her third floor dorm room. An 42 kg block is attached to a rope that passes over the outer diameter of a 0.7 m outer diameter, 15 kg, disk-shaped compound pulley. The student holds a second rope, wrapped around the inner 0.35 m diameter of the pulley. When the 42 kg block is released, the student is pulled up to her dorm room, 8 m off the ground. Motion Information Free-Body Diagrams Object: Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_1 = 1$  = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 = a_2 = 2$  = Mathematical Analysis A 60 kg student lifts herself from rest to a speed of 1.5 m/s in 2.1 s. The boson's chair has a mass of 35 kg. The 20 kg, disk-shaped pulley has a diameter of 1.1 m. Motion Information Free-Body Diagrams Object: Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_1 = 1$  = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 = a_2 = 2$  = Mathematical Analysis The 65 kg man at right is trapped inside a section of large pipe. If that's not bad enough, the pipe begins to roll, from rest, down a 35 m long, 180 incline! The pipe has a mass of 180 kg and a diameter of 1.2 m. (Assume the man's presence inside the pipe has a negligible effect on the pipe's rotational inertia.) The coefficient of friction between the pipe and the ground is (0.5, 0.4). Motion Information Free-Body Diagram Event 1:  $t_1 = r_1 = 1 = v_1 = 1 = a_1 = 1$  = Event 2:  $t_2 = r_2 = 2 = v_2 = 2 = a_2 = 2$  = Mathematical Analysis Conservation Laws Concepts and Principles The Angular Impulse-Angular Momentum Relation We've already used the impulse-momentum relation to analyze situations involving translations through three-dimensional space. The relation is typically applied in its component form: Arguing by analogy, if the change in momentum in the x-, y-, and z-directions is equal to the impulse applied to the object in each direction doesn't it seem plausible that a similar relation would hold in the "θ-direction"? If so, this relation should be constructed between the torques acting on an object, the time interval over which the torques act, and the change in angular velocity of an object. The relation should be: The product of rotational inertia and angular velocity is termed the angular momentum of the object, typically denoted L, and the product of torque and the time interval over which it acts is termed the angular impulse applied to the object. Thus, if no angular impulse is applied to an object, its angular momentum will remain constant. This special case is referred to as angular momentum conservation. However, if an angular impulse is applied to the object, the angular momentum will change by an amount exactly equal to the angular impulse applied. Angular momentum is changed through angular impulse. Incorporating Rotation in the Work-Energy Relation Our previous encounter with the work-energy relation resulted in: Recall from our previous discussion of work-energy that this is not a vector equation, meaning it is not applied independently in each of the coordinate directions. Generalizing this equation to include rotation will involve adding terms to this equation, not creating a separate "angular" energy equation. Recall that we model the motion of an arbitrary rigid body as a superposition of a pure translation of the CM and a pure rotation about the CM. Let's investigate the effect of this model on our calculation of the kinetic energy of the object. The translation portion of the motion is easy. We envision the object as a point particle, localized at the CM of the real object, traveling with the velocity of the CM of the object. Thus, this portion of the motion contributes a kinetic energy, What about the kinetic energy due to the pure rotation of the object about the axis through the CM? Every small chunk of the object, dm, moves in a circle around the CM. Each of these pieces of mass has a velocity magnitude given by  $v = r\omega$ . Thus, the kinetic energy of each piece is  $dK = \frac{1}{2} dm v^2 = \frac{1}{2} dm r^2 \omega^2$ . Therefore, the total kinetic energy due to rotation of all the little pieces is:  $K_{rot} = \int \frac{1}{2} dm r^2 \omega^2 = \frac{1}{2} \omega^2 \int dm r^2 = \frac{1}{2} I \omega^2$ . Combining the kinetic energy due to rotation with the kinetic energy due to translation leads to a total kinetic energy of:  $K = \frac{1}{2} M v_{CM}^2 + \frac{1}{2} I \omega^2$ . and a work-energy relation of:  $\Delta K = W_{ext}$ . Conservation Laws Analysis Tools Applying the Work-Energy Relation including Rotation - I A mischievous child releases his mother's bowling ball from the top of the family's 25 m long, 150 above horizontal driveway. The ball rolls without slipping down the driveway and at the bottom plows into the mailbox. The 6.4 kg ball has a diameter of 24 cm. Let's determine the speed of the ball when it hits the mailbox. To determine this value, we can apply the work-energy relation to the ball between: Event 1: The instant the ball is released. Event 2: The instant the ball hits the mailbox. For these two events, work-energy looks like this: where the bottom of the driveway is the zero for gravitational potential energy and the rotational inertia of the bowling ball is taken to be that of a sphere. Now we must carefully determine which, if any, of the forces on the bowling ball do work on the bowling ball. First, we don't have to worry about the force of gravity. The gravitational potential energy function was developed to automatically incorporate the work done by the force of gravity. Second, the contact force can do no work on the ball because the contact force is always perpendicular to the motion of the ball. Finally, what about the force of friction? It does appear, at first glance, that the force of friction is applied over the entire motion of the ball down the driveway. However, let's pay closer attention to the actual point at which the force acts. The frictional force is static in nature, because since the ball rolls without



slipping down the driveway the bottom of the ball is always in static contact with the ground (i.e., there is no relative velocity between the bottom of the ball and the ground). If the bottom of the ball has a velocity of zero, then the force that acts on the bottom of the ball (static friction) can act through no distance. During the instant at which the frictional force acts on a particular point on the bottom of the ball, that point is not moving. That point on the ball only moves when it is no longer in contact with the ground, but by that time the frictional force is acting on a different point. To summarize (and stop saying the same thing over and over), the force of static friction can do no work because it acts on a point that does not move. Thus, our equation simplifies to: The linear and angular velocity of the ball must be related. Since the ball rolls without slipping, the bottom of the ball has no linear velocity. Since the velocity of the bottom of the ball is the sum of the velocity due to translation of the CM and the velocity due to rotation about the CM, the velocity due to rotation must be equal in magnitude to the velocity of the CM: Substituting this into our equation yields: Applying the Work-Energy Relation including Rotation - II

The 60 kg crate is descending at 2.0 m/s when a brake shoe is applied to the 10 kg, 20 cm radius disk-shaped pulley. The coefficient of friction between the shoe and the pulley is (0.9, 0.7). The crate comes to rest after descending 3 m. To determine this function, let's apply the work-energy relation to both the crate and the pulley between: Event 1: The instant the brake shoe is applied. Event 2: The instant the crate comes to rest. In addition to defining the two instants of interest, we'll need free-body diagrams for both the crate and the pulley. Crate where the final position of the crate is the zero for gravitational potential energy. Pulley To find the work on the pulley, note that  $F_{\text{contact}}$  and  $F_{\text{pin}}$  do no work since the distance over which these forces act is zero. However, both  $F_{\text{rope}}$  and  $F_{\text{friction}}$  do work. If the crate falls 3.0 m, both of these forces act over a displacement of 3.0 m. However, note that this displacement is in the same direction as  $F_{\text{rope}}$  ( $= 0$ ) on the right side of the pulley but in the opposite direction on the other side of the pulley (opposite to  $F_{\text{friction}}$  and therefore  $= 180^\circ$ .) Now note that since the brake shoe does not accelerate, the external force applied to the shoe,  $F$ , is the same magnitude as the contact force between the shoe and the pulley. The brake shoe must have been pressed against the pulley with a force of 902 N in order to stop the crate in 3 m. Applying the Impulse-Momentum Relations (Linear and Angular)

A 30 kg child running at 4.0 m/s leaps onto the outer edge of an initially stationary merry-go-round. The merry-go-round is a flat disk of radius 2.4 m and mass 70 kg. Ignore the frictional torque in the bearings of the merry-go-round. Let's imagine we're interested in determining the final angular velocity of the merry-go-round after the child has safely come to rest on its surface. To determine this angular velocity, we can apply the impulse-momentum relations to both the child and the merry-go-round between: Event 1: The instant before the child lands on the merry-go-round Event 2: The instant after the child comes to rest on the merry-go-round. The child and the merry-go-round interact via some unknown magnitude force. (Contrary to the free-body diagram, this force does not necessarily act on the child's disembodied head.) There are also other forces due to gravity and supporting structures that act on both the child and the merry-go-round, however, these forces are in the vertical direction and supply no torque. Child x-direction Merry-Go-Round - direction

The final velocity of the child and final angular velocity of the merry-go-round are related because at this instant the child is at rest (hanging on to) the merry-go-round. Therefore, Substituting this into our angular equation yields: By Newton's Third Law, the force of the child on the merry-go-round and the force of the merry-go-round on the child must be equal in magnitude. Therefore our two equations can be summed and the impulses cancel. This gives: The child is slowed from 4 m/s to 1.85 m/s by jumping onto the merry-go-round. The merry-go-round, however, is accelerated from rest to:

Conservation Laws Activities For each of the scenarios described below, indicate the amount of linear kinetic energy, rotational kinetic energy, and gravitational potential energy in the system at each of the events listed. Use a consistent scale for each motion. Set the lowest point of each motion as the zero-point of gravitational potential energy.

a. A mischievous child releases his mother's bowling ball from the top of the family's 15 m long, 80 above horizontal driveway. The ball rolls without slipping down the driveway and at the bottom plows into the mailbox.

b. A yo-yo of inner diameter 0.70 cm and outer diameter 8.0 cm is released from rest. The string is 0.80 m long. Assume the string does not slip on the inner diameter of the yo-yo as the yo-yo falls.. For each of the scenarios described below, indicate the amount of linear kinetic energy, rotational kinetic energy, and gravitational potential energy of each object at each of the events listed. Use a consistent scale for each motion. Set the initial positions of the objects as the zero-points of gravitational potential energy.

a. Tired of walking up the stairs, an 80 kg engineering student (S) designs an ingenious device for reaching his third floor dorm room. An 84 kg block (B) is attached to a rope that passes over a 0.70 m diameter, 15 kg, disk-shaped pulley (P). The student holds the other end of the rope. When the block is released, the student is pulled up to his dorm room.

b. A 75 kg man (M) is falling at 20 m/s, 75 m above crocodile infested waters! He holds a rope attached to a 15 kg simple pulley (P). In an attempt to save him, a brake shoe is pressed against the spinning pulley. The man is saved, but barely. For each of the scenarios described below, indicate the amount of linear kinetic energy, rotational kinetic energy, and gravitational potential energy of each object at each of the events listed. Use a consistent scale throughout both motions. Set the initial positions of the objects as the zero-points of gravitational potential energy

a. In a horizontal skiing device, the skier begins from rest 35 m from the end of the skiing run. The skier (S) has a mass of 75 kg, the block (B) has a mass of 50 kg, and the simple pulley (P) has a mass of 15 kg. The coefficient of friction is extremely small.

b. In an inclined skiing device, the skier begins from rest 35 m from the end of the 200 above horizontal

inclined skiing run. The skier (S) has a mass of 75 kg, the block (B) has a mass of 50 kg, and the simple pulley (P) has a mass of 15 kg. The coefficient of friction is extremely small. A disk-shaped pulley of mass  $M$  and radius  $R$  is rotating at angular velocity  $\omega$ . The friction in its bearings is so small that it can be ignored. A brake shoe is pressed against the pulley in order to stop it. In all cases, the brake shoe is pressed against the pulley with the same force and the coefficient of friction between the brake shoe and the pulley is the same.  $M$   $R$   $\omega$  A 10 kg 0.8 m 12 rad/s B 10 kg 0.4 m 24 rad/s C 20 kg 0.4 m 6 rad/s D 40 kg 0.2 m 3 rad/s E 20 kg 0.2 m 3 rad/s F 30 kg 0.6 m 8 rad/s a. Rank the scenarios below on the amount of time it takes to stop the pulley. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank the scenarios below on the angle through which the pulley rotates before stopping. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A disk-shaped pulley of mass  $M$  and radius  $R$  is rotating at angular velocity  $\omega$ . The friction in its bearings is constant and ultimately causes the pulley to stop rotating.  $M$   $R$   $\omega$  A 10 kg 0.8 m 12 rad/s B 10 kg 0.4 m 24 rad/s C 20 kg 0.4 m 6 rad/s D 40 kg 0.2 m 3 rad/s E 20 kg 0.2 m 3 rad/s F 30 kg 0.6 m 8 rad/s a. Rank the scenarios below on the amount of time it takes to stop the pulley. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: b. Rank the scenarios below on the angle through which the pulley rotates before stopping. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A sphere of mass  $m$  and radius  $R$  is released from rest at the top of an incline and rolls without slipping down the incline. All spheres are released from rest from the same location on the incline.  $m$   $R$  A 1 kg 30 cm B 2 kg 60 cm C 3 kg 10 cm D 2 kg 15 cm E 1 kg 15 cm F 2 kg 30 cm Rank the velocity of the sphere at the bottom of the incline. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: Six different equal-mass objects are released from rest from the same location at the top of an incline and roll without slipping down the incline. Object  $R$  A solid sphere 30 cm B hollow sphere 60 cm C solid disk 10 cm D hoop 10 cm E solid cylinder 20 cm Rank the velocity of the objects at the bottom of the incline. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: A merry-go-round of radius  $R$  is rotating at constant angular speed. The friction in its bearings is so small that it can be ignored. A sandbag of mass  $m$  is dropped onto the merry-go-round, at a position designated by  $r$ . The sandbag does not slip or roll upon contact with the merry-go-round.  $m$   $r$  A 10 kg 0.50 R B 10 kg 0.25 R C 20 kg 0.25 R D 40 kg 0.25 R E 10 kg 1.00 R F 15 kg 0.75 R Rank the scenarios on the basis of the angular speed of the merry-go-round after the sandbag “sticks” to the merry-go-round. Largest 1. \_\_\_\_ 2. \_\_\_\_ 3. \_\_\_\_ 4. \_\_\_\_ 5. \_\_\_\_ 6. \_\_\_\_ Smallest \_\_\_\_ The ranking cannot be determined based on the information provided. Explain the reason for your ranking: The 75 kg man is falling at 20 m/s, 75 m above the crocodile infested waters below! In an attempt to save him, the brake shoe is pressed against the spinning pulley. The coefficient of friction between the brake shoe and the pulley is (0.9, 0.8), and the 35 kg disk-shaped pulley has inner and outer diameters of 0.60 m and 0.90 m, respectively. The man is saved, but barely. Free-Body Diagrams Mathematical Analysis Event 1: Event 2: The 75 kg man is falling at 15 m/s! In an attempt to stop his fall, the brake shoe is pressed against the spinning pulley with a force of 1000 N. The coefficient of friction between the brake shoe and the pulley is (0.9, 0.8), and the 35 kg disk-shaped pulley has inner and outer diameters of 0.60 m and 0.90 m, respectively. a. How long does it take to stop the man’s fall? b. How far does the man fall? Free-Body Diagrams Mathematical Analysis Event 1: Event 2: Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching his third floor dorm room. An 84 kg block is attached to a rope that passes over a 0.70 m diameter, 15 kg, disk-shaped pulley. The student holds the other end of the rope. When the 84 kg block is released, the student is pulled up to his dorm room in 5.3 s. Free-Body Diagrams Mathematical Analysis Event 1: Event 2: Tired of walking up the stairs, an 80 kg engineering student designs an ingenious device for reaching her third floor dorm room. An 42 kg block is attached to a rope that passes over the outer diameter of a 0.70 m outer diameter, 15 kg, disk-shaped compound pulley. The student holds a second rope, wrapped around the inner 0.35 m diameter of the pulley. When the 42 kg block is released, the student is pulled up to her dorm room, 8.0 m off the ground. Free-Body Diagrams Mathematical Analysis Event 1: Event 2: A 60 kg student lifts herself from rest to a speed of 1.5 m/s in 2.1 s. The boson’s chair has a mass of 35 kg. The 20 kg, disk-shaped pulley has a diameter of 1.1 m. Free-Body Diagrams Mathematical Analysis Event 1: Event 2: A mischievous child releases his mother’s bowling ball from the top of the family’s 15 m long, 80 above horizontal driveway. The ball rolls without slipping down the driveway and at the bottom plows into the mailbox. The 6.4 kg ball has a diameter of 21.6 cm. Free-Body Diagram Mathematical Analysis Event 1: Event 2: The 65 kg strange man at right is trapped inside a section of large pipe. If that’s not bad enough, the pipe begins to roll from rest down a 35 m long, 180 incline! The pipe rolls down the hill without slipping. The pipe has a mass of 180 kg and a diameter of 1.2 m. (Assume the man’s presence inside the pipe has a negligible effect on the pipe’s rotational inertia.) Free-Body Diagram Mathematical Analysis Event 1: Event 2: A 30 kg child running at 3.0 m/s leaps onto the outer edge of an initially stationary merry-go-round. The

merry-go-round is a flat disk of radius 2.4 m and mass 50 kg. Ignore the frictional torque in the bearings of the merry-go-round. Free-Body Diagrams Mathematical Analysis child (top view) merry-go-round (top view) Event 1: Event 2: A 40 kg child running at 4.0 m/s leaps onto the outer edge of a merry-go-round initially rotating in the opposite direction. The merry-go-round is brought to rest. The merry-go-round is a flat disk of radius 2.4 m and mass 50 kg. Ignore the frictional torque in the bearings of the merry-go-round. Free-Body Diagrams Mathematical Analysis child (top view) merry-go-round (top view) Event 1: Event 2: A 30 kg bag of sand is dropped onto a merry-go-round 1.6 m from its center. The merry-go-round was rotating at 3.2 rad/s before the bag was dropped. The merry-go-round is a flat disk of radius 2.4 m and mass 50 kg. Ignore the frictional torque in the bearings of the merry-go-round. Free-Body Diagrams Mathematical Analysis sandbag (top view) merry-go-round (top view) Event 1: Event 2: A 30 kg bag of sand is dropped onto a merry-go-round. The merry-go-round was rotating at 2.2 rad/s before the bag was dropped, and 1.3 rad/s after the bag comes to rest. The merry-go-round is a flat disk of radius 2.4 m and mass 60 kg. Ignore the frictional torque in the bearings of the merry-go-round. Free-Body Diagrams Mathematical Analysis sandbag (top view) merry-go-round (top view) Event 1: Event 2: Selected Answers

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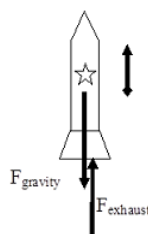
## 02. Time-Dependent Forces

### 1. Time-Dependent Forces

### Time-Dependent Forces

A 2.0 kg toy rocket is fitted with an engine that provides a thrust roughly modeled by the function  $F(t) = (60 \text{ N/s}) t - (15 \text{ N/s}^2) t^2$ , for  $0 < t < 4.0 \text{ s}$ , and zero thereafter. The rocket is launched directly upward.

Of course, the first step to analyzing any situation involving forces is to construct a free-body diagram. Below is a free-body diagram for the rocket during the time interval  $0 < t < 4.0 \text{ s}$ .



Given that you are not presented with information pertaining to the frictional force acting on the rocket, you must analyze the rocket's trajectory in a hypothetical, friction-free environment.

From Newton's second law:

$$\begin{aligned}\Sigma F &= ma \\ F_{\text{exhaust}} - F_{\text{gravity}} &= ma \\ 60t - 15t^2 - 2(9.8) &= 2a \\ -7.5t^2 + 30t - 9.8 &= a\end{aligned}$$

Thus, the acceleration of the rocket is not constant, and the velocity and position of the rocket must be determined by integrating the acceleration.

Before we do this, however, note that the acceleration of the rocket is initially directed downward. (At  $t = 0 \text{ s}$ ,  $a = -9.8 \text{ m/s}^2$ .) This is because the supporting force exerted on the rocket by the launch platform has been ignored. This force would hold the rocket in place until the force of the exhaust gases on the rocket is equal to, and then exceeds, the force of gravity on the rocket. In reality, the acceleration of the rocket is zero until  $F_{\text{exhaust}} = F_{\text{gravity}}$ . This occurs when:

$$\begin{aligned}F_{\text{exhaust}} &= F_{\text{gravity}} \\ 60t - 15t^2 &= 2(9.8) \\ -15t^2 + 60t - 19.6 &= 0 \\ t &= 0.36 \text{ s}\end{aligned}$$

The correct motion information is tabulated below.

Event 1: The engine is ignited	Event 2: The rocket leaves the launch pad	Event 3: The thrust ends	Event 4: The rocket reaches its apex
$t_1 = 0 \text{ s}$	$t_2 = 0.36 \text{ s}$	$t_3 = 4.0 \text{ s}$	$t_4 =$
$r_1 = 0 \text{ m}$	$r_2 = 0 \text{ m}$	$r_3 =$	$r_4 =$
$v_1 = 0 \text{ m/s}$	$v_2 = 0 \text{ m/s}$	$v_3 =$	$v_4 = 0 \text{ m/s}$
$a_1 = 0 \text{ m/s}^2$	$a_2 = 0 \text{ m/s}^2$	$a_3 =$	$a_4 = -9.8 \text{ m/s}^2$

Between event 2 and 3, the rocket's acceleration is given by the function above:

$$a(t) = -7.5t^2 + 30t - 9.8$$

We can integrate the acceleration to determine the velocity,

$$v(t) = \int a(t) dt$$

$$v(t) = \int (-7.5t^2 + 30t - 9.8) dt$$

$$v(t) = -2.5t^3 + 15t^2 - 9.8t + C$$

Since we know  $v = 0$  m/s when  $t = 0.36$  s, we can determine the integration constant:

$$v(0.36) = -2.5(0.36)^3 + 15(0.36)^2 - 9.8(0.36) + C = 0$$

$$C = 1.70 \text{ m/s}$$

Therefore,

$$v(t) = -2.5t^3 + 15t^2 - 9.8t + 1.7$$

and when the thrust ends, at  $t_3 = 4.0$  s,  $v_3 = 42.5$  m/s.

To find the position of the rocket when the thrust ends, integrate the velocity function:

$$r(t) = \int v(t) dt$$

$$r(t) = \int (-2.5t^3 + 15t^2 - 9.8t + 1.7) dt$$

$$r(t) = -0.625t^4 + 5t^3 - 4.9t^2 + 1.7t + D$$

Since we know  $r = 0$  m when  $t = 0.36$  s, we can determine the integration constant:

$$r(0.36) = -0.625(0.36)^4 + 5(0.36)^3 - 4.9(0.36)^2 + 1.7(0.36) + D = 0$$

$$D = -0.200$$

Therefore,

$$r(t) = -0.625t^4 + 5t^3 - 4.9t^2 + 1.7t - 0.200$$

and when the thrust ends, at  $t_3 = 4.0$  s,  $r_3 = 87.2$  m.

