

## 1.14: Legendre Polynomials

Consider the expression

$$(1 - 2rx + r^2)^{-1/2}, \quad (1.14.1)$$

in which  $|x|$  and  $|r|$  are both less than or equal to one. Expressions similar to this occur quite often in theoretical physics - for example in calculating the gravitational or electrostatic potentials of bodies of arbitrary shape. See, for example, Chapter 5, Sections 5.11 and 5.12.

Expand the expression 1.14.1 by the binomial theorem as a series of powers of  $r$ . This is straightforward, though not particularly easy, and you might expect to spend several minutes in obtaining the coefficients of the first few powers of  $r$ . You will find that the coefficient of  $r^l$  is a polynomial expression in  $x$  of degree  $l$ . Indeed the expansion takes the form

$$(1 - 2rx + r^2)^{-1/2} = P_0(x) + P_1(x)r + P_2(x)r^2 + P_3(x)r^3 \dots \quad (1.14.2)$$

The coefficients of the successive power of  $r$  are the *Legendre polynomials*; the coefficient of  $r^l$ , which is  $P_l(x)$ , is the Legendre polynomial of order  $l$ , and it is a polynomial in  $x$  including terms as high as  $x^l$ . We introduce these polynomials in this section because we shall need them in Section 1.15 on the derivation of Gaussian Quadrature.

If you have conscientiously tried to expand expression 1.14.1, you will have found that

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = \frac{1}{2}(3x^2 - 1), \quad (1.14.3)$$

though you probably gave up with exhaustion after that and didn't take it any further. If you look carefully at how you derived the first few polynomials, you may have discovered for yourself that you can obtain the next polynomial as a function of two earlier polynomials. You may even have discovered for yourself the following *recursion relation*:

$$P_{l+1} = \frac{(2l+1)xP_l - lP_{l-1}}{l+1}. \quad (1.14.4)$$

This enables us very rapidly to obtain higher order Legendre polynomials, whether numerically or in algebraic form. For example, put  $l = 1$  and show that Equation 1.12.4 results in  $P_2 = \frac{1}{2}(3x^2 - 1)$ . You will then want to calculate  $P_3$ , and then  $P_4$ , and more and more and more. Another way to generate them is from the Equation

$$P_{l+1} = \frac{1}{2^l l!} \frac{d^l}{dx^l} (x^2 - 1)^l. \quad (1.14.5)$$

Here are the first eleven Legendre polynomials:

$$\begin{aligned} P_0 &= 1 \\ P_1 &= x \\ P_2 &= \frac{1}{2}(3x^2 - 1) \\ P_3 &= \frac{1}{2}(5x^3 - 3x) \\ P_4 &= \frac{1}{8}(35x^4 - 30x^2 + 3) \\ P_5 &= \frac{1}{16}(63x^5 - 70x^3 + 15x) \\ P_6 &= \frac{1}{16}(231x^6 - 315x^4 + 105x^2 - 5) \\ P_7 &= \frac{1}{16}(429x^7 - 693x^5 + 315x^3 - 35x) \\ P_8 &= \frac{1}{128}(6435x^8 - 12012x^6 + 6930x^4 - 1260x^2 + 35) \\ P_9 &= \frac{1}{128}(12155x^9 - 25740x^7 + 18018x^5 - 4620x^3 + 315x) \\ P_{10} &= \frac{1}{256}(46189x^{10} - 109395x^8 + 90090x^6 - 30030x^4 + 3465x^2 - 63) \end{aligned} \quad (1.14.6)$$

The polynomials with argument  $\cos \theta$  are given in Section 5.11 of Chapter 5.

In what follows in the next section, we shall also want to know the roots of the Equations  $P_l = 0$  for  $l > 1$ . Inspection of the forms of these polynomials will quickly show that all odd polynomials have a root of zero, and all nonzero roots occur in

positive/negative pairs. Having read Sections 1.4 and 1.5, we shall have no difficulty in finding the roots of these Equations. The roots  $x_{l,i}$  are in the following table, which also lists certain coefficients  $c_{l,i}$ , that will be explained in Section 1.15.

Roots of  $P_l = 0$

$l$	$x_{l,i}$	$c_{l,i}$
2	$\pm 0.577\ 350\ 269\ 190$	1.000 000 000 00
3	$\pm 0.774\ 596\ 669\ 241$ $0.000\ 000\ 000\ 000$	0.555 555 555 56 0.888 888 888 89
4	$\pm 0.861\ 136\ 311\ 594$ $\pm 0.339\ 981\ 043\ 585$	0.347 854 845 14 0.652 145 154 86
5	$\pm 0.906\ 179\ 845\ 939$ $\pm 0.538\ 469\ 310\ 106$ $0.000\ 000\ 000\ 000$	0.236 926 885 06 0.478 628 670 50 0.568 888 888 89
6	$\pm 0.932\ 469\ 514\ 203$ $\pm 0.661\ 209\ 386\ 466$ $\pm 0.238\ 619\ 186\ 083$	0.171 324 492 38 0.360 761 573 05 0.467 913 934 57
7	$\pm 0.949\ 107\ 912\ 343$ $\pm 0.741\ 531\ 185\ 599$ $\pm 0.405\ 845\ 151\ 377$ $0.000\ 000\ 000\ 000$	0.129 484 966 17 0.279 705 391 49 0.381 830 050 50 0.417 959 183 68
8	$\pm 0.960\ 289\ 856\ 498$ $\pm 0.796\ 666\ 477\ 414$ $\pm 0.525\ 532\ 409\ 916$