Between event 3 and 4 the acceleration is constant, so we can use our constant-acceleration kinematic equations:

$$\begin{array}{ll} v_4 = v_3 + a_{34} (t_4 - t_3) & r_4 = r_3 + v_3 (t_4 - t_3) + \frac{1}{2} a_{34} (t_4 - t_3)^2 \\ 0 = 42.5 + (-9.8)(t_4 - 4) & r_4 = 87.2 + 42.5 (t_4 - 4) + \frac{1}{2} (-9.8)(t_4 - 4)^2 \\ t_4 = 8.34 \text{ s.} & r_4 = 180 \text{ m.} \end{array}$$

Thus, the rocket would achieve a maximum height of 180 m in a friction-free environment.

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## 03. Applying Newton's Second Law to Circular Motion

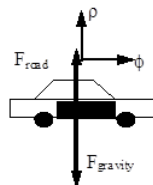
1. Applying Newton's Second Law to Circular Motion
2. Another Circular Motion Scenario

### Applying Newton's Second Law to Circular Motion

Investigate the situation below, in which an object travels on a circular path.

*During a high speed auto chase in San Francisco, a 1745 kg police cruiser traveling at constant speed loses contact with the road as it goes over the crest of a hill with radius of curvature 80 m.*

A free-body diagram for the car at the instant it's at the crest of the hill is sketched below.



Since the car is traveling along a circular path, it's best to analyze the situation using polar coordinates. Remember, in polar coordinates the radial direction points away from the center of the circular path.

In general, you should apply Newton's second law *independently* in the radial (r) and tangential (t) directions, although since the cruiser is moving at constant speed, there is no acceleration in the tangential direction and therefore no net force in the tangential direction.

$\rho$ -direction

$$\Sigma F = ma$$

$$F_{road} - F_{gravity} = ma_{\rho}$$

$$F_{road} - F_{gravity} = m(-R\omega^2)$$

Since the cruiser loses contact with the road,  $F_{road} = 0$  N.

$$0 - 1745(9.8) = -1745(80)\omega^2 \quad (1)$$

$$\omega = 0.35 \text{ rad/s} \quad (2)$$

Since  $v = R\omega$ ,

$$v = (80\text{m})(0.35 \text{ rad/s}) \quad (3)$$

$$v = 28 \text{ m/s} \quad (4)$$

For the car to lose contact with the road, it must be traveling at a speed of 28 m/s or greater when it reaches the crest of the hill. At lower speeds, the car would remain contact with the road.

Since it is often useful to directly relate the radial acceleration of an object to its velocity, let's construct a direct relationship between these two variables:

$$a_{\rho} = R\omega^2 \quad (5)$$

$$a_{\rho} = -R\left(\frac{v}{R}\right)^2 \quad (6)$$

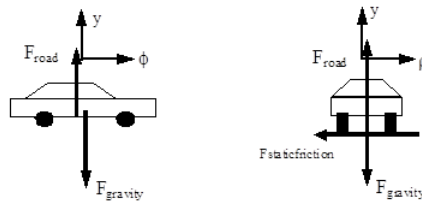
$$a_{\rho} = \frac{-v^2}{R} \quad (7)$$

This relationship, although mathematically equivalent to  $a_{\rho} = R\omega^2$ , is often a more "useful" form.

## Another Circular Motion Scenario

A 950 kg car, traveling at a constant 30 m/s, safely makes a lefthand-turn with radius of curvature 75 m.

First, let's draw a pair of free-body diagrams for the car, a side-view (on the left) and a rear-view (on the right)



The free-body diagram on the left is a side-view of the car. Notice that the upward direction is the y-direction, and the forward direction, tangent to the turn, is the f-direction.

The free-body diagram on the right is a rear-view of the car. This is what you would see if you stood directly behind the car. Notice that the upward direction is still the y-direction, and the horizontal direction, *perpendicular* to the direction of travel and hence directed radially outward, is the r-direction. The frictional force indicated is *perpendicular* to the tread on the tire. This force causes the car to accelerate *toward* the center of the turn. Remember, if the car is going to travel along a circular path, it must have an acceleration directed toward the center of the circle. *Something* has to be supplying the force that creates this acceleration. This something is the static friction between the tire and the road that acts to prevent the car from sliding out of the turn. Since the car has no velocity in the radial direction, the frictional force the points in this direction must be static!

Now that we have all that straightened out (maybe), let's apply Newton's Second Law.

<u>y-direction</u>	<u>ρ-direction</u>
$\Sigma F = ma$	$\Sigma F = ma$
$F_{road} - F_{gravity} = ma_y$	$-F_{staticfriction} = ma_\rho$
$F_{road} - 950(9.8) = 950(0)$	$-F_{staticfriction} = m(-\frac{v^2}{R})$
$F_{road} = 9310N$	$F_{staticfriction} = 950(\frac{30^2}{75})$
	$F_{staticfriction} = 11400N$

Thus, to safely make this turn requires at least 1140 N of static friction. Using this value, I should be able to compute the minimum coefficient of friction necessary for the car to safely round this turn at this speed.

$$F_{stiction} \leq \mu_s F_{road} \quad (8)$$

$$11400 \leq \mu_s (9310) \quad (9)$$

$$\mu_s \geq 1.22 \quad (10)$$

Although this is a large value for the coefficient of static friction, it is an attainable value for a sports car with performance tires.

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## 04. Activities

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## CHAPTER OVERVIEW

### Model 4: The Rigid Body Model

#### Topic hierarchy

- 01. Model Specifics
- 02. Kinematics
  - 01. Concepts and Principles
  - 02. Pure Rotation
  - 03. Rotating and Translating
  - 04. Activities
- 03. Dynamics
  - 01. Concepts and Principles
  - 02. Calculating the Rotational Inertia
  - 03. Applying Newton's Second Law in Translational and Rotational Form - I
  - 04. Applying Newton's Second Law in Translational and Rotational Form - II
  - 05. Applying Newton's Second Law to Rolling Motion
  - 06. Rotational Inertia of Common, Uniform Solids
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## 01. Model Specifics

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### 1. Model Specifics

1. The object is rigid
2. The rotation is planar.
3. The object is classical.

### Model Specifics

We are now ready to consider how the actual size and shape of an object affects its motion. The most immediate ramification of this Model is that the objects we study can now rotate. However, there are still three important restrictions to types of motions will we investigate.

#### The object is rigid

The size and shape of the object under investigation do not change during the motion.

#### The rotation is *planar*.

This restriction may take some explaining. Imagine an object spinning in place. The axis (real or imaginary) about which the object spins is referred to as the rotation axis. Note that all points that lie on the rotation axis do not move. All other points on the object, however, exhibit circular motion around the rotation axis.

If the object, in addition to spinning, is also moving, this model will restrict us to motions where the center-of-mass moves exclusively in the plane perpendicular to the rotation axis. I will refer to this as *planar rotation*.

For example, a yo-yo is an object that typically moves in a plane perpendicular to its rotation axis. Thus, we can study the motion of most yo-yos. The motion of a wheel is typically in a plane perpendicular to its rotation axis. Thus, we can study the motion of most wheels. The motion of the earth around the sun, however, is beyond this model's capabilities since the rotation axis of the earth is tilted relative to the plane of the earth's motion. A wobbling top is also beyond this model's capabilities (although a steadily spinning top is not).

#### The object is classical.

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## SECTION OVERVIEW

### 02. Kinematics

#### Topic hierarchy

#### 01. Concepts and Principles

#### 02. Pure Rotation

#### 03. Rotating and Translating

#### 04. Activities

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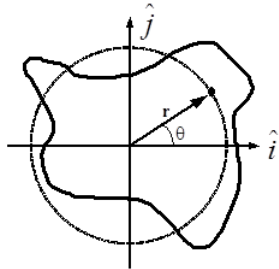
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## 01. Concepts and Principles

1. Rotation about a Fixed Axis (Spinning)
2. Rotation and Translation

### Rotation about a Fixed Axis (Spinning)

Imagine a rigid body constrained to rotate about a fixed axis. Non-physicists would say that the object is *spinning*.



Place the origin of a coordinate system at the location of the rotation axis. Examine an arbitrary point on the object, denoted by the position vector

$$\mathbf{P}(t) = r(\cos \theta(t)\hat{i} + \sin \theta(t)\hat{j})$$

or

$$\mathbf{P}(t) = r\hat{\rho}$$

Notice that as the object spins, this point undergoes circular motion (denoted by the dashed line). Although the actual object may be of irregular shape, as it spins *every point on the object undergoes circular motion*. Moreover, since each and every point on the object has to complete an entire cycle around the rotation axis in the same amount of time, every point must undergo circular motion with the same angular speed ( $\omega(t)$ ) and the same angular acceleration ( $\alpha(t)$ ).

Since every point on a rigid body must have the same angular speed and the same angular acceleration, we will speak of the angular speed and angular acceleration *of* the object, rather than the angular speed and angular acceleration of some point *on* the object.

Given that every point on a spinning object undergoes circular motion, the results from our study of circular motion will be very important in analyzing spinning objects. Recall that with  $Q(t)$  defined as the angular position of an arbitrary point on the rigid-body,

$$\omega(t) = \frac{d\theta(t)}{dt}$$

and

$$\alpha(t) = \frac{d\omega(t)}{dt}$$

Moreover, the velocity and acceleration of any point on a spinning, rigid body can be related to the angular quantities:

$$\mathbf{v}_\theta(t) = r\omega(t)$$

$$\mathbf{v}_\rho(t) = 0$$

and

$$\mathbf{a}_\theta(t) = r\alpha(t)$$

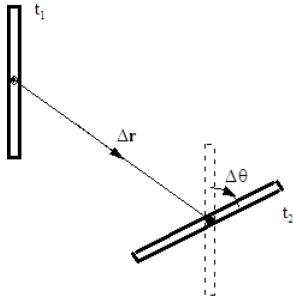
$$\mathbf{a}_\rho(t) = -r\omega^2(t)$$

### Rotation and Translation

To describe the motion of an object undergoing pure rotation (spinning), we have to describe the circular motion each point on the object undergoes. What must we do if the object is simultaneously rotating and translating (moving in a plane perpendicular to the

rotation axis), like a wheel rolling down an incline?

The answer lies in the independence of these two types of motion. In much the same manner as we attacked kinematics in two dimensions by independently analyzing the horizontal and vertical motions, we will attack rigid-body kinematics by independently analyzing the rotational and translational motions. In short, *we will model any motion of a rigid-body as a superposition of a translation of the object's center-of-mass (CM) (which we will analyze by particle kinematics) and a rotation about an axis passing through the CM (which we will analyze by the kinematics of spinning, detailed above).*



For example, examine the motion of the thin rod between  $t_1$  and  $t_2$ . Although the rod may have been spinning crazily through space between these two times, we can model its motion as a superposition of a simple translation of its CM without rotation (denoted by the vector  $\mathbf{Dr}$ ), which leaves the rod in the orientation denoted by the dashed lines, and a simple rotation about an axis through its CM without translation (denoted by  $\mathbf{Dq}$ ), which leaves the rod in its proper, final orientation. If we imagine the time difference ( $t_2 - t_1$ ) shrinking toward zero, hopefully it becomes plausible that we can model any motion through this method.

In summary, to describe the motion of an arbitrary rigid body we will break the motion down into a pure translation of the CM and a pure rotation about the CM. We will use particle kinematics to describe the rotational portion. The velocity (or acceleration) of any point on the object is then determined by the sum of the velocity ( or acceleration) due to the translation and the velocity ( or acceleration) due to the rotation.

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## 02. Pure Rotation

### 1. Pure Rotation

#### 1. Connecting Pure Rotation to Pure Translation

### Pure Rotation

After the off button is pressed, a ceiling fan takes 22 s to come to rest. During this time, it completes 18 complete revolutions.

To analyze this situation, we should first carefully determine and define the sequence of events that take place. At each of these instants, let's tabulate what we know about the motion. Since we are dealing with pure rotation, the relevant kinematic variables are the angular position, velocity, and acceleration of the fan. Also, let's take the direction that the fan is initially rotating to be the positive direction.

Event 1: The 'off' button is pressed	Event 2: The ceiling fan stops
$t_1 = 0 \text{ s}$	$t_2 = 22 \text{ s}$
$\theta_1 = 0 \text{ rad}$	$\theta_2 = 18 (2\pi) = 113 \text{ rad}$
$\omega_1 =$	$\omega_2 = 0 \text{ rad/s}$
$\alpha_1 =$	$\alpha_2 =$

Since no specific information concerning the angular acceleration of the fan is given, let's assume the angular acceleration is constant. (Thus,  $a_1 = a_2 = a_{12}$ .) Since the relationships between angular position, velocity, and acceleration are the same as the relationships between linear position, velocity and acceleration, the kinematic equations for constant linear acceleration must have direct analogies for constant angular acceleration. Thus,

$$\begin{aligned}\omega_2 &= \omega_1 + \alpha_{12}(t_2 - t_1) \\ 0 &= \omega_1 + \alpha_{12}(22) \\ \omega_1 &= -22\alpha_{12}\end{aligned}$$

Now substitute this expression into the other equation:

$$\begin{aligned}\theta_2 &= \theta_1 + \omega_1(t_2 - t_1) + \frac{1}{2}\alpha_{12}(t_2 - t_1)^2 \\ 113 &= 0 + \omega_1(22) + \frac{1}{2}\alpha_{12}(22)^2 \\ 113 &= 22(-22\alpha_{12}) + 242\alpha_{12} \\ 113 &= -242\alpha_{12} \\ \alpha_{12} &= -0.47 \text{ rad/s}^2\end{aligned}$$

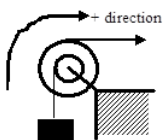
Substitute this result back into the original equation:

$$\begin{aligned}\omega_1 &= -22(-0.47) \\ \omega_1 &= 10.3 \text{ rad/s}\end{aligned}$$

Notice that the sign of the angular acceleration is negative. This indicates that the angular acceleration is in the opposite direction of the angular velocity, as it should be since the fan is slowing down.

### Connecting Pure Rotation to Pure Translation

The device at right is used to lift a heavy load. The free rope is attached to a truck which accelerates from rest at a rate of  $1.5 \text{ m/s}^2$ . The inner radius of the pulley is 20 cm and the outer radius is 40 cm. The load must be raised 15 m.



The coordinate system chosen indicates that the block moving upward, the pulley rotating clockwise, and the truck moving to the right are all positive.

There are three different objects that we should be able to describe kinematically; the truck, the pulley, and the block. Let's tabulate everything we know about each object:

<b>Truck</b> Event 1: Truck begins to move	Event 2: Block raised 15 m	<b>Pulley</b> Event 1: Truck begins to move	Event 2: Block raised 15 m	<b>Block</b> Event 1: Truck begins to move	Event 2: Block raised 15 m
$t_1 = 0 \text{ s}$	$t_2 =$	$t_1 = 0 \text{ s}$	$t_2 =$	$t_1 = 0 \text{ s}$	$t_2 =$
$r_1 = 0 \text{ m}$	$r_2 =$	$\theta_1 = 0 \text{ rad}$	$\theta_2 =$	$r_1 = 0 \text{ m}$	$r_2 = 15 \text{ m}$
$v_1 = 0 \text{ m/s}$	$v_2 =$	$\omega_1 = 0 \text{ rad/s}$	$\omega_2 =$	$v_1 = 0 \text{ m/s}$	$v_2 =$
$a_1 = 1.5 \text{ m/s}^2$	$a_2 = 1.5 \text{ m/s}^2$	$\alpha_1 =$	$\alpha_2 =$	$a_1 =$	$a_2 =$

The truck and block exhibit translational motion, so we only have to tabulate linear variables. The pulley spins, so the relevant variables are the angular variables. Notice that just because the block moves 15 m *doesn't* mean the truck moves 15 m. Also, the acceleration of the block is *not* equal to the acceleration of the truck. However, these variables are related to each other since both objects are attached to the same pulley.

Also notice that we can't currently solve this problem. Each of the objects has three unknown quantities. However, since the kinematics of the three objects are related, we will be able to solve the problem once we've worked out the exact relationship between each object's kinematics.

Let's start with the acceleration. Assuming the rope from the *truck* does not slip on the pulley, the point on the pulley in contact with the rope must be accelerating at the same rate as the truck. Notice that this acceleration is tangent to the pulley, and this rope is located 0.4 m from the center of the pulley. Therefore, from

$$\begin{aligned} a_{\theta} &= r\alpha \\ 1.5 &= 0.4\alpha \\ \alpha &= 3.75 \text{ rad/s}^2 \end{aligned}$$

The pulley must have an angular acceleration of  $3.75 \text{ rad/s}^2$  since it is attached to the truck.

In addition, assuming the rope from the *block* does not slip on the pulley, the point on the pulley in contact with this rope must be accelerating at:

$$\begin{aligned} a_{\theta} &= r\alpha \\ a_{\theta} &= 0.2(3.75) \\ a_{\theta} &= 0.75 \text{ m/s}^2 \end{aligned}$$

Since the point on the pulley attached to the block is accelerating at  $0.75 \text{ m/s}^2$ , the block itself must be accelerating at  $0.75 \text{ m/s}^2$ . In a nutshell, what we've done is used the acceleration of the truck to find the angular acceleration of the pulley, and then used the angular acceleration of the pulley to find the acceleration of the block. This chain of reasoning is very common.

Now that we know the block's acceleration, we can solve the problem. Applying the two kinematic equations to the block yields:

$$\begin{aligned} r_2 &= r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2 \\ 15 &= 0 + 0 + \frac{1}{2}(0.75)t_2^2 & v_2 &= v_1 + a_{12}(t_2 - t_1) \\ t_2 &= 6.32 \text{ s} & v_2 &= 0 + 0.75(6.32) \\ & & v_2 &= 4.74 \text{ m/s} \end{aligned}$$

Now that we know the final speed of the block, we can find the final angular speed of the pulley and the final speed of the truck:

Again assuming the rope from the *block* does not slip on the pulley, the point on the pulley in contact with the rope must be moving at the same rate as the block. Notice that this velocity is tangent to the pulley, and this rope is located 0.2 m from the center of the pulley. Therefore, from

$$\begin{aligned} v_{\theta} &= r\omega \\ 4.74 &= 0.2\omega \\ \omega &= 23.7 \text{ rad/s} \end{aligned}$$

The pulley must have an angular velocity of  $23.7 \text{ rad/s}$  since it is attached to the block.

In addition, assuming the rope from the *truck* does not slip on the pulley, the point on the pulley in contact with this rope must be moving at:

$$\begin{aligned}v_{\phi} &= r\omega \\v_{\phi} &= 0.4(23.7) \\v_{\phi} &= 9.48\text{ m/s}\end{aligned}$$

The truck is moving at 9.48 m/s when the block reaches 15 m.

We can relate the displacement of the block to the angular displacement of the pulley and the displacement of the truck using

$$s = r\theta$$

where  $s$  is (technically) the arc length over which a rope is wrapped around the pulley and  $\theta$  is the angular displacement of the pulley. Of course, the amount of rope wrapped around a pulley is exactly equal to displacement of the object attached to the rope. Therefore, relating the block to the pulley yields

$$\begin{aligned}s &= r\theta \\15 &= 0.2\theta \\ \theta &= 75\text{ rad}\end{aligned}$$

The pulley must have turned through 75 rad since it is attached to the block.

Relating the pulley to the truck results in:

$$\begin{aligned}s &= r\theta \\s &= 0.4(75) \\s &= 30\text{ m}\end{aligned}$$

The truck moved 30 m while the block moved 15 m.

Notice that in all cases, the values of the kinematics variables for the truck are exactly twice the values for the block. This is not a coincidence. Since the truck is attached, all of its kinematic variables will have twice the value. Once you understand why this is the case, you can use this insight to simplify your analysis.

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## 03. Rotating and Translating

### 1. Rotating and Translating

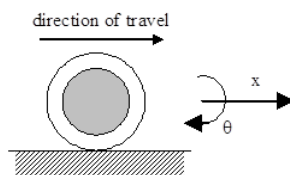
### Rotating and Translating

Accelerating from rest, a Cadillac Sedan de Ville can reach a speed of 25 m/s in a time of 6.2 s. During this acceleration, the Cadillac's tires do not slip in their contact with the road. The diameter of the Cadillac's tires is 0.80 m.

Let's examine the motion of one of the Cadillac's tires. Obviously, the tire both translates and rotates. We will imagine the motion to be a superposition of a pure translation of the CM of the tire and a pure rotation about the CM of the tire.

To analyze the situation, we should define the sequence of events that take place and tabulate what we know about the motion at each event. Since we are dealing with translation as well as rotation, we will need to keep track of both the linear kinematic variables and the angular kinematic variables.

Let's take the positive x direction to be the direction that the Cadillac translates and the positive  $\theta$  direction to be the direction in which the tires rotate.



Event 1: The Cadillac begins to accelerate	Event 2: The Cadillac reaches 25 m/s
$t_1 = 0 \text{ s}$	$t_2 = 6.2 \text{ s}$
$r_1 = 0 \text{ m}$	$r_2 =$
$\theta_1 = 0 \text{ rad}$	$\theta_2 =$
$v_1 = 0 \text{ m/s}$	$v_2 = 25 \text{ m/s}$
$\omega_1 = 0 \text{ rad/s}$	$\omega_2 =$
$a_1 =$	$a_2 =$
$\alpha_1 =$	$\alpha_2 =$

First, let's examine the translational portion of the tire's motion.

$$\begin{aligned}
 r_2 &= r_1 + v_1(t_2 - t_1) + \frac{1}{2}a_{12}(t_2 - t_1)^2 \\
 v_2 &= v_1 + a_{12}(t_2 - t_1) \\
 25 &= 0 + a_{12}(6.2) \\
 a_{12} &= 4.03 \text{ m/s}^2
 \end{aligned}
 \quad
 \begin{aligned}
 r_2 &= 0 + 0 + \frac{1}{2}(4.03)(6.2)^2 \\
 r_2 &= 77.5 \text{ m}
 \end{aligned}$$

What about the rotational kinematics of the tire? The motion of the tire is the superposition of the pure translational motion of the CM and the pure rotational motion about the CM.

$$\begin{aligned}
 \vec{v}_{\text{any point tire}} &= \vec{v}_{\text{due to translation of CM}} + \vec{v}_{\text{due to rotation about CM}} \\
 \vec{v}_{\text{any point tire}} &= \vec{v}_{\text{CM}} + r\omega\hat{\phi}
 \end{aligned}$$

where  $r$  is the distance between the point of interest and the rotation axis, i.e., the distance from the CM.

The key insight into studying the rotation of the tire is to realize that *at any instant the velocity of the point on the tire in contact with the road is zero because the tire never slips in its contact with the road*. Thus, if the CM of the tire is moving forward at

$$\vec{v}_{\text{CM}} = (25 \text{ m/s})\hat{i}$$

the point on the bottom of the tire must be moving backward *relative to the ground* to be zero! Thus, only a very particular value for  $\omega$  will allow the tire to roll without slipping.

From above,

$$\begin{aligned}\vec{v}_{\text{any point on tire}} &= \vec{v}_{\text{CM}} + r\omega\hat{\phi} \\ 0 &= (25\text{ m/s})\hat{i} + r\omega\hat{\phi}\end{aligned}$$

At the bottom of the tire,  $\hat{\phi} = -\hat{i}$ , since motion tangent to the tire at the bottom of the tire is motion to the left. Thus,

$$\begin{aligned}0 &= (25\text{ m/s})\hat{i} - r\omega\hat{i} \\ r\omega\hat{i} &= (25\text{ m/s})\hat{i} \\ (0.4\text{ m})\omega &= (25\text{ m/s}) \\ \omega &= 62.5\text{ rad/s}\end{aligned}$$

Therefore, at the end of the acceleration, the tire has an angular speed of 62.5 rad/s.

I included all of the vector notation to ensure I did the calculation correctly, but hopefully we can understand the result without getting bogged down in notation. Basically, the bottom of the tire can't be moving relative to the ground if the wheel rolls without slipping. Since the CM of the tire is moving forward at 25 m/s, the bottom of the tire must be moving backward at 25 m/s relative to the CM. From the CM frame of reference, the tire is just spinning, and I can apply  $v = r\omega$  to calculate the angular speed of the tire.

$$\begin{aligned}v_{\phi} &= r\omega \\ 25 &= 0.4\omega \\ \omega &= 62.5\text{ rad/s}\end{aligned}$$

We can find the angular acceleration and angular position of the tire by using the linear variables determined above and the method described in the previous example, or by using the equations for constant angular acceleration.

$$\begin{aligned}\theta_2 &= \theta_1 + \omega_1(t_2 - t_1) + \frac{1}{2}\alpha_{12}(t_2 - t_1)^2 \\ \omega_2 &= \omega_1 + \alpha_{12}(t_2 - t_1) & \theta_2 &= 0 + 0 + \frac{1}{2}(10.1)(6.2)^2 \\ 62.5 &= 0 + \alpha_{12}(6.2) & \theta_2 &= 194\text{ rad} \\ \alpha_{12} &= 10.1\text{ rad/s}^2\end{aligned}$$

Thus, the wheel rotates through 194 rad, or 30.8 revolutions, while accelerating.

[1]As stated in the text, to describe the general motion of the tire we will break the motion down into a pure translation of the CM and a pure rotation about the CM. Another way to envision the motion is that the tire is rotating about the point on the tire in contact with the ground. The rotation axis of the tire passes through this contact point. (Imagine the point on the tire in contact with the ground as being glued to the ground and the tire rotating about it.) However, an instant later this point is no longer in contact with the ground, and the rotation axis passes through the *next* point on the tire in contact with the ground. There are advantages and disadvantages to both conceptualizations of the tire's motion, however, we will restrict ourselves to the view that the tire undergoes pure rotation about the CM.

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## 04. Activities

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## 03. Dynamics

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- [01. Concepts and Principles](#)
  - [02. Calculating the Rotational Inertia](#)
  - [03. Applying Newton's Second Law in Translational and Rotational Form - I](#)
  - [04. Applying Newton's Second Law in Translational and Rotational Form - II](#)
  - [05. Applying Newton's Second Law to Rolling Motion](#)
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## 01. Concepts and Principles

1. Concepts and Principles
2. Torque
3. Rotational Inertia

### Concepts and Principles

To study the dynamics of an arbitrary rigid body we will break the motion down into a pure translation of the CM and a pure rotation about the CM. We will use particle dynamics, i.e., Newton's second law applied to the CM of the object, to study the translational portion of the motion. The study of the rotational portion of the motion requires a pair of new concepts. We will "invent" these concepts through the use of an analogy with linear dynamics.

In linear dynamics, Newton's second law states that the linear acceleration of an object is proportional to the total force acting on the object and inversely proportional to the mass, or inertia, of the object. It would seem *plausible* that the angular acceleration of an object would depend on analogous concepts in the same manner.

We will replace the concept of force, often thought of as the push or pull applied to an object, with a quantity measuring the *twist* applied to an object. We will call this new quantity *torque*, symbolized  $\tau$ .

We will replace the concept of mass, the measure of the resistance of the object to changes in its linear velocity, with a quantity measuring the resistance of the object to changes in its angular velocity. We will call this new quantity *rotational inertia*, symbolized  $I$ .

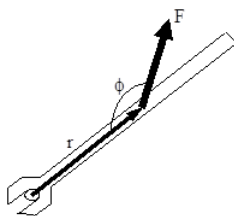
In summary,

$$\Sigma \tau = I\alpha$$

Before we go any further, however, let's define these new concepts more clearly.

### Torque

In simple English, torque measures the twist applied to an object. The question remains, however, how do we *quantify* twist?



Let's examine a common device used to generate twist, a wrench. The magnitude, location, and orientation of the force applied to the wrench by the person's hand are indicated. Each of these three parameters effects the amount of twist the person delivers to the wrench (and therefore to the bolt).

If you've turned many bolts in your life, two things about this person's bolt-turning technique should grab you. First, why is this person applying the force at such a silly angle? She would generate much more twist if she applied the same magnitude force perpendicular to the wrench, rather than at an angle far from  $90^\circ$ . Second, why is she not applying the force at the far edge of the wrench? She would generate far more twist if she applied the same magnitude force at the far edge of the wrench.

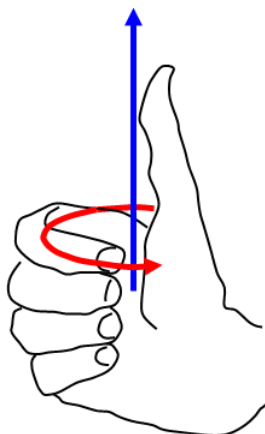
If the preceding paragraph makes sense to you, you understand how to quantify torque. To maximize torque, you should:

1. Apply the force far from the axis of rotation (the bolt).
2. Apply the force perpendicular to the position vector between the axis of rotation and the force.
3. Apply a large magnitude force.

Mathematically, this is summarized by:

$$\tau = rF \sin \phi$$

Note that this function has a maximum when  $r$  is large,  $F$  is large and  $= 90^\circ$ . Torque will also be assigned a direction, either clockwise or counterclockwise, depending upon the direction of the twist applied to the object. Use the Right Hand Rule to help determine the direction of the torque.



## Rotational Inertia

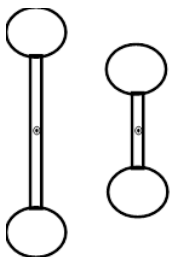
We have constructed a rotational analogy to Newton's second law,

$$\Sigma \tau = I\alpha$$

Our next task is to better define what we mean by  $I$ , the rotational inertia.

The rotational inertia is a measure of the resistance of the object to changes in its angular velocity. Imagine applying the same torque to two objects, initially at rest. After applying the torques for some set amount of time, measure the angular velocity of the objects. The object with the smaller angular velocity has the larger rotational inertia, because it has the larger resistance to angular acceleration.

Since more massive objects are harder to get moving linearly, it seems plausible that rotational inertia should depend on the mass of the object. It also seems plausible that rotational inertia should depend on the shape of the object.



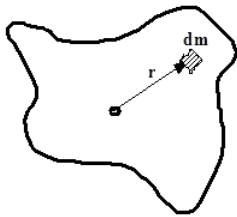
For example, it would be easier to get the smaller dumbbell spinning than the larger dumbbell, even though they have the same total mass.

Let's try to get more quantitative. Examining

$$\Sigma \tau = I\alpha$$

we can see that the units of  $I$  must be the units of torque ( $\text{N}\cdot\text{m}$ ) divided by the units of angular acceleration ( $\text{s}^{-2}$ ). Remembering that a Newton is equivalent to  $\text{kg}\cdot\text{m}\cdot\text{s}^{-2}$  leads to the units of rotational inertia being  $\text{kg}\cdot\text{m}^2$ . Thus, rotational inertia must be the product of a mass and a distance squared.

It seems plausible (there's that word again!) that the distance that is squared in the relationship for rotational inertia is the distance from the rotation axis. The farther a piece of mass is from the rotation axis, the more difficult it is to give the object an angular acceleration.



Enough with the "plausibilities", let's finally just define the rotational inertia to be:

$$I = \int r^2 dm$$

Imagine the object of interest divided into a large number of infinitesimally small chunks of mass, each with mass  $dm$ . Each chunk of mass is a distance  $r$  from the rotation axis. If you take the product of the mass of each chunk and the distance of the chunk from the rotation axis, squared, and sum this quantity over all the chunks of mass, you have the rotational inertia of the object of interest.

In summary, we now have quantitative relationships for measuring torque and rotational inertia,

$$\tau = rF \sin \phi$$

$$I = \int r^2 dm$$

and the hypothesis that these two quantities are related in a manner analogous to Newton's second law,

$$\Sigma \tau = I\alpha$$

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## 02. Calculating the Rotational Inertia

1. [Calculating the Rotational Inertia](#)
2. [The Parallel-Axis Theorem](#)

### Calculating the Rotational Inertia

To fully utilize Newton's second law in rotational form, we must be able to set up and evaluate the integral that determines the rotational inertia. (To be honest, this is a lie. For the vast majority of common shapes, and many quite uncommon shapes, these integrals have already been evaluated. A table of selected results is at the end of this section.) To test our understanding of the relationship for rotational inertia for a thin disk about an axis passing through its center of mass and perpendicular to its circular face.

pic

1. Choose the chunks of mass,  $dm$ , to be ring-shaped. This is because you must multiply each  $dm$  by the distance of the chunk from the rotation axis squared. To facilitate doing this, it's crucial that every point in the chunk be the same distance from the axis, i.e., have the same  $r$ . If the ring-shaped chunk is thin enough, for example  $dr$  (infinitesimally) thick, then this is true.
2. Realize that the mass of the little chunk is directly proportional to its volume, assuming the disk has a constant density. If it does have a constant density, the ratio of the chunk's mass to its volume must be the same as the ratio of the total mass of the disk to its volume.

pic

where  $R$  is the radius of the disk and  $T$  is its thickness. The volume of the ring-shaped chunk,  $dV$ , is equal to the product of the circumference of the ring ( $2\pi r$ ), the thickness of the disk ( $T$ ), and the thickness of the ring ( $dr$ ). Thus,

pic

3. Plug the expression for  $dm$  into

pic

To include all the chunks of mass, the integral must go from  $r = 0$  m up to  $r = R$ .

pic

Thus, the rotational inertia of a thin disk about an axis through its CM is the product of one-half the total mass of the disk and the square of its radius. Notice that the thickness of the disk does not effect its rotational inertia. A consequence of this fact is that a cylinder has the same rotational inertia as a disk, when rotated about an axis through its CM and perpendicular to its circular face.

### The Parallel-Axis Theorem

Now let's imagine we need to calculate the rotational inertia of a thin disk about an axis perpendicular to its circular face and along the edge of the disk. It would be convenient if we could determine the rotational inertia about an axis along the edge using the rotational inertia about an axis through the CM (which we've already calculated). In fact, there is a very convenient method to determine the rotational inertia about any axis *parallel* to an axis through the CM if we know the rotational inertia about an axis through the CM.

Imagine you want to determine the rotational inertia of an arbitrarily shaped object about an arbitrary axis. The solid circle denotes an axis of interest. The two axes are parallel.

Notice that

pic

with

pic

and

pic



Thus,

pic

The rotational inertia about the axis of interest is given by:

pic

Note that  $x_{CM}$ ,  $y_{CM}$ , and  $r_{CM}$  are constants that depend only on the distance between the two axes. Thus,  $x_{CM}$ ,  $y_{CM}$ , and  $r_{CM}$  can be brought outside of the integral.

pic

Now comes the key observation in the derivation. Examine the term  $\int x \, dm$ . Remember that  $x$  is the horizontal distance from the CM. If this distance is integrated over all the chunks of mass,  $dm$ , throughout the entire object, this integral must equal zero because the CM is defined to be in exactly the spot where a mass-weighted average over distance is equal to zero.  $\int x \, dm$  and  $\int y \, dm$  are equal to zero by the definition of CM! (Pretty cool, huh?)

Thus,

pic

Noting that  $\int dm$  is the total mass of the object,  $M$ , and  $\int r^2 \, dm$  is the rotational inertia about the CM,  $I_{CM}$ , then

pic

This result states that the rotational inertia about an axis parallel to an axis through the CM,  $I$ , is equal to the rotational inertia about an axis through the CM,  $I_{CM}$ , plus the product of the total mass of the object and the distance between the axes,  $r_{CM}$ , squared.

To answer the original question, let's determine the rotational inertia of a thin disk about an axis perpendicular to its circular face and along the edge of the disk using the parallel-axis theorem.

We know that the rotational inertia for a thin disk about an axis passing through its center of mass and perpendicular to its circular face is  $\frac{1}{2} MR^2$  and the axis of interest is  $R$ . Thus,

pic

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## 03. Applying Newton's Second Law in Translational and Rotational Form - I

### 1. Applying Newton's Second Law in Translational and Rotational Form - I

#### Applying Newton's Second Law in Translational and Rotational Form - I

Investigate the scenario described below.

pic

*The robotic arm at right consists of a pair of hydraulic jacks, each attached 10 cm from the pin "elbow". The elbow is 15 cm from the extreme edge of the 40 kg, 0.90 m long "forearm". The forearm and attached 90 kg load are held stationary at an angle of 20° below horizontal. Since the jacks primarily exert forces through extension, the front, or "biceps", jack is exerting negligible force.*

To study the dynamics of this situation, we will first need a free-body diagram of the forearm (and the attached load).

pic

The forces due to the jack, the rope, and gravity should not need much explanation. However, the forces due to the pin elbow may need some explanation.

First, the free-body diagram drawn at left is of the forearm, not including the pin that serves as the elbow. Thus, the pin is external to the object of interest, and thus its interactions with the object are forces that must be indicated on the diagram. If you were to include the pin as part of the forearm, then the upper part of the arm would be external to the object of interest (pin plus forearm) and its interactions with the "pin plus forearm" would have to be indicated as forces on the diagram. Both of these approaches are completely valid.

Second, the direction of the force of the pin on the forearm may not be obvious. Does the pin push straight down on the forearm, pull up on the elbow, or push at some unspecified angle? The only thing that's clear is that this force is directed somewhere in the xy-plane. If the force is in the xy-plane, then it must have components along both the x- and y-axis (although one of these components may turn out to have a magnitude of zero). Thus, to handle the generality of the situation you should include both an x- and y-component for the force of the pin on the forearm. For convenience, I'll draw these forces as pointing in the positive x and y directions. (If my guess is wrong, the mathematics will tell me.)

Now that we have a free-body diagram, we can apply Newton's second law, both the linear form (in the x- and y-directions) and the rotational form (in the Q-direction).

#### x-direction y-direction

pic and pic

Since the forearm is held stationary, both  $a_x$  and  $a_y$  are equal to zero. Also,  $F_{\text{rope}}$  is equal to the force of gravity on the load, 882 N.

pic and pic

Obviously, we need another equation in order to determine the two unknown forces. The obvious choice is Newton's second law in rotational form.

Before we begin, we should determine the rotational inertia for a thin rod (the closest thing to a forearm in our table) rotated about an axis not at its CM. A thin rod rotated about its CM has rotational inertia  $\frac{1}{12} ML^2$ . We are interested in its rotational inertia about an axis not at its CM, so we must use the parallel-axis theorem with  $r_{\text{CM}} = 0.30$  m. (The CM is at the center of the forearm, 0.45 m from either end. Since the elbow is 0.15 m from one end, the distance between the elbow and the CM is 0.30 m.) Thus,

pic

We must also determine the torque due to each force. (Remember, the angle,  $\phi$ , in the relation for torque is the angle between  $\mathbf{r}$  (oriented along the forearm) and  $\mathbf{F}$ .)

#### **Pin, Y Jack Gravity Rope**

pic and pic and pic and pic

Finally, let's apply Newton's second law in rotational form, paying careful attention to the algebraic sign of each torque. Note that our coordinate system indicates that all torques acting counterclockwise are positive. Therefore, all torques acting clockwise are negative.

pic

Since the forearm is held stationary,  $a$  is equal to zero.

pic

The jack must exert a force of 7796 N to hold the forearm stationary.

Plugging this value into our  $y$ -equation yields:

pic

The pin must exert a force of 9070 N upwards, since the mathematics determined that the algebraic sign of  $F_{\text{pin } y}$  was positive.

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## 04. Applying Newton's Second Law in Translational and Rotational Form - II

### 1. Applying Newton's Second Law in Translational and Rotational Form - II

#### Applying Newton's Second Law in Translational and Rotational Form - II

A 75 kg man is attached to a rope wrapped around a 35 kg disk-shaped pulley, with inner and outer diameters 0.60 m and 0.90 m, respectively. The man is initially at rest. The brake shoe is pressed against the pulley with a force of 500 N. The coefficient of friction between the brake shoe and the pulley is (0.9, 0.8).

pic

First, we have to determine whether the force applied to the brake shoe is sufficient to hold the man stationary. If it's not, we have to find the man's acceleration. To accomplish this, we need free-body diagrams for both the man and the pulley.

pic

We'll use counterclockwise as the positive  $Q$ -direction and down as the positive  $y$ -direction.

The pin, or axle, at the center of the pulley exerts forces in  $x$ - and  $y$ -direction, although the exact directions may be unknown. We'll just label these forces as  $F_{\text{pin } x}$  and  $F_{\text{pin } y}$ .

The frictional force that acts on the right edge of the pulley acts in a way to prevent the man from falling (or, if he does fall, to slow down the man's fall).

To begin my analysis, I'll assume that the man is stationary, solve for the value of the frictional force required to keep him stationary, and then determine whether this frictional force is possible given the applied force of 500 N and the coefficient of static friction.

Let's apply Newton's second law to the man in the  $y$ -direction. (Remember, we are *assuming* he is not falling.)

pic

Now look at Newton's second law applied to the pulley in the  $Q$ -direction.

pic

Notice that all of the other forces acting on the pulley do *not* exert torques on the pulley. They are either located at  $r = 0$  m or have angular orientations of  $\theta = 180^\circ$ .

pic

Therefore, to hold the man stationary requires 490 N of friction. However,

pic

Therefore, the man must accelerate downward.

Now that we know the man must fall, let's write the same two equations as before, although this time the accelerations (both angular for the pulley and linear for the man) are not zero.

For the pulley:

pic

This assumes that the rotational inertia of a "compound" pulley, one with more than one location where a rope can be wrapped, is the same as a regular pulley, and that the outermost radius of the pulley determines the inertia.

Since the man is falling, the frictional force is now kinetic, and

pic

Thus our "pulley" equation becomes:

pic

Examining Newton's second law for the man,

pic

We now have two equations with three variables. However, the angular acceleration of the pulley and the linear acceleration of the man are directly related. Since the rope that the man is attached to is wrapped 0.3 m from the center of the pulley, the man accelerates at the same rate as the tangential acceleration of a point on the pulley 0.3 m from the center. Thus,

pic

Using this relationship allows me to rewrite the two Newton's second law equation as:

pic

This pair of equations has the solution

pic

The man falls with this acceleration and the rope exerts this force (upward on the man and downward on the pulley).

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## 05. Applying Newton's Second Law to Rolling Motion

### 1. Applying Newton's Second Law to Rolling Motion

#### Applying Newton's Second Law to Rolling Motion

pic

The child at the top is pretending to be a steamroller by pushing a 24 kg, 0.60 m radius cylindrical object around his backyard. The boy pushes horizontally with a force of 70 N. The coefficient of friction between the cylinder and the ground is (0.5, 0.4).

Let's draw a free-body diagram and write Newton's second law in the x-, y-, and  $\mathbf{q}$ -directions.

pic

**x-direction y-direction  $\mathbf{q}$ -direction**

pic and pic and pic

Note that the description does not specify whether the cylinder *rolls* or *skids* when the child pushes it. We will have to make an assumption, continue with the calculation, and then check our assumption for validity. Let's assume that the bottom of the cylinder *does not slip* in its contact with the ground, which means the cylinder rolls without slipping around the backyard.

If the cylinder bottom does not slip in its contact with the ground, the horizontal acceleration (and velocity) of the cylinder bottom must equal zero. Since the acceleration of the cylinder bottom is the sum of the acceleration due to translation of the CM and the acceleration due to rotation about the CM, for the bottom to have zero horizontal acceleration means that the acceleration due to rotation,

pic

must be equal in magnitude, and opposite in direction, to the translational acceleration of the CM. Thus,

pic

Combining this with our y- and  $\mathbf{q}$ -equations yields:

pic

This pair of equation has the solution

pic

Now we must check the validity of our assumption. If the cylinder rolls without slipping, the frictional force is static. Thus,

pic

Since our calculated value for  $F_{\text{friction}}$  is less than 118 N, the cylinder does remain in static contact with the ground during its motion, our assumption is validated, and our numerical results are correct.

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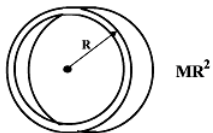
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## 06. Rotational Inertia of Common, Uniform Solids

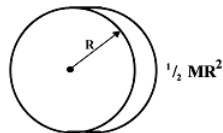
No headers

General Rule  $I = cMR^2$

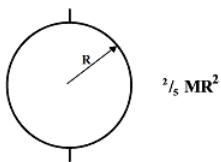
Hoop about axis through CM



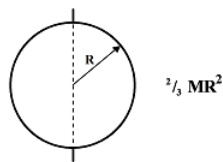
Cylinder about axis through CM



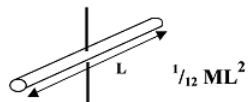
Solid sphere about axis through CM



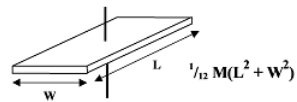
Hollow sphere about axis through CM



Thin rod about axis through CM



Rectangular plate about axis through CM



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## 07. Activities

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## 04. Conservation

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- [01. Concepts and Principles](#)
  - [02. Applying the Work-Energy Relation including Rotation - I](#)
  - [03. Applying the Impulse-Momentum Relations \(Linear and Angular\) - I](#)
  - [04. Activities](#)
- 

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## 01. Concepts and Principles

1. [The Angular Impulse-Angular Momentum Relation](#)
2. [Incorporating Rotation in the Work-Energy Relation](#)

### The Angular Impulse-Angular Momentum Relation

We've already used the impulse-momentum relation to analyze situations involving translations through three-dimensional space. The relation is typically applied in its component form:

pic 1

Arguing by analogy, if the change in momentum in the x-, y-, and z-directions is equal to the impulse applied to the object in each direction doesn't it seem plausible that a similar relation would hold in the "θ"-direction"? If so, this relation should be constructed between the torques acting on an object, the time interval over which torques act, and the change in angular velocity of an object.<sup>1</sup> The relation should be:

pic

The product of rotational inertia and angular velocity is termed the *angular momentum* of the object, typically denoted  $L$ , and the product of torque and the time interval over which it acts is termed the *angular impulse* applied to the object.

Thus, if no angular impulse is applied to an object, its angular momentum will remain constant. This special case is referred to as *angular momentum conservation*. However, if an angular impulse is applied to the object, the angular momentum will change by an amount exactly equal to the angular impulse applied. *Angular momentum is changed through angular impulse.*

### Incorporating Rotation in the Work-Energy Relation

Our previous encounter with the work-energy relation resulted in:

pic

Recall from our previous discussion of work-energy that this is *not* a vector equation, meaning it is not applied independently in each of the coordinate directions. Generalizing this equation to include rotation will involve adding terms to this equation, not creating a separate "angular" energy equation.

Recall that we model the motion of an arbitrary rigid body as a superposition of a pure translation of the CM and a pure rotation about the CM. Let's investigate the effect of this model on our calculation of the kinetic energy of the object.

The translation portion of the motion is easy. We envision the object as a point particle, localized at the CM of the real object, traveling with the velocity of the CM of the object. Thus, this portion of the motion contributes a kinetic energy,

pic

What about the kinetic energy due to the pure rotation of the object about the axis through the CM?

Every small chunk of the object,  $dm$ , moves in a circle around the CM. Each of these pieces of mass has a velocity magnitude given by

pic

Thus, the kinetic energy of each piece is

pic

Therefore, the total kinetic energy due to rotation of all the little pieces is:

pic

Combining the kinetic energy due to rotation with the kinetic energy due to translation leads to a total kinetic energy of:

pic

and a work-energy relation of:

pic

Unless a scenario involves springs or other elastic material, I'll typically write this relationship as:

pic

---

1 Please remember that this model, and hence the relationships derived under it, are restricted to rigid bodies and motions in which the rotation axis is perpendicular to the plane in which the center-of-mass moves. Since objects must be rigid, their rotational inertia must remain constant. In addition, motions must either be about a stationary rotation axis or, if not, the rotation axis is taken to be through the CM.

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## 02. Applying the Work-Energy Relation including Rotation - I

1. Applying the Work-Energy Relation including Rotation - I
2. Applying the Work-Energy Relation including Rotation - II

### Applying the Work-Energy Relation including Rotation - I

*A mischievous child releases his mother's bowling ball from the top of the family's 25 m long, 15° above horizontal driveway. The ball rolls without slipping down the driveway and at the bottom plows into the mailbox. The 6.4 kg ball has a diameter of 24 cm.*

Let's determine the speed of the ball when it hits the mailbox. To determine this value, we can apply the work-energy relation to the ball between:

Event 1: The instant the ball is released.

Event 2: The instant the ball hits the mailbox.

For these two events, work-energy looks like this:

pic and pic

where the bottom of the driveway is the zero for gravitational potential energy and the rotational inertia of the bowling ball is taken to be that of a sphere.

Now we must carefully determine which, if any, of the forces on the bowling ball do work on the bowling ball.

First, we don't have to worry about the force of gravity. The gravitational potential energy function was developed to automatically incorporate the work done by the force of gravity.

Second, the contact force can do no work on the ball because the contact force is always perpendicular to the motion of the ball.

Finally, what about the force of friction? It does appear, at first glance, that the force of friction is applied over the entire motion of the ball down the driveway. However, let's pay closer attention to the actual point at which the force acts.

The frictional force is static in nature, because since the ball rolls without slipping down the driveway the bottom of the ball is always in static contact with the ground (i.e., there is no relative velocity between the bottom of the ball and the ground). If the bottom of the ball has a velocity of zero, then the force that acts on the bottom of the ball (static friction) can act through no distance. During the instant at which the frictional force acts on a particular point on the bottom of the ball, that point is not moving. That point on the ball only moves when it is no longer in contact with the ground, but by that time the frictional force is acting on a *different* point. To summarize (and stop saying the same thing over and over), *the force of static friction can do no work because it acts on a point that does not move.*

Thus, our equation simplifies to:

pic

The linear and angular velocity of the ball must be related. Since the ball rolls without slipping, the bottom of the ball has no linear velocity. Since the velocity of the bottom of the ball is the sum of velocity due to translation of the CM and the velocity due to rotation about the CM, the velocity due to rotation must be equal in magnitude to the velocity of the CM:

pic

Substituting this into our equation yields:

pic

### Applying the Work-Energy Relation including Rotation - II

pic

The crate is descending with speed  $v_0$  when a brake shoe is applied to the disk-shaped pulley of mass  $M$  and radius  $R$ . The coefficient of friction between the shoe and the pulley is  $(\mu_s, \mu_k)$ . Determine the distance ( $D$ ) the crate moves before stopping as a function of  $M$ ,  $m$ ,  $F$ ,  $R$ ,  $v_0$ ,  $g$ , and the appropriate coefficient of friction.

To determine this function, let's apply the work-energy relation to both the crate and the pulley between:

Event 1: The instant the brake shoe is applied.

Event 2: The instant the crate comes to rest.

In addition to defining the two instants of interest, we'll need free-body diagrams for both the crate and the pulley.

### ***Pulley***

pic and pic

To find the work on the pulley, note that  $F_{\text{contact}}$  and  $F_{\text{pin}}$  do no work since the distance over which these forces act is zero. However, both  $F_{\text{rope}}$  and  $F_{\text{friction}}$  do work.

If the crate falls a distance  $D$ , both of these forces act over a distance  $D$ . However, note that this displacement is in the same direction as  $F_{\text{rope}}$  ( $f = 180$ ) on the right side of the pulley but in the opposite direction on the other side of the pulley (opposite to  $F_{\text{friction}}$  and therefore  $f = 180$ .) Since these forces are constant, there's no need to actually integrate:

pic

Note that since the brake shoe does not accelerate, the external force applied to the shoe,  $F$ , is the same magnitude as the contact force between the shoe and the pulley.

pic

### ***Crate***

pic

where the final position of the crate is the zero for gravitational potential energy.

The two equations can be added together to yield:

pic

Thus, in order for  $D$  to have a physical (positive) value, the frictional force on the pulley must be greater in magnitude than the force of gravity on the crate, which agrees with common sense. Note that the numerator is simply the kinetic energy of the pulley plus crate. The larger this sum, the larger the distance needed to stop the crate's fall, which again agrees with common sense.

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### 03. Applying the Impulse-Momentum Relations (Linear and Angular) - I

1. Applying the Impulse-Momentum relations (Linear and Angular) - I
2. Applying the Impulse-Momentum Relations (Linear and Angular) - II

#### Applying the Impulse-Momentum relations (Linear and Angular) - I

*A bowling ball of mass  $M$  and radius  $R$  leaves a bowler's hand with CM velocity  $v_0$  and no rotation. The ball skids down the alley before it begins to roll without slipping. The coefficient of friction is  $(m_s, m_k)$ . Determine the elapsed time ( $T$ ) before the ball begins to roll without slipping as a function of  $M$ ,  $R$ ,  $v_0$ ,  $g$ , and the appropriate coefficient of friction.*

To determine this function, let's apply the impulse-momentum relations to the ball between:

Event 1: The instant the ball leaves the bowler's hand.

Event 2: The instant the ball begins to roll without slipping.

pic

Since the ball both translates and rotates, we must write both the linear and rotational forms of the impulse-momentum relation. Remember, the rotation of the bowling ball is modeled to be about an axis through the CM. (Since all the forces are constant, we'll write the relations without the use of the integral.)

**x-direction (linear momentum)**,  **$\phi$ -direction (angular momentum)**

pic and pic

The final CM velocity and final angular velocity are related because at this instant the ball begins to roll without slipping. When the ball rolls without slipping, the bottom of the ball is in static contact with the ground, i.e., it has no linear velocity. Since the velocity of the bottom of the ball is the sum of the velocity due to translation of the CM and the velocity due to rotation about the CM, the velocity due to rotation must be equal in magnitude to the velocity of the CM:

pic

Substituting this into our angular equation yields:

pic

Plugging this into the linear equation:

pic

Thus, the faster you throw the ball or the smaller the kinetic coefficient of friction, the longer it will take for the ball to begin to roll.

#### Applying the Impulse-Momentum Relations (Linear and Angular) - II

*A 30 kg child running at 4.0 m/s leaps onto the outer edge of an initially stationary merry-go-round. The merry-go-round is a flat disk of radius 2.4 m and mass 70 kg. Ignore the frictional torque in the bearings of the merry-go-round.*

Let's imagine we're interested in determining the final angular velocity of the merry-go-round after the child has safely come to rest on its surface. To determine this angular velocity, we can apply the impulse-momentum relations to both the child and the merry-go-round between:

Event 1: The instant before the child lands on the merry-go-round

Event 2: The instant after the child comes to rest on the merry-go-round.

pic

The child and the merry-go-round interact via some unknown magnitude force. (Contrary to the free-body diagram, this force does not necessarily act on the child's disembodied head.) There are also other forces due to gravity and supporting structures that act on both the child and the merry-go-round, however, these forces are in the vertical direction and supply no torque.

### *Child Merry-Go-Round*

#### **x-direction q-direction**

pic and pic

The final velocity of the child and final angular velocity of the merry-go-round are related because at this instant the child is at rest (hanging on to) the merry-go-round. Therefore,

pic

Substituting this into our angular equation yields:

pic

By Newton's Third Law, the force of the child on the merry-go-round and the force of the merry-go-round on the child must be equal in magnitude. Therefore our two equations can be summed and the impulses cancel. This gives:

pic

The child is slowed from 4 m/s to 1.85 m/s by jumping onto the merry-go-round. The merry-go-round, however, is accelerated from rest to:

pic

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## 04. Activities

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## CHAPTER OVERVIEW

### Model 05 - Electric Fields

#### Topic hierarchy

- 01. Concepts and Principles
- 02. Concepts and Principles 2
- 03. Analysis Tools
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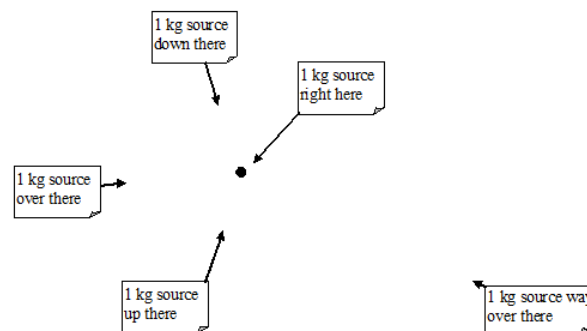
## 01. Concepts and Principles

### 1. Electric Charge, Electric Field and a Goofy Analogy

#### Electric Charge, Electric Field and a Goofy Analogy

We all know that electrons and protons have electric charge. But what *is* electric charge and what does it mean for a particle, like a neutron, to not have electric charge[1]? On one level, the answer is that electric charge is the ability to create and interact with an electric field. Of course, this begs the question, what is an electric field? To try to answer this question, let's look at mass and the gravitational field.

In Newton's theory of gravitation, every object that has mass creates a gravitational field. The object with mass is termed the *source* of the gravitational field. The source's gravitational field, which fills all of space, encodes information as to the location and mass of the source into space itself. It's almost as if an infinite number of little business cards have been printed and distributed throughout space with detailed information concerning the source's characteristics. The gravitational field can be thought of as a huge number of business cards embedded into the fabric of space. (No, I am not joking, it looks like this.)



Typically when we draw the gravitational field we translate the message on the business cards into a mathematically equivalent message. Here's a picture of the gravitational field near the surface of the earth.

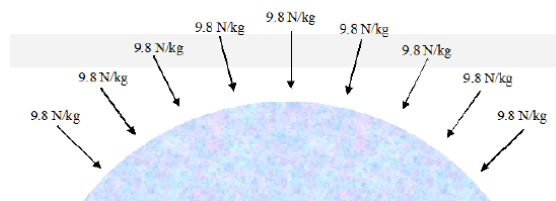


Figure 2:

The relationship that allows you to "translate the business cards" is Newton's expression for the gravitational field surrounding a mass,

$$g = \frac{Gm}{r^2} \quad (1)$$

where

- $G$  is the *gravitational constant*, equal to  $6.67 \times 10^{-11} \text{ N m}^2/\text{kg}^2$ .
- $m$  is the source mass, the mass that creates the gravitational field.
- $r$  is the distance between the source mass and the location of the business card.
- $(\hat{r})$  is the unit vector that points *from* the source mass *to* the business card. Notice that the direction of the gravitational field,  $g$ , is opposite to this direction.

Every mass in the universe has a gravitational field described by this formula surrounding it. Additionally, every mass in the universe has the ability to "read the business cards" of *every other mass*. This transfer of information from one mass to another is

the gravitational force that attracts masses together. (Once one mass reads another masses' business card, the first mass feels a strange urge to go visit the second mass. The more enticing the business card (the larger the value of  $g$ ) the stronger the urge.)

Let's drop the business card analogy and return to electric charges and electric fields. Basically, every object that has electric charge surrounds itself with an electric field given by a formula incredibly similar to the one for the gravitational field. How incredibly similar? Below is the formula for the electric field at a particular point in space, from a single source charge:

$$E = \frac{kq}{r^2} \quad (2)$$

where

- $k$  is the *electrostatic constant*, equal to  $9.00 \times 10^9 \text{ N m}^2/\text{C}^2$ .
- $q$  is the source charge, the electric charge that creates the electric field. Its value can be positive or negative, and is measured in coulombs (C).
- $r$  is the distance between the source charge and the point of interest.
- $\hat{r}$  is the unit vector that points *from* the source charge *to* the point of interest.

Every charge object is surrounded by a field given by this relationship. Every *other* charged object in the universe can "read" this field and will respond to its information by feeling an electric force. Object without electric charge neither create nor interact with electric fields (they can't read the business cards).

In this chapter you'll learn how to calculate the electric field produced by charge objects. In the next chapter, you'll learn how the electric field can be "sensed" by other electric charges resulting in the electric force.

---

[1] Actually, neutrons are constructed from smaller particles called quarks that do have electric charge. Neutrons are neutral because the sum of the quark charges inside of them is zero. The distinction between an object having a *net* charge of zero and an object having *no* charge whatsoever is important when considering polarization.

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## 02. Concepts and Principles 2

1. Charge and Charge Density
2. Perfect Conductors and Perfect Insulators

### Charge and Charge Density

Macroscopic objects are normally neutral (or very close to neutral) because they contain equal numbers of protons and electrons. All charged objects are charged because of either an excess or lack of electrons. (It's much easier to add or remove electrons from an object than trying to add or remove the protons tightly bound inside the nuclei of its atoms.) Thus, the electric charge of any object is always an integer multiple of the electric charge on an electron.

Because of its fundamental importance, the magnitude of the charge on an electron is termed the *elementary charge* and denoted by the symbol  $e$ . In a purely logical world, the charge on any object would be reported as a multiple of  $e$ . However, since the charge on a macroscopic system can be *many* multiples of  $e$ , a more user-friendly unit, the coulomb (C), typically used to quantify electric charge. In this system,

$$e = 1.6 \times 10^{-19} \text{ C} \quad (1)$$

Thus, you can consider the charge on an electron as an incredibly small fraction of a coulomb, or a coulomb of charge as an incredibly large number of electrons.

In many applications, in addition to knowing the total charge on an object you will need to know how the charge is distributed. The distribution of charge on an object can be defined in several different ways. For objects such as wires or other thin cylinders, a *linear charge density*,  $\lambda$ , will often be defined. This is the amount of charge per unit length of the object. If the charge is uniformly distributed, this is simply

$$\lambda = \frac{Q}{L} \quad (2)$$

where  $Q$  is the total charge on the object[1] and  $L$  its total length. However, if the charge density varies over the length of the object, its value at any point must be defined as the ratio of the charge on a differential element at that location to the length of the element:

$$\lambda(x) = \frac{dQ}{dx} \quad (3)$$

For objects such as flat plates or the surfaces of cylinders and spheres, a *surface charge density*,  $\sigma$ , can be defined. This is the amount of charge per unit area of the object. If the charge is uniformly distributed, this is

$$\sigma = \frac{Q}{\text{Area}} \quad (4)$$

or if the charge density varies over the surface:

$$\sigma(x, y) = \frac{dQ(x, y)}{dxdy} \quad (5)$$

Lastly, for objects that have charge distributed throughout their volume, a *volume charge density*,  $\rho$ , can be defined. This is the amount of charge per unit volume of the object. If the charge is uniformly distributed, this is

$$\rho = \frac{Q}{V} \quad (6)$$

or if the charge density varies inside the object:

$$\rho = \frac{dQ}{dV} \quad (7)$$

To add to the confusion, you must realize that the same object can be described as having two different charge densities. For example, consider a plastic rod with charge distributed throughout its volume. Obviously, the charge per unit volume,  $\rho$ , can be

defined for this object. However, you can also define the object as having linear charge density,  $\lambda$ , reporting the amount of charge present per meter of length. These two parameters will have different values but refer to exactly the same object.

## Perfect Conductors and Perfect Insulators

Determining how electric charges in real materials respond to electric fields is incredibly important but also incredibly complicated. In light of this, we will restrict ourselves to two types of hypothetical materials.

In a perfect conductor, electric charges are free to move without any resistance to their motion. Metals provide a reasonable approximation of perfect conductors, although, of course, in a real metal a small amount of resistance to motion is present. When I refer to a material as a metal, we will approximate the metal as a perfect conductor.

In a perfect insulator electric charges can not move, regardless of the amount of force applied to them. Many materials act as insulators, but all real materials experience electrical breakdown if the forces acting on charges become so great that the charges begin to move. When I refer to an insulating material, like plastic, for example, we will approximate the material as a perfect insulator.

Since electric fields create forces on electric charges, there can not be static electric fields present inside perfect conductors. If a field was present inside a perfect conductor, the charges inside the conductor would feel an electric force and hence move in response to that force. They would continue to move until they redistributed themselves inside of the conductor in such a way as to cancel the electric field. The system could not reach equilibrium as long as an electric field was present. This re-arranging process would typically occur *very* quickly and we will always assume our analysis takes place after it is completed.

---

[1] I will use lowercase  $q$  to designate the charge on a point particle and uppercase  $Q$  to designate the total charge distributed on macroscopic objects.

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## 03. Analysis Tools

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### 1. Point Charges

#### Point Charges

Find the electric field at the indicated point. The charges are separated by a distance  $4a$ .

pic

The electric field at this point will be the vector sum of the electric field from the left charge (special  $e_1$ ) and the electric field from the right charge (special  $e_2$ ).

Let's examine the left charge first. Since the problem is expressed symbolically, the charge is simply " $2q$ ", and the distance,  $r$ , between the charge and the indicated point of interest should be expressed in terms of " $a$ ". Since each square in the diagram has width and height  $a$ , this distance can be expressed as:

pic

All that's to determine is the unit vector. The unit vector is simply a mathematical description of how to get from the source charge to the point of interest. In English, to get from the source charge to the point of interest you should move  $3a$  in the  $x$ -direction and  $a$  in the  $y$ -direction<sup>[1]</sup>. This can be written as:

pic

This is the vector that points from the source charge to point of interest. However, this is not a *unit* vector that points from the source charge to the point of interest. However, this is not a *unit* vector since its magnitude isn't 1. (A unit vector should convey a *direction* in space without altering the *magnitude* of the rest of the equation. It can accomplish this only if its magnitude is equal to 1.)

However, it's simple to convert a regular vector into a unit vector, just divide the vector by its magnitude. This leads to:

pic

Putting this all together yields:

pic

Repeating for the right charge gives:

pic

Adding these two contributions together yields

pic

Thus, the electric field at the location indicated points to the right and slightly downward.

---

<sup>[1]</sup>For the sake of consistency, we will use a common coordinate system with the  $+x$ -direction pointing to the right and the  $+y$ -direction pointing to the top of the page. For three dimensional systems, the  $+z$ -direction will point directly out of the page. The coordinate axes will be indicated in the vast majority of diagrams.

## 04. Analysis Tools 2

### 1. Continuous Charge Distribution

#### Continuous Charge Distribution

The plastic rod of length  $L$  below has uniform charge density  $\lambda$ . Find the electric field at all points on the  $x$ -axis to the right of the rod.

pic

Electric charge is discrete, meaning it comes in integer multiples of electron and proton charge. Therefore, the electric field can always be calculated by summing the electric field from each of the electrons and protons that make up an object. However, macroscopic objects contain a *lot* of electrons and protons, so this summation has many, many terms:

pic

As described earlier, we will assume that the charge on macroscopic objects is continuous, and distributed throughout the object. Mathematically, this means we will replace a summation over a very large number of discrete charges with an integral over a hypothetically continuous distribution of charge. This leads to a relationship for the electric field at a particular point in space, from a continuous distribution of charge, of:

pic

where  $dq$  is the charge on an infinitesimally small portion of the object, and the integral is over the entire physical object.

Finding the electric field from a continuous distribution of charge involves several distinct steps. Until you become very comfortable setting up and evaluating electric field integrals, I would suggest you systematically walk through these steps.

1. Carefully identify and label the location of the differential element on a diagram of the situation.
2. Carefully identify and label the location of the point of interest on a diagram of the situation.
3. Write an expression for  $dq$ , the charge on the differential element.
4. Write an expression for  $r$ , the distance between the differential element and the point of interest.
5. Write an expression for  $(\hat{r})$ , the unit vector representing the direction from the element to the point of interest.
6. Insert your expressions into the integral for the electric field.
7. Carefully choose the limits of integration.
8. Evaluate the integral.

I'll demonstrate below each of these steps for the scenario under investigation.

1. Carefully identify and label the location of the differential element on a diagram of the situation.

The differential element is a small (infinitesimal) piece of the object that we will treat like a point charge. The location of this differential element must be *arbitrary*, meaning it is not at a "special" location like the top, middle, or bottom of the rod. Its location must be represented by a variable, where this variable is the variable of integration and determines the limits of the integral.

For this example, select the differential element to be located a distance " $y$ " above the center of the rod. The length of this element is " $dy$ ". (Later, you will select the limits of integration to go from  $-L/2$  to  $+L/2$  to allow this arbitrary element to "cover" the entire rod.)

pic

2. Carefully identify and label the location of the point of interest on a diagram of the situation.

You are interested in the electric field at all points along the  $x$ -axis to the right of the rod. Therefore, select an *arbitrary* location along the  $x$ -axis and label it with its location.

pic

3. Write an expression for  $dq$ , the charge on the differential element.

The rod has a uniform charge density  $\lambda$ , meaning the amount of charge per unit length along the rod is constant. Since the differential element has a length  $dy$ , the total charge on this element ( $dq$ ) is the product of the density and the length:

pic

4. Write an expression for  $r$ , the distance between the differential element and the point of interest.

By Pythagoras' theorem, the distance between the differential element and the point of interest is:

pic

5. Write an expression for  $(\hat{r})$ , the unit vector representing the direction from the element to the point of interest.

The vector going from the element to the point of interest heads "down" a distance  $y$  and then to the right a distance  $x$ . This leads to a unit vector:

pic

6. Insert your expressions into the integral for the electric field.

pic

7. Carefully choose the limits of integration.

The limits of integration are determined by the range over which the differential element must be "moved" to cover the entire object. The location of the element must vary between the bottom of the rod  $(-L/2)$  and the top of the rod  $(+L/2)$  in order to include every part of the rod. The two ends of the rod form the two limits of integration.

pic

8. Evaluate the integral.

pic

It's typically easier to think of a vector integral as two, separate scalar integrals, one "in the x-direction" and one "in the y-direction", as above.

Examining the x-integral (and using an integral table):

pic

I'll leave it to you to evaluate the y-integral, but you should find that it equals zero. (Without doing calculation whatsoever, you should *know* that the y-component of the electric field along the x-axis *must* be zero because of the symmetry of the situation. If it's not clear why the y-component of the field must be zero, talk to your instructor!)

Thus, the electric field along the x-axis is given by:

pic

Does this function make sense? Two ways to check whether this field is reasonable are to determine how the function behaves for very small and very large values of  $x$ .

As  $x$  gets very small, you are getting closer and closer to the charged rod. This *should* lead to an electric field that grows larger and larger (without bound). Notice that the limit of the function as  $x$  approaches zero does "go to infinity", so the function does have the proper behavior for small  $x$ .

As  $x$  gets very large, you are getting farther and farther to the charged rod. Not only should the field decrease, but as you get very far from the rod should begin to look like a point charge. Notice that as  $x$  gets large, the term  $(L^2/4)$  becomes negligible compared to  $x^2$ . This leads to a field

pic

Since the total charge on the rod ( $Q$ ) is simply the product of the charge density and the total length of the rod, this reduces to

pic

This is exactly the expression for the electric field from a point charge. Thus, as you move farther and farther from the rod, the rod does indeed begin to look like a point charge.

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## 05. Analysis Tools 3

### 1. Gauss' Law

### Gauss' Law

The long, hollow plastic cylinder at right has inner radius  $a$ , outer radius  $b$ , and uniform charge density  $r$ . Find the electric field at all points in a plane perpendicular to the cylinder near its midpoint.

pic

For certain situations, typically ones with a high degree of symmetry, Gauss' Law allows you to calculate the electric field relatively easily. Gauss' Law, mathematically, states:

pic

Let's describe what this means in English. The left side of the equation involves the vector dot product between the electric field and an infinitesimally small area that is a piece of a larger closed surface (termed the *gaussian surface*). This dot product between electric field and area is termed *electric flux*, and is often visualized as the amount of field that "passes through" the little piece of area. The integral simply tells us to add up all of these infinitesimal electric fluxes to get the total flux through the entire closed surface.

The gist of Gauss' Law is that this total electric flux is exactly equal to the total amount of electric charge enclosed within the gaussian surface, divided by a constant,  $\epsilon_0$ . ( $\epsilon_0$  is the *permittivity of free space*, a constant equal to  $8.85 \times 10^{-12} \text{ C}^2/\text{Nm}^2$ .)

Somewhat counter intuitively, the key to applying Gauss' Law is to choose a gaussian surface such that you never really have to do the integral on the left side of the equation! To try to help you understand what I'm talking about, let's walk through the solution of the above problem. The following sequence of steps will help you understand the process of applying Gauss's law:

1. Choose the appropriate gaussian surface.
2. Carefully draw the hypothetical gaussian surface at the location of interest.
3. Carefully draw the electric field at all points on the gaussian surface.
4. Write an expression for the surface area parallel to the electric field.
5. Write an expression for  $q_{\text{enclosed}}$ , the charge enclosed within the gaussian surface.
6. Apply Gauss' Law and determine the electric field at all points on this hypothetical surface.

Since we want to find the electric field at all points in space, there are three distinct regions we will have to investigate:

- the region inside the "hole" in the cylinder ( $r < a$ ),
- the region within the actual material of the cylinder ( $a < r < b$ ),
- and the region outside of the cylinder ( $r > b$ ).

Let's start outside of the cylinder.

#### Outside of the cylinder: $r > b$

1. Choose the appropriate gaussian surface.

Gauss' law is primarily useful when the objects under investigation have a high degree of symmetry. Gauss' law basically exploits that symmetry to make the calculation of the electric field relatively painless. To exploit the symmetry of the situation, you should always choose a gaussian surface to mimic the symmetry of the object you are investigating. In this example, since the object is a cylinder, my gaussian surfaces will all be cylinders.

2. Carefully draw the hypothetical gaussian surface at the location of interest.

Since we are trying to determine the electric field for all points outside of the cylinder, draw a cylindrical gaussian surface with radius  $r$ . The value of  $r$  is *variable*, and can take on any value greater than  $b$ , the radius of the *real* cylinder. Remember, the gaussian surface is *hypothetical*; it's a mathematical "object" that only exists to help you solve the problem. Try not to confuse it with the real cylinder of radius  $b$ .

The gaussian surface is represented from two separate viewpoints below. (The gray cylinder is the actual, charged cylinder while the dashed cylinder is the gaussian surface.) The gaussian surface has radius  $r$ , where  $r$  can be any value greater than  $b$ , and length

L. It is located near the midpoint of the actual cylinder. The gaussian surface also includes the circular "end caps" of the cylinder since a gaussian surface must be a *closed* surface.

pic

3. *Carefully draw the electric field at all points on the gaussian surface.*

Although I have no idea what the *magnitude* of the electric field is at any point on my gaussian surface, the symmetry of the situation tells me that the *direction* of the electric field must be either radially away from or radially toward the central axis of the cylinder. If the charge density is positive, the field will be directed radially outward from the cylinder axis.

Moreover, even though I don't know the magnitude of the field, I *do* know that the magnitude is the same at every point on my surface.

pic

4. *Write an expression for the surface area parallel to the electric field.*

The left side of Gauss' law requires us to evaluate an integral over the surface of our gaussian cylinder. The integral requires us to find the dot product between the electric field and the area, and then integrate this dot product over the entire surface. I mentioned earlier that you should never have to actually do this integral (assuming you chose the "correct" gaussian surface). So why don't we have to do this integral?

The vector dot product can be re-written as:

pic

where  $\theta$  is the angle between the electric field and the area of the gaussian surface. You may recall from calculus that the "direction" associated with area is perpendicular to its surface. Thus, the direction of each infinitesimal area is indicated on the diagrams below.

pic

Notice that the electric field vector and the area vector are parallel at every point on the cylindrical portion of the gaussian surface and perpendicular at every point on the circular end caps. Thus, breaking the integral into two parts, one over the cylindrical portion of the gaussian surface and one over the end caps, yields:

pic

Now note that the magnitude of the electric field is the same at every point on the cylindrical portion of the gaussian surface since every point is equal distance from the charge distribution. Thus, the electric field is constant and can be brought outside of the integral, leaving a pretty easy integral to evaluate.

pic

Notice that the entire left-hand side of Gauss' law reduces to the product of the electric field magnitude and the magnitude of the surface area parallel to this field. Thus, because of our wise choice of gaussian surface, all we really need to calculate is the magnitude of the surface area parallel to the electric field. In this case, the parallel area is given by:

pic

5. *Write an expression for  $q_{\text{enclosed}}$  the charge enclosed within the gaussian surface.*

Since the gaussian surface is outside of the real cylinder, *all* of the charge on the cylinder within the length  $L$  is enclosed. Since we know the volume charge density on the cylinder,  $\rho$ , the total charge on the cylinder within the length  $L$  is the product of the charge density and the volume of the cylinder.

pic

6. *Apply Gauss' Law and determine the electric field at all points on this hypothetical surface.*

pic

Thus, the electric field outside of the cylinder is inversely proportional to the distance from the cylinder.

Now we have to repeat this analysis for the other two regions.

**Within the cylinder:  $a < r < b$**

1. *Choose the appropriate gaussian surface.*

We will again use a cylindrical surface.

2. *Carefully draw the hypothetical gaussian surface at the location of interest.*

Since we are interested in the electric field within the actual cylinder, the radius of our gaussian surface is larger than  $a$  but less than  $b$ .

pic

3. *Carefully draw the electric field at all points on the gaussian surface.*

As before, the magnitude of the electric field must be constant on the gaussian surface and directed radially outward due to the symmetry of the situation.

pic

4. *Write an expression for the surface area parallel to the electric field.*

The area parallel to the electric field is again the surface area of the cylindrical portion of the gaussian surface:

pic

5. *Write an expression for  $q_{\text{enclosed}}$  the charge enclosed within the gaussian surface.*

Since the gaussian surface is inside the real cylinder, not all of the charge on the real cylinder is enclosed by the gaussian surface. All of the charge between  $a$  and  $r$  is enclosed, but the charge between  $r$  and  $b$  is not enclosed by the gaussian surface. Thus, the amount enclosed is the product of the volume charge density and the volume of the portion of the cylinder enclosed by the gaussian surface.

pic

6. *Apply Gauss' Law and determine the electric field at all points on this hypothetical surface.*

pic

Notice that this electric field increases with increasing  $r$ , since as  $r$  increases, more and more charge is available to produce the electric field.

The last region we have to investigate is inside of the "hole" in the cylinder.

**Inside the "hole":  $r < a$**

Since we must choose our gaussian surface to have a radius less than  $a$ , it is located inside the hollow center of the cylinder. Since there is *no* charge enclosed by this gaussian surface, the electric field in this region must be zero.

pic

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## 06. Activities

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## CHAPTER OVERVIEW

### Model 06 - Electric Forces

#### Topic hierarchy

- 01. Concepts and Principles
- 02. Analysis Tools
- 03. Activities

*Thumbnail: A generator with a single rectangular coil rotated at constant angular velocity in a uniform magnetic field produces an emf that varies sinusoidally in time. Note the generator is similar to a motor, except the shaft is rotated to produce a current rather than the other way around. Image used with permission (CC BY; OpenStax).*

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## 01. Concepts and Principles

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### 1. The Gravitational Analogy

#### The Gravitational Analogy

In introducing the concept of the electric field, I tried to illustrate it by drawing an analogy with the gravitational field, ( $\vec{g}$ ). This analogy can be extended to electric force and gravitational force.

From mechanics, the relationship for the gravitational force on an object is:

$F_g = m\vec{g}$

where

- $m$  is the mass of the particle of interest,
- and  $\vec{g}$  is the net gravitational field at the location of the particle of interest. (This field was typically approximated as that of the earth, but should actually be the field created by all of the massive particles in the universe, other than the particle of interest.)

Mass is the property that allows particles to *create* gravitational fields and it is also the property that allow them to *interact* with other particles' gravitational fields. This interaction is termed the gravitational force.

Wouldn't it be great if charge played the same role with regard to the electric field? Well, it is great! The electric force on a particle is given by the relation,

$F_e = q\vec{E}$

where

- $q$  is the charge on the particle of interest,
- and  $\vec{E}$  is the net electric field at the location of the particle of interest (created by all of the *other* charged particles in the universe).

All charged particles create electric fields, but this is only half of the story. All charged particles also interact with *other* particles' electric fields. This interaction is termed the electric force.

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## 02. Analysis Tools

1. Point Charges
2. Force and Motion

### Point Charges

The three charges below are located as shown. Each grid square has width  $a$ . Find the net electric force on the positive charge.

pic

To find the electric force on the positive charge, you must first find the electric field at its location in space. Thus, this is really just more practice in calculating the electric field!

The electric field at the location of the positive charge will be the vector sum of the electric field from the  $-2q$  charge ( $E_{-2Q}$ ) and the electric field from the  $-2q$  charge ( $E_{-2q}$ ). Let's calculate these two fields separately and then add them together.

From the  $-2q$  charge:

pic

and from the  $+2q$  charge:

pic

Adding these two contributions together yields

pic

So the force on the positive charge is:

pic

The force on the positive charge is to the left, and slightly downward, attracted to the two negative charges.

### Force and Motion

In many applications, oppositely charged parallel plates are used to "steer" beams of charged particles. In this example, a proton is injected at  $2.0 \times 10^6$  m/s into the space between the plates. The plates are 2.0 cm long. What charge density is needed on the plates to give proton a  $y$ -velocity of  $2.0 \times 10^5$  m/s as it exits the plates?

pic

Since this problem involves the motion of a particle between two distinct events, let's complete a motion table.

Event 1: The proton enters the device. Event 2: The proton exits the device.

$$t_1 = 0 \text{ s } t_2 =$$

$$r_{1x} = 0 \text{ m } r_{2x} = 0.02 \text{ m}$$

$$r_{1y} = 0 \text{ m } r_{2y} =$$

$$v_{1x} = 2 \times 10^6 \text{ m/s } v_{2x} =$$

$$v_{1y} = 0 \text{ m/s } v_{2y} = 2 \times 10^5 \text{ m/s}$$

Between the plates, the proton will experience an electrical force, and hence acceleration, in the  $y$ -direction. Remember from mechanics that this  $y$ -acceleration will not affect the kinematics of the proton in the  $x$ -direction. Thus, in the  $x$ -direction the acceleration of the proton is zero.

Applying the kinematic equations in the  $x$ -direction yields:

pic

Applying the same kinematics equations in the  $y$ -direction yields:

pic

Now, using Newton's Second Law and the relation for electric field from parallel conducting plates, find the necessary charge density.

pic

The bottom plate should be made positive and the top plate negative, both with this charge density. (Note that the force of gravity acting on the proton is completely insignificant compared to the electrical force. This is generally true and we will typically ignore the force of gravity acting on individual particles such as protons and electrons.)

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### 03. Activities

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## CHAPTER OVERVIEW

### Model 07 - Magnetic Fields

- 01. Concepts and Principles
- 02. Analysis Tools
- 03. Analysis Tools 2
- 04. Activities

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## 01. Concepts and Principles

1. Moving Charges
2. Permanent Magnets
3. Electric Current

### Moving Charges

All charged particles create electric fields, and these fields can be detected by other charged particles resulting in electric force. However, a completely *different* field, both qualitatively and quantitatively, is created when charged particles move. This is the magnetic field. All *moving* charged particles create magnetic fields, and all *moving* charged particles can detect magnetic fields resulting in magnetic force. *This is in addition to the electric field that is always present surrounding charged particles.*

This should strike you as rather strange. Whenever a charged particle begins to move a completely new field springs into existence (distributing business cards throughout the universe). Other charged particles, if they are at rest relative to this new field, *do not notice this new field* and do not feel a magnetic force. Only if they move relative to this new field can they sense its existence and feel a magnetic force. It's as if only while in motion can they read the business cards distributed by the original moving charge, and only while in motion does the original charge distribute these business cards in the first place! Does it sound strange yet?

Why the magnetic field exists, and its relationship to the electric field and relative motion will be explored later in the course. For now, we will concentrate on learning how to calculate the value of the magnetic field at various points surrounding moving charges. Next chapter, we will learn how to calculate the value of the magnetic force acting on *other* charges moving relative to a magnetic field.

### Permanent Magnets

I claimed above that the magnetic field only exists when the source charges that create it are moving. But what about permanent magnets, like the ones holding your favorite physics assignments to your refrigerator? Where are the moving charges in those magnets?

The simplest answer is that the electrons in "orbit" in each of the atoms of the material create magnetic fields. In most materials, these microscopic magnetic fields are oriented in random directions and therefore cancel out when summed over all of the atoms in the material. In some materials, however, these microscopic magnetic field (large enough to interact with the microscopic magnetic fields present in your refrigerator door). Although this is a gross simplification of what actually takes place, it's good enough for now.

The magnetic properties of real materials are extremely complicated. In addition to the orbital contribution to magnetic field, individual electrons and protons have an intrinsic magnetic field associated with them due to a property called *spin*. Moreover, even neutrons, *with no net electric charge*, have an intrinsic magnetic field surrounding them. To learn more about the microscopic basis of magnetism, consider becoming a physics major...

### Electric Current

Moving electric charges form an electric current. We will consider the source of all magnetic fields to be electric current, whether that current is macroscopic and flows through a wire or whether it is microscopic and flows "in orbit" around an atomic nucleus.

The simplest source of magnetic field is electric current flowing through a long, straight wire. In this case, the magnetic field at a particular point in space is given by the relation,

pic

where

- $\mu_0$  is the *permeability of free space*, a constant equal to  $1.26 \times 10^{-6} \text{ Tm/A}$ .
- $i$  is the source current, the electric current that creates the magnetic field, and is measured in amperes (A). One ampere is equal to one coulomb of charge flowign through the wire per second.(We will always enter the current as a positive value in this relation. The definition of ( $\hat{t}$  hat) below will be used to specify the correct direction of the field.)
- $r$  is the distance between the source current and the point of interest.

- $\hat{t}$  is the unit vector that is *tangent* to a circle centered on the source current, and located at the point of interest. To determine the sense of  $\hat{t}$ , place your thumb in the direction of current flow. The sense in which the fingers of your *right* hand curl is the sense of  $\hat{t}$ . (For example, for current flowing out of the page,  $\hat{t}$  is counterclockwise.)

If the source of the magnetic field is more complicated than long, straight current-carrying wires, then a more general relationship is needed. The magnetic field at a particular point in space, from any current distribution, is given by the relation,

$\vec{B}$

where

- $d\vec{r}$  is an infinitesimally small portion of current-carrying wire.
- $\hat{r}$  is the unit vector that points *from* the small portion of current-carrying wire *to* the point of interest.
- The integral is over the entire length of wire.

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## 02. Analysis Tools

1. Long, Parallel Wires
2. More General Current Distribution

### Long, Parallel Wires

Find the magnetic field at the indicated point. The long, parallel wires are separated by a distance  $4a$ .

The magnetic field at this point will be the vector sum of the magnetic field from the left wire ( $\mathbf{B}_L$ ) and the magnetic field from the right wire ( $\mathbf{B}_R$ ).

For the left wire, I've indicated the direction of the tangent vector. Remember, with your thumb pointing in the direction of the current (out of the page), the direction in which the fingers of your right hand curl is the direction of the tangent vector (counterclockwise).

pic

By definition, the tangent vector is perpendicular to the radial vector. You should know how to specify the radial unit vector from the chapter on electric field; to specify the tangent unit vector, you have to construct a vector *perpendicular* to the radial vector.

To do this, simply flip the x- and y-components of the radial vector and add the appropriate algebraic signs. (Ask your math professor to prove that this always results in a vector perpendicular to the original vector.) Using this trick, the magnetic field from the left wire is:

pic

For the right wire your thumb should point into the page, making your right-hand fingers curl clockwise.

pic

pic

Adding these two contributions together yields

pic

### More General Current Distribution

Find the magnetic field at the origin. The wire forms a circle of radius  $R$ .

PIC

Finding the magnetic field from a current-carrying wire involves several distinct steps. Until you become very comfortable setting up and evaluating magnetic field integrals, I would suggest you systematically walk through these steps.

1. Carefully identify and label the location of the differential element on a diagram of the situation.
2. Carefully identify and label the location of the point of interest on a diagram of the situation.
3. Write an expression for ( $pic$ ), the vector differential element.
4. Write an expression for  $r$ , the distance between the differential element and the point of interest.
5. Write an expression for ( $\hat{r}$ ), the unit vector representing the direction from the element to the point of interest.
6. Insert your expressions into the integral for the magnetic field.
7. Carefully choose the limits of integration.
8. Evaluate the integral.

I'll demonstrate each of these steps for the scenario under investigation.

1. Carefully identify and label the location of the differential element on a diagram of the situation.

The differential element is a small (infinitesimal) piece of the current-carrying wire. The location of this differential element must be *arbitrary*, meaning it is not at a "special" location like the top or bottom of the loop. Its location must be represented by a variable, where this variable is the variable of integration and determines the limits of the integral.

For this example, select the differential element to be located at an angle " $q$ " counter-clockwise from the x-axis. (Later, you will select the limits of integration to go from 0 to  $2\pi$  to allow this arbitrary element to "cover" the entire loop.)



pic

2. Carefully identify and label the location of the point of interest on a diagram of the situation.

The point of interest is the center of the loop of wire.

pic

3. Write an expression for  $(pic)$ , the vector differential element.

The differential element in the integral for magnetic field is a vector, meaning it has both a magnitude (its length) and a direction. The direction of the differential element is the direction in which the current is flowing through the element. Thus, we need an expression for the vector illustrated at right (and greatly magnified below).

The length of the differential element is  $Rd\theta$ , since the element forms an arc on a circle of radius  $R$ . Its direction can be determined by noting that it makes an angle of  $\theta$  with respect to the  $y$ -axis, and is directed in the  $-x$ -direction and  $+y$ -direction. This results in:

pic

4. Write an expression for  $r$ , the distance between the differential element and the point of interest.

The distance between the differential element and the point of interest is just the radius of the loop;

pic

5. Write an expression for  $\hat{r}$ , the unit vector representing the direction from the element to the point of interest.

Since this vector points from the element to the point of interest, it is directed in the  $-x$ -direction and  $-y$ -direction. This results in:

pic

pic

6. Insert your expressions into the integral for the magnetic field.

pic

7. Carefully choose the limits of integration.

The limits of integration are determined by the range over which the differential element must be "moved" to cover the entire object. In this case, the element must move all the way around the circular loop:

pic

8. Evaluate the integral.

pic

The magnetic field at the center of any current-carrying loop of radius  $R$  is given by the expression above.

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## 03. Analysis Tools 2

### 1. Ampere's Law

#### Ampere's Law

The long, hollow-core wire below has inner radius  $a$ , outer radius  $b$  and current  $i$  uniformly distributed across its area. The current flows in the  $+z$ -direction. Find the magnetic field at all points in space.

pic

For certain situations, typically ones with cylindrical symmetry, Ampere's Law allows you to calculate the magnetic field relatively easily. Ampere's Law, mathematically, states:

pic

Let's describe what this means in English. The left side of the equation involves the vector dot product between the magnetic field and an infinitesimally small length that is a piece of a larger closed path (termed the *amperian loop*). This dot product determines the amount of magnetic field that is parallel to this very small piece of a larger closed path. The integral simply tells us to add up all of these contributions around the entire closed path.

The gist of Ampere's Law is that this integral is exactly equal to the total amount of electric current flowing within the amperian loop, multiplied by the constant  $\mu_0$ .

Somewhat counter intuitively, the key to applying Ampere's Law is to choose an amperian loop such that you never really have to do the integral on the left side of the equation! (Does this sound familiar?) To try to help you understand what I'm talking about, let's walk through the solution of the above problem. The following sequence of steps will help you understand the process of applying Ampere's law:

1. Choose the appropriate amperian loop.
2. Carefully draw the hypothetical amperian loop at the location of interest.
3. Carefully draw the magnetic field at all points on the amperian loop.
4. Write an expression for the path length parallel to the magnetic field.
5. Write an expression for  $i_{\text{enclosed}}$ , the current inside the amperian loop.
6. Apply Ampere's Law and determine the magnetic field at all points on this hypothetical loop.

There are three distinct regions we will investigate:

- the region inside the "hole" in the wire ( $r < a$ ),
- the region within the actual material of the wire ( $a < r < b$ ),
- and the region outside of the wire ( $r > b$ ).

Let's start outside of the wire.

#### Outside of the wire: $r > b$

1. Choose the appropriate amperian loop.

The key to using Ampere's Law is to try to exploit the symmetry of the current-carrying wire. Since wires typically have circular cross-section, circular amperian loops are the norm.

2. Carefully draw the hypothetical amperian loop at the location of interest.

Since we are trying to determine the magnetic field for all points outside of the wire, draw a circular amperian surface with radius  $r$ . The value of  $r$  is *variable*, and can take on any value greater than  $b$ , the radius of the *real* wire. Remember, the amperian loop is *hypothetical*; it's a mathematical "object" that only exists to help you solve the problem. Try not to confuse it with the real wire of radius  $b$ .

pic

3. Carefully draw the magnetic field at all points on the amperian loop.

Although I have no idea what the *magnitude* of the magnetic field is at any point on my amperian loop, the symmetry of the situation tells me that the *direction* of the magnetic field must be either clockwise or counterclockwise. If the current is flowing out

of the page, the field will be directed counterclockwise around the amperian loop.

pic

Moreover, even though I don't know the magnitude of the field, I *do* know that the magnitude is the same at every point on my loop.

4. *Write an expression for the path length parallel to the magnetic field.*

The left side of Ampere's law requires us to evaluate an integral around our amperian loop. The integral requires us to find the dot product between the magnetic field and the differential line element, and integrate this dot product around the entire loop. I mentioned earlier that you should never have to actually do this integral (assuming you chose the "correct" amperian loop). So why don't we have to do this integral?

The vector dot product can be re-written as:

pic

where  $\theta$  is the angle between the magnetic field and the differential line element of the amperian loop. Traversing our amperian loop counterclockwise leads to the diagram at right

pic

Notice that the magnetic field vector and the vector representing the differential element are parallel at every point on the amperian loop. This simplifies the left-hand side of Ampere's law:

pic

Now note that the magnitude of the magnetic field is the same at every point on the amperian loop since every point is equal distance from the current distribution. Thus, magnetic field is constant and can be brought outside of the integral, leaving a pretty easy integral to evaluate.

pic

Notice that the entire left-hand side of Ampere's Law reduces to the product of the magnetic field magnitude and the length of the amperian loop, assuming the loop is parallel to this field. Thus, because of our wise choice of amperian loop, all we really need to calculate is the length of the loop parallel to the magnetic field. For a circular loop:

pic

5. *Apply Ampere's Law and determine the magnetic field at all points on this hypothetical loop.*

Since the amperian loop is outside of the real wire, *all* of the current flowing through the wire is enclosed by the amperian loop. Thus,

pic

6. *Apply Ampere's Law and determine the magnetic field at all points on this hypothetical loop.*

pic

Thus, the magnetic field outside of the hollow-core wire looks exactly the same as the magnetic field outside of a "normal" current-carrying wire.

Now we have to repeat this analysis for the other two regions.

**Within the wire:  $a < r < b$**

1. *Choose the appropriate amperian loop.*

Again choose a circular amperian loop.

2. *Carefully draw the hypothetical amperian loop at the location of interest.*

Since we are trying to determine the magnetic field within the actual wire, the radius of our amperian loop is greater than  $a$  but less than  $b$ .

pic

3. *Carefully draw the magnetic field at all points on the amperian loop.*

As before, the magnitude of the magnetic field must be constant at all points on the amperian loop and directed counterclockwise.

pic

4. Write an expression for the path length parallel to the magnetic field.

The length parallel to the magnetic field is again the length of the circular amperian loop:

pic

5. Write an expression for  $i_{\text{enclosed}}$ , the current inside the amperian loop.

Since the amperian loop is within the wire, not all of the current in the wire is enclosed by the loop. The amount enclosed can be expressed as the product of a *current density* ( $J$ ) and the cross-sectional area enclosed. First, since the current is uniformly distributed throughout the wire I can define the current density as:

pic

Then, the current enclosed by the amperian loop is the product of the current density and the area enclosed by the loop:

pic

6. Apply Ampere's Law and determine the magnetic field at all points on this hypothetical loop.

pic

Thus, the magnetic field inside the hollow-core wire actually increases with increasing  $r$ , since as  $r$  increases, more and more current is available to produce the magnetic field.

**Inside the "hole":  $r < a$**

Since we must choose our amperian loop to have a radius less than  $a$ , it is located inside the hollow center of the wire. Since there is *no* current enclosed by this loop, the magnetic field in this region must be zero.

pic

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## 04. Activities

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## CHAPTER OVERVIEW

### Model 08 - Magnetic Forces

#### Topic hierarchy

- 01. Concepts and Principles
- 02. Analysis Tools
- 03. Analysis Tools 2
- 04. Activities

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## 01. Concepts and Principles

1. From Magnetic Field to Magnetic Force
2. Magnetic Force on a Current-Carrying Wire

### From Magnetic Field to Magnetic Force

As mentioned previously, only *moving* charged particles can interact with a magnetic field. Stationary electric charges are completely oblivious to the presence of magnetic fields. The magnetic force (also known as **Lorentz force**) on a moving electric charge is given by the relation,

$$\vec{F} = q (\vec{v} \times \vec{B}) \quad (1)$$

where

- $q$  is the charge on the particle of interest,
- $\vec{v}$  the velocity of the particle of interest,
- and  $\vec{B}$  is the net magnetic field at the location of the particle of interest (created by all of the *other* moving charge particles in the universe).

Thus, the direction of the magnetic force on a moving charge is more complicated than in the analogous case of the electric force. In addition to determining the magnetic field at the location of the particle, you must know its velocity and perform a vector cross-product.

### Magnetic Force on a Current-Carrying Wire

In many cases, instead of considering moving electric charges individually we will focus our attention on the collection of moving charges that make up an electric current. In an electric current, consider a small amount of charge,  $dq$ , moving with velocity  $\vec{v}$ . This velocity can be represented as:

$$\vec{v} = \frac{d\vec{l}}{dt} \quad (2)$$

where  $\vec{l}$  is the instantaneous displacement of this small collection of charge. Based on this observation, we can calculate the force acting on this small amount of charge by:

$$d\vec{F} = (dq) \vec{v} \times \vec{B} \quad (3)$$

$$d\vec{F} = (dq) \frac{d\vec{l}}{dt} \times \vec{B} \quad (4)$$

$$d\vec{F} = \left( \frac{dq}{dt} \right) d\vec{l} \times \vec{B} \quad (5)$$

Since the current in a wire is the amount of charge passing through any cross-section of the wire per second,

$$d\vec{F} = i (d\vec{l}) \times \vec{B} \quad (6)$$

Therefore the magnetic force on the entire current-carrying wire is given by the relation,

$$\vec{F} = \int i (d\vec{l}) \times \vec{B} \quad (7)$$

where

- $i$  is the current in the wire of interest,
- $d\vec{l}$  is an infinitesimal length of the wire of interest,
- $d\vec{B}$  is the net magnetic field at the location of  $d\vec{l}$  (created by all of the *other* moving charged particles and currents in the universe),
- and the integral is over the entire length of the wire.

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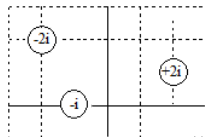


## 02. Analysis Tools

1. Long, Parallel Wires
2. Curved Wires

### Long, Parallel Wires

Three long, parallel wires are located as shown. Each grid square has a width  $a$ . Find the net magnetic force per unit length on the rightmost wire.



The magnetic force on a current-carrying wire is given by the relation,

$$\vec{F} = \int i (d\vec{l}) \times \vec{B} \quad (1)$$

However, for this scenario the relationship simplifies. First, the magnetic field from the two source wires is constant along the entire length of the wire of interest. Second, since the wire of interest is straight,  $d\vec{l}$  is also constant. Therefore, instead of an integral you have:

$$\vec{F} = i\vec{L} \times \vec{B}$$

where  $L$  is the total length of the wire of interest.

To solve this problem, first find the net magnetic field at the location of the wire of interest, and then perform the cross-product. The magnetic field will be the vector sum of the magnetic field for the leftmost wire ( $\vec{B}_L$ ) and the magnetic field from the central wire ( $\vec{B}_C$ ).

$$\begin{aligned} \vec{B}_L &= \frac{\mu_0 i}{2\pi r} \hat{t} & \vec{B}_C &= \frac{\mu_0 i}{2\pi r} \hat{t} \\ \vec{B}_L &= \frac{\mu_0 (2i)}{2\pi \sqrt{(4a)^2 + (a)^2}} \left( \frac{-a\hat{i} - 4a\hat{j}}{\sqrt{(4a)^2 + (a)^2}} \right) & \vec{B}_C &= \frac{\mu_0 (i)}{2\pi \sqrt{(3a)^2 + (a)^2}} \left( \frac{a\hat{i} - 3a\hat{j}}{\sqrt{(3a)^2 + (a)^2}} \right) \\ \vec{B}_L &= \frac{2\mu_0 i}{2\pi (17a^2)} (-a\hat{i} - 4a\hat{j}) & \vec{B}_C &= \frac{\mu_0 i}{2\pi (10a^2)} (a\hat{i} - 3a\hat{j}) \\ \vec{B}_L &= \frac{\mu_0 i}{17\pi a} (-\hat{i} - 4\hat{j}) & \vec{B}_C &= \frac{\mu_0 i}{20\pi a} (\hat{i} - 3\hat{j}) \\ \vec{B}_L &= (-0.0588\hat{i} - 0.235\hat{j}) \frac{\mu_0 i}{\pi a} & \vec{B}_C &= (0.05\hat{i} - 0.15\hat{j}) \frac{\mu_0 i}{\pi a} \end{aligned}$$

Adding these two contributions together yields

$$\vec{B} = (-0.0088\hat{i} - 0.385\hat{j}) \frac{\mu_0 i}{\pi a}$$

So the magnetic force on the rightmost wire is:

$$\begin{aligned} \vec{F} &= i\vec{L} \times \vec{B} \\ \vec{F} &= (2i)(L\hat{k}) \times (-0.0088\hat{i} - 0.385\hat{j}) \frac{\mu_0 i}{\pi a} \\ \vec{F} &= (0.385\hat{i} - 0.0088\hat{j}) \frac{2\mu_0 i^2 L}{\pi a} \\ \vec{F} &= (0.245\hat{i} - 0.0056\hat{j}) \frac{\mu_0 i^2 L}{a} \end{aligned}$$

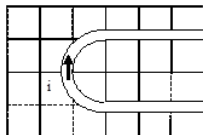
and the force per unit length is

$$\frac{\vec{F}}{L} = (0.245\hat{i} - 0.0056\hat{j}) \frac{\mu_0 i^2}{a}$$

The force is to the right, and slightly downward. Notice that wires carrying currents in opposite directions tend to *repel* each other, while wires carrying currents in the same direction tend to *attract* each other.

## Curved Wires

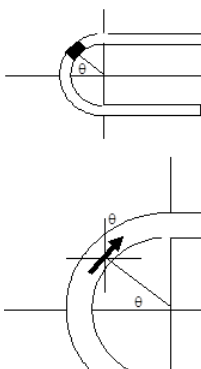
The bent wire below carries current  $i$ , consists of two straight segments of length  $3l$  and a half-circle of radius  $L$ , and lies in a region of uniform magnetic field,  $B$ , in the  $+x$ -direction. Find the total force acting on the wire.



To find the total magnetic force on this bent wire, treat the wire as three separate wires. The force on the two straight sections is zero, because the current in these sections is parallel to the field, resulting in no magnetic force. Therefore, the force on the entire wire is equal to the force on the curved section, which will require setting up and evaluating an integral.

The differential element is located at an angle  $\theta$ , measured clockwise from the  $-y$ -axis. At this location,  $d\vec{l}$  is directed in the  $+x$ -direction and  $+y$ -direction. This results in:

$$d\vec{l} = Ld\theta(\sin \hat{\theta} + \cos \hat{\theta})$$



Using this differential element, the force on the curved section of wire is given by:

$$\begin{aligned}\vec{F} &= \int id\vec{l} \times \vec{B} \\ \vec{F} &= \int_{-\pi/2}^{\pi/2} i(Ld\theta(\sin \hat{\theta} + \cos \hat{\theta})) \times (B\hat{i}) \\ \vec{F} &= -iLB\hat{k} \int_{-\pi/2}^{\pi/2} d\theta \cos \theta \\ \vec{F} &= -iLB\hat{k}(\sin(\frac{\pi}{2}) - \sin(-\frac{\pi}{2})) \\ \vec{F} &= -2iLB\hat{k}\end{aligned}$$

The magnetic force will cause the curved end of the wire to sink into the plane of the page.

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## 03. Analysis Tools 2

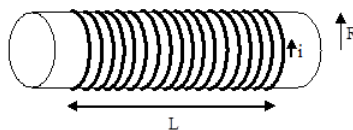
### 1. Force and Motion

### Force and Motion

Protons are injected at  $2.0 \times 10^5$  m/s into a 2000 turn-per-meter solenoid carrying 3.0 A clockwise in the diagram at right. Determine the protons' orbit radius.



A solenoid is a very useful device for generating a uniform magnetic field. A solenoid consists of a wire carrying current  $i$  wrapped  $N$  times around a hollow core of radius  $R$  and length  $L$ . Typically,  $L$  is substantially larger than  $R$  (much larger than illustrated below). Contrast this with a *coil* of wire, in which  $R$  is typically larger than  $L$ .



In a solenoid, the magnetic field inside the core is extremely uniform. In a sense, the small radius and long length "concentrate" the magnetic field within the core leading to an approximately constant value. Again, this contrasts with a coil of wire, in which the field varies at different locations inside and outside the coil.

By using **Ampere's Law**, and a few simplifying approximations, it can be shown that the field inside the solenoid is given by:

$$B = \mu_0 \left( \frac{N}{L} \right) i$$

$$B = \mu_0 n i$$

where  $n$  is the *turn density*, the number of loops of wire, or turns, per meter. Therefore, the solenoid described in the problem creates a magnetic field

$$B = \mu_0 n i$$

$$B = (1.26 \times 10^{-6})(2000)(3)$$

$$B = 7.56 \times 10^{-3} \text{ T}$$

This field is directed into the page since the current flows clockwise.

This magnetic field will create a magnetic force on the proton. When the proton first enters the device the magnetic force will be directed upward, causing the path of the proton to bend upward. As the proton begins to move upward the direction of the magnetic force changes, and when the proton is moving directly upward the force will be to the left. This will cause the proton to bend toward the left. When the proton is moving directly leftward the magnetic force will be directed downward, causing the proton to begin to bend downward. And so on...

Since the magnetic force is always perpendicular to the direction of travel of the proton, the magnetic force causes the proton to make a non-stop left-hand turn! The proton will begin moving in circles due to this force, and since the force has no component along the direction of travel of the proton it does no work on the proton and the proton moves at constant speed. Basically, magnetic fields "steer" charged particles but don't make them speed up or slow down.

With this in mind, let's apply Newton's Second Law to the proton the instant it enters the solenoid:

$$\vec{F} = m\vec{a}$$

$$q\vec{v} \times \vec{B} = m\vec{a}$$

It enters the solenoid traveling in the +x-direction, and it will accelerate toward the center of its circular path, the +y-direction. (Remember that this radial acceleration can be expressed as  $v^2/r$ .)

$$\begin{aligned}
 q(v\hat{i}) \times (-\mu_0 ni\hat{k}) &= m\left(\frac{v^2}{r}\hat{j}\right) \\
 qv\mu_0 ni\hat{j} &= m\frac{v^2}{r}\hat{j} \\
 q\mu_0 ni &= m\frac{v}{r} \\
 r &= \frac{mv}{q\mu_0 ni} \\
 r &= \frac{(1.67 \times 10^{-27})(2 \times 10^5)}{(1.6 \times 10^{-19})(7.56 \times 10^{-3})} \\
 r &= 0.28m
 \end{aligned}$$

The proton will circle at constant speed at this radius.

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## 04. Activities

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## CHAPTER OVERVIEW

### Model 09 - Electromagnetic Induction

#### Topic hierarchy

- 01. Concepts and Principles
- 02. Analysis Tools
- 03. Activities

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## 01. Concepts and Principles

### 1. Creating Electrical Energy

#### Creating Electrical Energy

When electric charges move, their electric fields vary. In the previous two chapters we considered moving electric charges as the source of magnetic fields, but we could just as easily have considered the variation in the electric field as the source of the magnetic field. This leads to an interesting question: If varying electric fields can create magnetic fields, can varying magnetic fields create electric fields? The answer is yes, and this process, termed *electromagnetic induction*, is at the heart of almost all electrical power generation worldwide. In addition to incredible technological importance, electromagnetic induction hints at a deep inter-relationship and symmetry between electric and magnetic fields that will be explored more fully in later chapters.

Imagine a region of space with a magnetic field. Surrounding a portion of this region is a hypothetical closed path. (Often, a real loop of wire will be the closed path of interest, but induction occurs whether or not a real wire loop is present.)

pic 1

First, let me define *magnetic flux*. In analogy to electric flux, introduced in conjunction with Gauss' Law, magnetic flux is defined to be:

pic 2

This equation involves the vector dot product between the magnetic field and an infinitesimally small area within the area bounded by the closed path. This dot product between magnetic field and area is often visualized as the amount of field that "passes through" the little piece of area. The integral simply tells us to add up all of these infinitesimal magnetic fluxes to get the total flux through the area enclosed by the path.

Second, let me define *emf*. (Actually, let me apologize. *Emf* used to stand for electromotive force, even though it is not a force. In light of this misleading name, *emf* now, officially, stands for *emf*. It's not short for anything. I'm not making this up.) *Emf* is the name for the electrical energy per unit charge created by changing magnetic flux. In general, any process that generates electrical energy "creates" *emf*. In addition to changing magnetic flux, chemical batteries and some solar cells create *emf*. Mathematically, *emf* is defined by integrating the electric field around the closed path described above:

pic 3

(Don't worry, you will never actually calculate this integral.) All you need to know is that *emf* is the energy created per unit charge. The units of *emf*, joules per coulomb, are given the name *volts* (V). In crude language, an *emf* is a "voltage".

Let's put this all together. The central relationship describing electromagnetic induction, termed *Faraday's Law*, claims that:

pic 4

where

- $\mathcal{E}$  is the *emf* induced (the voltage created) in the closed path,
- $\Phi_B$  is the magnetic flux that passes through the closed path,
- and the negative sign indicates that the *emf*'s direction in the closed path is to *oppose the change in magnetic flux*. (If the closed path is a real loop of wire, the *emf* will drive an induced current whose direction is such that the magnetic field produced by this induced current is opposite to the change in magnetic field that produces the induced current. Crystal clear?)

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## 02. Analysis Tools

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1. [Abrupt Change in Flux](#)
2. [Continuous Change in Flux](#)

### Abrupt Change in Flux

*An  $N$ -turn loop of radius  $r$  is  $R$  from a very long straight wire carrying current  $i$ . If the current is reduced to zero in a time  $T$ , what is the average induced emf in the loop? What is the direction of the induced current? Assume  $r \gg R$ .*

pic 5

When the current is reduced, the magnetic flux through the small loop will change. By Faraday's Law, this change in flux will create an emf in the loop. The first step toward finding the emf involves finding the magnetic flux through the loop of interest. For this example, the direction of the magnetic field and the direction of the loop's area are parallel, so the dot product reduces to:

pic 6

Moreover, since the loop is far from the wire, and the loop's diameter is small, the magnetic field from the wire is approximately constant over the area of the loop, and the flux is given by:

pic 7

The  $N$  indicates that there are  $N$  loops of wire through which the flux passes.

This flux is reduced to zero over some finite time interval. Interpreting Faraday's Law over a finite time interval results in:

pic 8

This is the average emf induced during the time period under investigation.

Since initially the magnetic field through the loop was in the  $+z$ -direction, and was then removed, the induced emf will drive an induced current counterclockwise around the loop in an attempt to counter the reduction in flux. In other words, the induced current will try to maintain a constant magnetic flux through the loop.

### Continuous Change in Flux

*A 1000-turn secondary coil of radius 2.0 cm is concentric with a 400-turn primary coil of radius 20 cm carrying AC current at 60 Hz with peak current 2.0 A. (AC current can be modeled as  $i(t) = i_0 \cos(2\pi ft)$ .) What is the induced emf in the secondary coil as a function of time?*

pic 9

Again, the first step toward finding the emf involves finding the magnetic flux through the secondary loop. For this example, the direction of the magnetic field and the direction of the loop's area are parallel, flux reduces to:

pic 10

Moreover, since the magnetic field is approximately constant over the relatively small area of the secondary loop, the flux is given by:

pic 11

Plugging in values yields:

pic 12

Applying Faraday's Law:

pic 13

Thus, the secondary coil of wire has 60 Hz AC voltage  $90^\circ$  out of phase with the voltage in the primary coil. (Sine and cosine functions are  $90^\circ$  out of phase with each other.) The maximum emf in the secondary coil is 1.19 V.

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## 03. Activities

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## CHAPTER OVERVIEW

### Model 10 - Electric Potential

#### Topic hierarchy

- 01. Concepts and Principles
- 02. Analysis Tools
- 03. Analysis Tools 2
- 04. Activities

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## 01. Concepts and Principles

1. An Alternative Approach
2. A Gravitational Analogy
3. Relating the Electric Field and the Electric Potential

### An Alternative Approach

The electric field surrounding electric charges and the magnetic field surrounding moving electric charges can both be conceptualized as information embedded in space. In both cases, the information is embedded as *vectors*, detailing both the magnitude and direction of each field. Moreover, when this information is "read" by other moving electric charges, the result is a *force* acting on the charge. These forces can be calculated to allow us to determine the subsequent motion of the charge.

Just as in mechanics, there is an alternative to this force-based approach to analyzing the behavior of electric charges. In this chapter I will define a new field, the *electric potential*, which surrounds every electric charge. This field differs from the electric field because it is a *scalar* field, meaning that this field has only a magnitude at every point in space and no associated direction. Moreover, the information in this field, when read by other electric charges, does not result in a force on the charge but rather determines the *electric potential energy* the charge possesses at that point in space.

Understanding how to calculate this new field, how this field relates to energy, and how this field is related to the electric field will be the focus of this chapter.

### A Gravitational Analogy

Rather than thinking in terms of the gravitational force and Newton's Second Law, an alternative way to examine mechanics scenarios is by using the concept of gravitational potential energy and the conservation of energy.

In the force approach, we envision a vector field surrounding the earth, regardless of whether a second mass is nearby to interact with this field. If a mass is present, the mass interacts with this field and feels a gravitational force. This idea is captured in the equation:

pic 1

In the energy approach, we can envision a scalar field, the *gravitational potential*, which is present regardless of whether a second mass is nearby to interact with this field. If a mass is present, the mass interacts with this field and has gravitational potential energy. Near the surface of the earth, the familiar expression for gravitational potential energy:

pic 2

can be thought of as the product of the mass of the object and this pre-existing gravitational potential,  $V_G$ :

pic 3

if we define the gravitational potential by:

pic 4

Although we didn't use the concept of gravitational potential while studying mechanics, it will prove to be a very useful concept in our study of electrical phenomenon. The general expression for gravitational potential, valid regardless of distance from a massive object, is:

pic 5

In summary, just as a mass will interact with the vector gravitational field as a force, a mass will interact with the scalar gravitational potential field as potential energy.

The situation is very similar for electrical phenomenon. We can envision a scalar field, the *electric potential*, which is present regardless of whether a second charge is nearby to interact with this field. If a charge is present, the charge interacts with this field and has *electric potential energy*. The electric potential,  $V_E$ , is defined by the relationship:

pic 6

where

- $q$  is the source charge, the electric charge that creates the field,
- and  $r$  is the distance between the source charge and the point of interest.

This leads to an expression for the electric potential energy of:

pic 7

where

- $q$  is the charge on the particle of interest, the charge that is interacting with the field,
- and  $V_E$  is the net electric potential at the location of the particle of interest (created by all of the *other* charged particles in the universe).

We will typically leave the subscript off the electric potential and electric potential energy unless the possibility of confusion with the gravitational potential and potential energy are present.

## Relating the Electric Field and the Electric Potential

The electric field and the electric potential are not two, independent fields. They are two independent ways of conceptualizing the effect that an electric charge has on the space surrounding it. Just as problems in mechanics can be analyzed using a force-approach or an energy-approach, problems dealing with electrical phenomenon can be analyzed by focusing on the electric field or on the electric potential.

Additionally, just as it is sometimes necessary in mechanics to transfer between force and energy representations, it is sometimes necessary to transfer between the electric field and electric potential representations. The relationship between two fields can be understood by examining the expression for work, which relates force to transfer of energy.

Utilizing the dot product, the work done in moving a particle from an initial point,  $i$ , to a final point,  $f$ , can be written as:

pic 8

where (pic 9) is an infinitesimal portion of the path along which the particle moves. You may also recall that the difference in potential energy between initial and final locations is defined as the opposite of the work needed to move the particles between the two points:

pic 10

Putting these two ideas together yields:

pic 11

We can now use this result to relate the electric potential and the electric field. Substituting in expressions for potential energy and force in terms of the fields that convey them leads to:

pic 12

In English, this final result states that the electric potential difference between any two points is defined as the negative of the integral of the electric field along a path connecting the two points. (I'm sure that doesn't seem like particularly *clear* English, but this idea will become more tangible once you get to work on some of the fun activities in this chapter.) The bottom line is that the electric potential can be determined by integrating the electric field, and, conversely, the electric field can be determined by differentiating the electric potential.

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## 02. Analysis Tools

### Point Charges

*Find the electric potential at the indicated point. The charges are separated by a distance  $4a$ .*

Pic 1

The electric potential at the point specified will be the sum of the electric potential from the left charge ( $V_L$ ) and the electric potential from the right charge ( $V_R$ ).

Pic 2

Notice that since the electric potential is a scalar, calculating the electric potential is often much easier than calculating the electric field.

### Continuous Charge Distribution

*The plastic rod of length  $L$  at the bottom has uniform charge density  $\lambda$ . Find the electric potential at all points to the right of the rod on the  $x$ -axis.*

Pic 3

Since electric charge is discrete, the electric potential can always be calculated by summing the electric potential from each of the electrons and protons that make up an object. However, macroscopic objects contain a *lot* of electrons and protons, so this summation has many, many terms:

pic 4

As described earlier, we will replace this summation over a very large number of discrete charges with an integral over a hypothetically continuous distribution of charge. This leads to a relationship for the electric potential at a particular point in space, from a continuous distribution of charge, of:

pic 5

where  $dq$  is the charge on an infinitesimally small portion of the object, and the integral is over the entire physical object.

The steps for finding the electric potential from a continuous distribution of charge are:

1. Carefully identify and label the location of the differential element on a diagram of the situation.
2. Carefully identify and label the location of the point of interest on a diagram of the situation.
3. Write an expression for  $dq$ , the charge on the differential element.
4. Write an expression for  $r$ , the distance between the differential element and the point of interest.
5. Insert your expressions into the integral for the electric potential.
6. Carefully choose the limits of integration.
7. Evaluate the integral.

I'll demonstrate below each of these steps for the scenario under investigation.

1. Carefully identify and label the location of the differential element on a diagram of the situation.

The differential element is a small (infinitesimal) piece of the object that we will treat like a point charge. The location of this differential element must be *arbitrary*, meaning it is not at a "special" location like the top, middle, or bottom of the rod. Its location must be represented by a variable, where this variable is the variable of integration and determines the limits of the integral.

For this example, select the differential element to be located a distance " $y$ " above the center of the rod. The length of this element is " $dy$ ". (Later, you will select the limits of integration to go from  $-L/2$  to  $+L/2$  to allow this arbitrary element to "cover" the entire rod.)

Pic 6

2. Carefully identify and label the location of the point of interest on a diagram of the situation.

You are interested in the electric potential at all points along the x-axis to the right of the rod. Therefore, select an *arbitrary* location along the x-axis and label it with its location.

pic 7

3. Write an expression for  $dq$ , the charge on the differential element.

The rod has a uniform charge density  $\lambda$ , meaning the amount of charge per unit length along the rod is constant. Since the differential element has a length  $dy$ , the total charge on this element ( $dq$ ) is the product of the density and the length:

pic 8

4. Write an expression for  $r$ , the distance between the differential element and the point of interest.

By Pythagoras' theorem, the distance between the differential element and the point of interest is:

pic 9

6. Insert your expressions into the integral for the electric potential.

pic 10

7. Carefully choose the limits of integration.

The limits of integration are determined by the range over which the differential element must be "moved" to cover the entire object. The location of the element must vary between the bottom of the rod ( $-L/2$ ) and the top of the rod ( $+L/2$ ) in order to include every part of the rod. The two ends of the rod form the two limits of integration.

pic 11

8. Evaluate the integral.

pic 12

This expression looks daunting, but its limiting behaviors (as  $x$  approaches zero and  $x$  approaches infinity) are correct. Consider the case where  $x$  gets smaller and smaller, approaching zero:

pic 13

This expression diverges (goes to infinity) as  $x$  becomes close to zero, as you should expect since the electric potential directly "on" an electric charge is infinite.

As  $x$  gets larger and larger, the  $L/2$  terms become insignificant:

pic 14

This expression becomes zero as  $x$  becomes infinite, as you should expect since the electric potential extremely far from an electric charge is zero.

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## 03. Analysis Tools 2

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1. [Potential Difference](#)
2. [Electric Potential Energy](#)

### Potential Difference

*The long, hollow plastic cylinder at the bottom has inner radius  $a$ , outer radius  $2a$ , and uniform charge density  $r$ . Find the electric potential difference between the inner and outer radius.*

pic 1

An alternative method for calculating the electric potential at a point, or the electric potential difference between two points, is by using knowledge of the electric field. The following relation,

pic 2

states that the potential difference between two points can be determined by integrating the component of the electric field that lies along the path connecting the two points. This means that if you know the electric field in a region of space, you can easily (more or less) find the potential difference between any two points that lie in that region. This relation is particularly useful in conjunction with Gauss' Law for situations with cylindrical or spherical symmetry.

To find the potential difference between points  $a$  and  $b$  (i.e., what a voltmeter would read if connected across points  $a$  and  $b$ ), we need the electric field in this region. Gauss' Law can be used to find that (review Gauss' Law if this step is a little fuzzy):

pic 3

This electric field points radially away from the center of the cylinder.

For simplicity, choose a path that directly connects  $a$  to  $2a$ , i.e., a radial path.

pic 4

pic 5

Notice that this method directly calculates the *difference* in electric potential between two points, without actually determining the *value* of the electric potential at either point. Since electric potential is related to electric potential energy, this method allows you to find the difference in energy between two points but not the actual value of the energy of an electric charge.

This should strike you as quite similar to the gravitational case. For gravitational potential energy, the choice of the zero-point is arbitrary and only energy differences lead to changes in kinetic energy. For electrical potential energy, the situation is identical. The zero-point of electric potential energy (and electric potential) is typically taken at infinity, although you can "zero" the potential at a more convenient point by "grounding" the system at that point. The physical act of grounding a point on an electrical device is mathematically equivalent to setting the potential equal to zero at that point.

### Electric Potential Energy

*In many applications, oppositely charged parallel plates (with small holes cut for the beam to pass through) are used to accelerate beams of charged particles. In this example, a proton is injected at  $v_1$  into the space between the plates. The potential difference between the plates is  $DV$  (the left plate is at higher potential). What is the velocity of the proton as it exits the device?*

pic 6

Since the electric potential changes as the proton moves from the left plate to the right plate, its potential energy changes. This change in potential energy results in a change in kinetic energy of the proton by energy conservation.

pic 7

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## 04. Activities

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## SECTION OVERVIEW

### Model 11 - Electric Circuits

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#### 02. Analysis Tools

#### 03. Analysis Tools 2

#### 04. Activities

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## 01. Concepts and Principles

1. [Electric Circuits as Applied Physics](#)
2. [The Resistor](#)
3. [The Capacitor](#)
4. [The Inductor](#)

### Electric Circuits as Applied Physics

Electric circuits are one of the most practical applications of our understanding of electric and magnetic fields. In general, an electric circuit is any device that consists of a closed path for charges to move (a current), a source of energy to "drive" the motion of the charges (a potential difference, or voltage, often in the form of a battery), and various circuit elements that can either convert (resistors) or store (capacitors and inductors) the energy supplied by the energy source.

The study of circuits is incredibly broad, since there are limitless ways to combine these elements into an electric circuit. We will restrict ourselves to studying circuits with only a limited number of elements, and with a source that supplies a constant voltage<sup>[1]</sup>.

### The Resistor

In general, a *resistor* is any device that converts electrical energy into another form of energy, often heat. For example, a fluorescent light bulb converts electrical energy into light (with about a 20% efficiency, the remaining energy is converted into heat) and an incandescent light bulb converts electrical energy very efficiently into heat (with only about 5% of the incident energy converted to light). Since a conversion of electrical energy takes place in these devices, they are resistors.

In all resistors, the electric potential energy of the charges entering the device is larger than the electric potential energy of the charges exiting the device, because some of the potential energy has been converted to other forms. This decrease in potential energy is due to a decrease in electric potential between the two ends of the device and is directly proportional to the *resistance* of the device.

The definition of resistance for a device is:

pic 1

where

- $\Delta V$  is the potential difference between the two ends of the device, often termed the *voltage drop* across the device,
- and  $i$  is the current that flows through the device.

The unit of resistance, pic 2, is defined as the ohm ( $\Omega$ ).

The previous expression relates the resistance of a resistor to properties of the circuit it is part of. However, it is also sometimes useful to directly relate the resistance to the actual physical parameters of the device itself. For simple, passive resistors (basically blocks of material connected to a voltage source), resistance is defined as:

pic 3

where

- $\rho$  is the *resistivity* of the material from which the resistor is constructed,
- $L$  is the length of the resistor in the direction of current flow,
- and  $A$  is the cross-sectional area of the resistor.

Resistivity can range from 0 for a perfect conductor to for a perfect insulator.

One final note on the properties of resistors concerns their rate of energy conversion. Since electric potential is the electric potential per unit of charge, and current is the charge flowing through the device per second, the product of change in electric potential and current is the change in electric potential energy per second. Thus, the rate of energy conversion, or *power*, in a resistor is given by:

pic 4

### The Capacitor

A *capacitor* is a device that stores energy in the electric field between two closely spaced conducting surfaces. When connected to a voltage source, electric charge accumulates on the two surfaces but, since the conducting surfaces are separated by an insulator, the charges cannot travel from one surface to the other. The charges create an electric field in the space between the surfaces, and the two surfaces have a difference in electrical potential.

pic 5

Once "charged", if the capacitor is removed from the original circuit and connected to a second circuit it can act as a voltage source and "drive" its collected charge through the second circuit. When used in this way, the capacitor clearly acts as a temporary storehouse of energy.

To determine the energy stored in a capacitor, we first need to define the *capacitance* of the capacitor. The capacitance of a capacitor is defined as:

pic 6

where

- $Q$  is the magnitude of the electric charge stored on either conducting surface,
- and  $\Delta V$  is the potential difference between the surfaces.

The unit of capacitance, pic 7, is defined as the *farad* (F).

The amount of energy that can be stored on a capacitor is a function of both its capacitance and the potential difference between its surfaces. The relationship between stored energy and these parameters is:

pic 8

## The Inductor

An *inductor* is a device that stores energy in the magnetic field created when current passes through a coil of wire. When connected to a voltage source, current will flow through the inductor, establishing a magnetic field.

pic 9

If the voltage source is suddenly removed, current will continue to flow in the coil because of electromagnetic induction. This induced current will act to replace the disappearing source current. The energy needed to drive this current comes from the energy stored in the magnetic field, so in this case the inductor acts as a temporary storehouse of energy.

To determine the energy stored in an inductor, we first need to define the *inductance* of the inductor. The inductance of the inductor is defined as:

pic 10

where

- $\Phi$  is the magnetic flux within the inductor,
- and  $i$  is the current flowing through the inductor.

The unit of capacitance, pic 11, is defined as the *henry* (H).

The amount of energy that can be stored in an inductor is a function of both its inductance and the current flowing through it. The relationship between stored energy and these parameters is:

pic 12

---

[1] Circuits with constant voltage sources are referred to as DC, or direct current, circuits.

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## 02. Analysis Tools

1. Resistors in Circuits
2. Capacitor Properties
3. Capacitors in Circuits

### Resistors in Circuits

The circuit at right represents a 12 V car battery and two mismatched headlights,  $R_1 = 1.9 \text{ W}$  and  $R_2 = 2.1 \text{ W}$ .

- a. Determine the magnitude of the potential difference across and the current through each circuit component.
- b. If the battery has a total stored energy of 800 W hr, and produce a constant potential difference until discharged, how long will the bulbs stay lit?

pic 1

The potential difference across the car battery is given as 12 V. This means that the electric potential in the wire coming out of the "top" of the battery is 12 V larger than the potential in the wire coming from the "bottom". Since each of the resistors are attached to these same two wires, the top of each resistor is 12 V higher in potential than the bottom. Therefore the potential difference across each resistor is 12 V. When circuit elements are connected such that the elements all have the same potential difference, the elements are said to be in *parallel*.

Since the potential difference across each resistor is known, we can use the definition of resistance to calculate the current through each branch of the circuit. Analyzing branch #1 yields

pic 2

and branch #2

pic 3

The current that flows through  $R_1$  and the current that flows through  $R_2$  must also flow through both the top and bottom wires connected to the battery. To complete the mental image of a closed circuit of current, we will say the current flows "through" the battery as well, although this is not technically true. Therefore, the current that flows through the battery (the total current flowing in the circuit) is:

pic 4

We can summarize this information in a simple table:

	$V_{\text{across}} \text{ (V)}$	$i_{\text{through}} \text{ (A)}$
battery	12	12.0
$R_1$	12	6.32
$R_2$	12	5.71

To determine how long the headlights will stay lit, we must calculate the total power of the circuit (the total amount of electrical energy converted per second). We can do this separately for each headlight and then add the results:

pic 5

and

pic 6

so the total power of the circuit is:

pic 7

Therefore, the battery will last for

pic 8

## Capacitor Properties

Imagine a pair of long, hollow nested cylinders of inner radius  $a$  and outer radius  $b$ . Calculate the capacitance, per meter, for these nested cylinders.

pic 9

Since capacitance is defined by the relation

pic 10

we need to determine the potential difference that would develop between these cylinders if charges  $Q$  (and  $-Q$ ) were placed on the two surfaces. To do this, imagine that a charge  $+Q$  (per meter) was placed on the inner cylinder. Using Gauss' Law, this leads to an electric field between the cylinders of:

pic 11

This field is directed radially away from the central axis of the cylinders.

Once the electric field between the cylinders is known, the magnitude of the potential difference between the cylinders can be calculated by:

pic 12

Substituting this result into definition of capacitance yields:

pic 13

Thus, the capacitance per meter of a set of nested cylinders depends on the natural logarithm of the ratio of the cylinder radii. Notice that if the cylinders are very close together ( $b$  is not much larger than  $a$ ), the capacitance is very large. The capacitance of a capacitor is always enhanced by having the two charged surfaces very close together. However, as the surfaces get closer together, the possibility of electrical breakdown (charges "jumping" across the gap) becomes larger. For this reason, and several others, the space between the surfaces in a capacitor is typically filled with a type of material, called a *dielectric*, which both enhances the capacitance of the capacitor and inhibits electrical breakdown.

## Capacitors in Circuits

The device at right represents a simplified camera flash circuit. With  $V = 3 \text{ V}$  and  $R = 100 \text{ }\Omega$ , find  $C$  such that the flash reaches 80% of its final voltage in 1.0 s.

pic 14

The circuit above, termed an *RC circuit*, can best be analyzed by considering the changes in electric potential experienced by a hypothetical charge "journeying" around the circuit:

- as it "passes through" the battery the potential increases by  $V$ ,
- as it passes through the resistor the potential decreases by

pic 15

- and as it "passes through" the capacitor the potential decreases by

pic 16

Putting these changes in potential together results in:

pic 17

Note that the total change in potential (and potential energy) must be zero since the energy given to the charge by the battery is partially converted by the resistor and partially stored by the capacitor.

If we take a time derivative of the above equation (noting that  $V$ ,  $R$ , and  $C$  are constants, but that  $Q$ , the charge on the capacitor, is changing) we are left with a differential equation for the current in the circuit:

pic 18

This equation says that the time derivative of the current is equal to the product of the current and the numerical factor pic 19. The only mathematical function that has the property that its derivative is proportional to itself is the exponential function. Therefore,

the current must be given by the function:

pic 20

where  $i_0$  is the current at  $t = 0$  s.

If we assume that the capacitor is uncharged when the switch is first closed, then

pic 21

so the final expression for the current in the circuit as a function of time is:

pic 22

Using this expression we can determine the time-dependence of any other circuit parameter.

For example, the question asks about the voltage across the capacitor. Since the voltage across the resistor can be expressed as:

pic 23

the voltage across the capacitor is the amount of the source voltage that "remains":

pic 24

This function shows that after a long time ( $t \rightarrow \infty$ ), the voltage across the capacitor will equal the voltage of the source.

Therefore,

pic 25

Thus, a 6.21 mF capacitor will reach 80% of its final voltage in 1.0 s.

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## 03. Analysis Tools 2

1. Inductor Properties
2. Inductors in Circuits

### Inductor Properties

Imagine a pair of long, hollow nested wires of inner radius  $a$  and outer radius  $b$ , designed to carry current into and out of the page. Calculate the inductance, per meter, for these nested wires.

pic 1

Since inductance is defined by the relation

pic 2

we need to determine the flux that would develop between these wires if current  $i$  (and  $-i$ ) flowed along the two wires:

To do this, imagine that current  $i$  flowed out of the page along the inner wire. Using Ampere's Law, this leads to a magnetic field between the cylinders of:

pic 3

To help calculate the flux between the wires, the diagram at right is a top view of the nested wires. The dashed area is the area over which we will calculate the flux. (The current along the inner wire flows toward the top of the page, resulting in magnetic field pointing directly out of the page in the area of interest.)

The shaded sliver is the differential element, located a distance  $r$  from the center of the wires, with thickness  $dr$  and length  $l$ . The magnetic flux is then:

pic 4

Substituting this result into the definition of inductance yields:

pic 5

The inductance per meter is then:

pic 6

Thus, the inductance per meter of a set of nested wires depends on the natural logarithm of the ratio of the wire radii. Notice that if the wires are very far apart, the inductance is larger. However, as the wires get farther apart, the size of the device gets larger and may become impractical. For this reason, and several others, the space containing the magnetic flux in an inductor is typically filled with a material with a high magnetic permeability, like iron, in order to "concentrate" the magnetic flux into a smaller region of space.

### Inductors in Circuits

The device below represents a simplified electromagnet. With  $V = 100 \text{ V}$  and  $R = 15 \text{ W}$ , find  $L$  such that the current reaches  $5.0 \text{ A}$  in  $0.5 \text{ s}$ .

pic 7

The circuit above, termed an *RL circuit*, can best be analyzed by considering the changes in electric potential experienced by a hypothetical charge "journeying" around the circuit:

- as it "passes through" the battery the potential increases by  $V$ ,
- as it passes through the resistor the potential decreases by

pic 8

- and as it "passes through" the inductor's potential changes by

pic 9

Since by Faraday's Law of Induction,



pic 10

The emf induced by the inductor is the potential drop across it, so

pic 11

Putting these changes in potential together results in:

pic 12

Again, note that the total change in potential (and potential energy) must be zero since the energy given to the charge by the battery is partially converted by the resistor and partially stored by the inductor.

If we take a time derivative of the above equation (noting that  $V$ ,  $R$ , and  $L$  are constant) we are left with a differential equation for the current in the circuit:

pic 13

This equation says that the time derivative of the *derivative* of the current is equal to the product of the *derivative* of the current and a numerical factor. This means that the derivative of the current must be exponential function. Therefore, the derivative of the current must be given by the function:

pic 14

where  $A$  is an arbitrary constant. Integrating this result leads to a current of the form:

pic 15

where  $B$  and  $D$  are arbitrary constants.

To determine these constants, consider the current in the circuit after a very long time ( $t \rightarrow \infty$ ). After this amount of time the circuit will have reached an equilibrium value, so the change in the current will be zero. Thus,

pic 16

Therefore,

pic 17

Now consider the current in the circuit the instant you first close the switch ( $t \rightarrow 0$ ). At this instant, no current can be flowing in the circuit. This is because if there was current flowing instantaneously after the switch was closed, this would be discontinuous change in current and the inductor would create an infinite emf to oppose this "infinite" increase in current. Therefore,

pic 18

Now that we know the values of the two constants, the final expression for the current in the circuit as a function of time is:

pic 19

Using this expression we can determine the time-dependence of any other circuit parameter.

Since the question asks about the current directly,

pic 20

Therefore, if the electromagnet has an inductance of 5.4 H, it will take 0.5 s for the current to rise to 5.0 A.

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## 04. Activities

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## CHAPTER OVERVIEW

### Model 12 - Electromagnetic Waves

#### Topic hierarchy

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## 01. Concepts and Principles

1. Putting it All Together
2. Electromagnetic Waves
3. Energy of Electromagnetic Waves

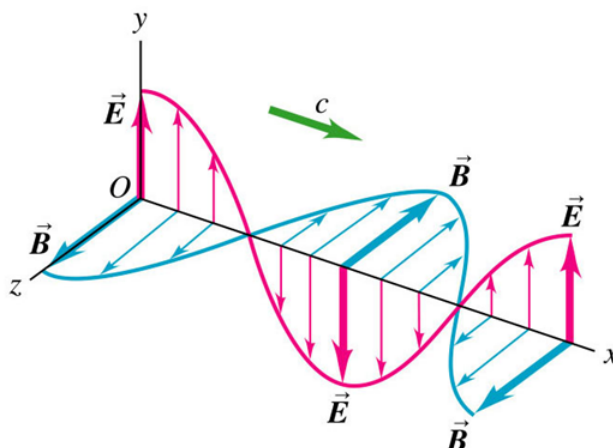
### Putting it All Together

Electric charges are surrounded by electric fields. When these charges move, the electric field changes. Additionally, these moving charges (or, more importantly, these changing electric fields) create magnetic fields. Since before the motion no magnetic field existed, this is a change in the value of the magnetic field (from zero to non-zero). By electromagnetic induction, this change in magnetic field creates emf, which is a change in the value of the electric field in a region of space. But this change in electric field must cause further change in the magnetic field, which must cause further change in the electric field, which must...

Sorting through this interrelationship between electric and magnetic fields, which involves the simultaneous solution of a set of coupled, partial-differential equations, is the grand accomplishment of James Maxwell. This work of Maxwell's, in the 1860s, is generally regarded as on par with the work of Isaac Newton and Albert Einstein. Although we won't actually solve this set of equations, we will study several of the consequences of Maxwell's work.

### Electromagnetic Waves

By far the most important conclusion from Maxwell's work is that the changes in electric and magnetic fields, coupled together as described above, propagate through space as an electromagnetic wave. Near the moving charge that created the wave, the mathematical description of the wave is very complicated, but once the wave has moved a "reasonable" distance from the charge that created it, it can be visualized as below:



and mathematically represented as:

$$\vec{E} = E_{\max} \cos\left(\frac{2\pi x}{\lambda} - 2\pi f t\right) \hat{j}$$

$$\vec{B} = B_{\max} \cos\left(\frac{2\pi x}{\lambda} - 2\pi f t\right) \hat{k}$$

where

- $\lambda$  is the wavelength of the wave, the distance between maximum values of the field vectors,
- $f$  is the frequency of the wave, the number of complete cycles of oscillation that wave passes through per second,
- the electric and magnetic field vectors are perfectly in phase, meaning they pass through their maximum and minimum values at the same points in space and time,
- the field vectors are oriented at  $90^\circ$  from each other,
- and the wave propagates in the direction, given by  $\vec{E} \times \vec{B}$ , which in this case is the +x-direction.

Of course, a similar mathematical description could be given for an electromagnetic wave propagating in any direction.

It's important to realize that the "picture" of the wave above is for a specific instant of time. This is a *traveling* wave, in that it moves through space, in this case in the +x-direction. Since it moves through space, an important characteristic of the wave is its speed. While solving the set of equations, Maxwell found that the speed of electromagnetic waves, denoted  $c$ , is given by the expression:

$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

where  $\epsilon_0$  and  $\mu_0$  are the familiar permittivity and permeability of free space.

What is somewhat odd about this result is that it states that the speed of the wave is a universal constant, since both  $\epsilon_0$  and  $\mu_0$  are constant. This means that regardless of the wavelength of the wave, or its frequency, or whether the charges that created the wave were wiggling in a stationary radio antenna on top of a mountain or in the headlight of a spaceship moving at a billion miles per hour (or if you viewed the wave from a second spaceship moving a billion mph in the other direction!) the wave always moves at exactly the same speed. *The wave speed is completely independent of how it was created.* Once the wave is created, its motion through space is completely determined by how its electric and magnetic fields interact with each other, and its speed has no relationship to the physical charge that created it. (This fact will have some curious repercussions about 40 years after Maxwell's discovery.)

Plugging in the known values for  $\epsilon_0$  and  $\mu_0$  yields:

$$c = 3.0 \times 10^8 \text{ m/s}$$

Maxwell realized that this was the known speed for visible light, and soon came to the conclusion that light was, in fact, an electromagnetic wave. Electromagnetic waves of other frequencies, such as radio waves, microwaves, x rays, etc, all propagate at this same speed and obey the same mathematical framework described above.

One final interesting conclusion from Maxwell's solution was that the ratio of the electric field to the magnetic field at any point in the wave has a constant value, equal to the wave speed:

$$\frac{E}{B} = c$$

## Energy of Electromagnetic Waves

Energy can be stored in electric and magnetic fields, as you learned when you analyzed capacitors and inductors in electric circuits. Since electromagnetic waves involve the propagation of changes in these fields, they should involve the flow of energy in the direction of the wave's motion.

The *intensity*, or energy flow per unit area, of an electromagnetic wave is given by:

$$\vec{S} = \frac{1}{\mu_0} \vec{E} \times \vec{B}$$

where  $\vec{S}$  is termed the *Poynting vector*, and indicates both the direction of the energy flow and its magnitude per unit area. Since the electric and magnetic fields are oriented at  $90^\circ$  from each other, the magnitude of the intensity is given by:

$$S = \frac{EB}{\mu_0}$$

This expression gives you the intensity at any point in the electromagnetic wave at one specific instant of time. However, in many ways this isn't particularly useful. If you are dealing with a visible light wave, for example, the frequency of the wave is a approximately  $10^{14}$  cycles per second. This means that the intensity of the wave cycles through maximum and minimum values every  $10^{-14}$  s!

A much more useful expression would be for the *average* intensity of the wave. To find this expression, rearrange

$$\begin{aligned} \frac{E}{B} &= c \\ B &= \frac{E}{c} \end{aligned}$$

and substitute into the expression for the instantaneous intensity

$$S = \frac{EB}{\mu_0}$$

$$S = \frac{E^2}{c\mu_0}$$

Substituting our description of the electric field-portion of an electromagnetic wave

$$E = E_{\max} \cos\left(\frac{2\pi x}{\lambda} - 2\pi ft\right)$$

into the above expression yields

$$S = \frac{E^2}{c\mu_0}$$

$$S = \frac{(E_{\max} \cos(\frac{2\pi x}{\lambda} - 2\pi ft))^2}{c\mu_0}$$

$$S = \frac{E_{\max}^2 \cos^2(\frac{2\pi x}{\lambda} - 2\pi ft)}{c\mu_0}$$

To find the average value of this function we need only to find the average value of the cosine-squared function, since all of the other terms are constant. It's common knowledge<sup>[1]</sup> that the average value of the cosine-squared function is  $\frac{1}{2}$ . Therefore, the average intensity of an electromagnetic wave is:

$$\bar{S} = \frac{E_{\max}^2}{2c\mu_0}$$

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<sup>[1]</sup>It is now.

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## 02. Concepts and Principles 2

1. Momentum of Electromagnetic Waves
2. Polarization of Electromagnetic Waves
3. Interference of Electromagnetic Waves

### Momentum of Electromagnetic Waves

Just as electromagnetic waves involve the flow of energy, they also carry momentum in the direction of motion. Although it may seem confusing that mass-less, non-material fields, or at least changes in these fields, can have momentum, the momentum carried by electromagnetic waves can be directly measured. The long tail of a comet, pointing directly away from the sun, is a result of the momentum carried by the sun's electromagnetic waves.

The simplest phenomenon directly involving the momentum of electromagnetic waves occurs when the waves are absorbed or scattered from an object. Since the wave has momentum, some of this momentum is transferred to the object the wave "collides" with. This interaction can be directly studied by applying conservation of momentum, but an easier shortcut involves defining the concept of *radiation pressure*. Radiation pressure is the force per unit area that the wave exerts on the object during the collision. Radiation pressure is defined as:

pic 15

where  $a$  is a number between 1 and 2 that depends on the portion of the electromagnetic wave absorbed during the collision. If the wave is fully absorbed,  $a = 1$ . If the wave is perfectly reflected,  $a = 2$ . (We won't analyze situations between these two extremes.)

### Polarization of Electromagnetic Waves

In many cases, the electric field component of an electromagnetic wave oscillates in a well-defined direction. When this is the case, the wave is said to be *polarized*. The polarization direction is the direction in which the electric field oscillates. For example, in the diagram below the wave is polarized in the  $y$ -direction.

pic 16

Certain materials allow only one specific direction of electric field vector to propagate through them. These materials, called *polarizers*, typically absorb (or reflect) all waves with electric field vectors not aligned with their transmission axis. For example, a polarizer with transmission axis along the  $+x$ -direction would completely block the propagation of the wave illustrated above.

If the incident electric field vector is not perfectly aligned with the transmission axis, only the component of the field vector along the transmission axis can pass through the polarizer. Thus the electric field magnitude that passes through the polarizer,  $E$ , is given by:

pic 17

where

- $E_0$  is the electric field magnitude incident on the polarizer,
- and  $\theta$  is the angle between the incident polarization direction and the transmission axis.

After passing through the polarizer, the electromagnetic wave is now polarized along the transmission axis.

Rather than concentrating on the electric field magnitude, in most cases it's more useful to focus on the intensity of the wave. Since the average intensity of the wave is given by

pic 18

the intensity is proportional to the square of the electric field vector. Thus, a more useful version of the polarization equation is:

pic 19

where

- $S_0$  is the intensity incident on the polarizer,
- and  $\theta$  is the angle between the incident polarization direction and the transmission axis.



## Interference of Electromagnetic Waves

Since electromagnetic waves consist of alternating electric and magnetic field vectors, if two or more waves pass through the same point in space (at the same time) their field vectors must add according to the basic principles of vector addition. This combination of waves through the vector addition of their field vectors is called *interference*.

In general, the interference of two electromagnetic waves can be incredibly difficult to analyze. We will restrict ourselves to a specific sub-class of interference phenomenon, in which there are only two sources of waves and the two sources produce waves of exactly the same frequency perfectly in phase. Typically, the easiest way to accomplish this is by having a single source of waves and then somehow dividing the waves and sending each portion of the wave along a separate path to a final common location. (This is a lot simpler than it sounds.)

The interference that results when the two waves are recombined depends only on the *path length difference*,  $\Delta d$ , the waves traveled to reach the recombination location. The path length difference is, as the name implies, simply the difference between the distance traveled by wave #1 and the distance traveled by wave #2 between separation and recombination,

pic 20

If this difference is exactly equal to zero, one wavelength, or any integer number of wavelengths, the waves will be perfectly in phase when they recombine and exhibit *constructive interference*, resulting in a locally maximum value for the electric field vector. Therefore, for constructive interference:

pic 21

where  $m$  is zero or any integer.

If the path length difference is exactly equal to one-half a wavelength, or any half-integer number of wavelengths, the waves will be perfectly out of phase when they recombine and exhibit *destructive interference*, resulting in a locally minimum value for the electric field vector. Therefore, for destructive interference:

pic 22

where  $m$  is zero or any integer.

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## 03. Analysis Tools

1. [Energy, Power and Intensity](#)
2. [Pressure](#)
3. [Polarization](#)
4. [Two Source Interference](#)

### Energy, Power and Intensity

*Future lunar colonists may want to watch their favorite earthly TV shows. Suppose a television station on Earth has a power of 1.0 MW and broadcasts isotropically.*

- a. What is the intensity of this signal on the moon?
- b. What total power would be received by a 20 cm radius "satellite" dish on the moon?
- c. What is the maximum electric field strength of this signal on the moon?

If the signal is broadcast isotropically (the same in all directions), its intensity should decrease as the surface area of the sphere over which the signal spreads increases. Since the distance to the moon is  $3.82 \times 10^8$ ,

pic 1

The power received by the dish would be

pic 2

This corresponds to a maximum electric field of

pic 3

### Pressure

*High-power lasers are used to compress a small hydrogen pellet in an attempt to initiate fusion. A laser generating pulses of radiation of power 1.5 GW is focused isotropically onto the 0.5 mm radius pellet. What is the pressure exerted on the pellet if the pellet absorbs the incident light?*

When an electromagnetic wave is absorbed by a surface, it exerts a pressure on the surface given by

pic 4

The intensity of the light on the surface of the pellet is

pic 5

Thus, the pressure on the pellet is

pic 6

### Polarization

*Initially unpolarized light is sent through three polarizing sheets with transmission axes oriented at  $q_1 = 0^\circ$ ,  $q_2 = 45^\circ$ ,  $q_3 = 90^\circ$  measured counterclockwise from the x-axis. What percentage of the initial intensity is transmitted by the system of the three sheets?*

If polarized light of intensity  $S_0$  is incident on a polarizing sheet, the intensity transmitted by the sheet depends on the angle between the polarization direction of the light and the transmission axis of the sheet,  $q$ , via

pic 7

But what if the incident light is *unpolarized*?

If the incident wave is unpolarized, we can imagine that it is a combination of waves of all possible polarizations. Then, to determine the intensity that passes through the polarizer we can average over all of these different hypothetical polarization directions. Since the intensity that passes through the polarizer depends on the square of the cosine function, this means we must find the average value of the cosine-squared function. Since it's common knowledge<sup>[1]</sup>

that the average value of the cosine-squared functions is  $\frac{1}{2}$ , the intensity after the first sheet is one-half of the initial intensity:

pic 8

and the wave is now polarized along the x-axis.

After the second sheet,

pic 9

and the wave is now polarized at  $45^\circ$  from the x-axis.

After the third sheet,

pic 9

and the wave is now polarized along the y-axis. (Notice that  $\theta$  is **not** the angle of the transmission axis ( $90^\circ$ ), but rather the angle *between* the previous polarization direction and the transmission axis.) Thus, 12.5% of the initial light intensity passes through the three sheets and the resulting light is vertically polarized.

---

[1] It should be by now.

## Two Source Interference

*Two radar sources are separated by  $d = 10$  m. During testing, the two sources broadcast a test frequency perfectly in phase. If destructive interference occurs at  $x = 40$  m, what are the possible values for the source wavelength?*

pic 10

Destructive interference occurs when the two waves have a relative phase difference of  $\pi$  (one wave is "flipped over" relative to the other). Since the two waves are emitted in phase, this phase difference must be due to the different distances the waves travel to reach the point of destructive interference. For their phase difference to equal  $\pi$ , their path length difference must be a half-integer multiple of their wavelength.

With the top source labeled #1 and the bottom source #2:

pic 11

For destructive interference,

pic 12,13,14,15

The source could have any of the above (or many other) wavelengths.

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## 04. Analysis Tools 2

1. Double Slit Interference
2. Thin Film Interference

### Double Slit Interference

Monochromatic green light of wavelength  $550 \text{ nm}$  illuminates two parallel narrow slits  $d = 8.00$  apart. If the interference pattern is projected on a screen  $D = 5.0 \text{ m}$  from the slits, where are the bright fringes (constructive maximas) located on the screen?

pic 1

Constructive interference occurs when the two waves have a relative phase difference of  $0^\circ$ . Since the two waves are emitted in phase, any phase difference must be due to the different distances the waves travel to reach the various points on the screen. For their phase difference to be  $0^\circ$ , their path length difference must be an integer multiple of their wavelength.

We could proceed exactly as in the previous example and directly calculate the path length difference between the two waves that recombine on the screen. However, since the distance to the screen ( $D$ ) is much, much larger than the separation between the slits ( $d$ ) we can make use of a simplifying observation.

Let's imagine we are interested in the location on the screen indicated below.

pic 2

Although the paths from each slit toward the screen are not precisely parallel, if the drawing was to scale they would be *extremely* close to parallel. Below is a blow-up of the region near the slits:

pic 3

By adding a line perpendicular to the two parallel paths, and noting that this line makes the same angle  $\theta$  with the vertical that the paths make with the horizontal, the path length difference between these two paths is just the small distance indicated in the diagram below:

pic 4

so

pic 5

where

pic 6

from the diagram on the previous page. Although these results are only approximate, when  $D \gg d$  they are very useful.

Since we are looking for constructive interference in this example,

pic 7

These angles correspond to

pic 8

### Thin Film Interference

A camera lens with index of refraction  $1.40$  is coated with a thin transparent film of index of refraction  $1.20$  to eliminate the reflection of blue light ( $\lambda = 480 \text{ nm}$ ) normal to its surface. What is the minimum thickness of the film?

If blue light is not reflected from this coated camera lens, this means that waves that reflect from the *front* of the lens coating (the air-coating interface) and waves that reflect from the *rear* of the lens coating (the coating-lens interface) destructively interfere for  $\lambda = 481 \text{ nm}$  light. These waves may have a difference in phase for two, independent reasons.

First, when a wave reflects from a material having a larger index of refraction than the medium it is currently traveling through, it will reflect with a complete phase inversion (i.e., the wave will "flip over" when reflecting from something with a larger index).

Second, the wave reflecting from the rear of the coating will have traveled a greater distance than the wave reflecting from the front of the coating. This difference in path length will also cause a phase shift between the two waves.

For this particular example, both the front-reflecting wave and rear-reflecting wave reflect from surfaces having a larger index than they currently are traveling through, so both waves are "flipped" upon reflection. Since both waves are flipped, this effect will not influence the relative phase of the two waves.

Calling the thickness of the coating  $d$ , the rear-reflecting wave travels a distance  $2d$  further than the front-reflecting wave. Thus,

pic 9

If this distance is equal to one-half (or three-halves, etc.) wavelength, the reflected interference will be destructive.

pic 10

However, there is one last complication. The wavelength of the light *in the coating* is not equal to the wavelength of the light in air, and  $2d$  must equal one-half of the wavelength of the light in the coating for destructive reflection. The wavelength of light in a material other than vacuum is given by:

pic 11

Combining this with the previous relationship yields

pic 12

Thus, the coating must have a minimum thickness of 100 nm.

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## 05. Activities

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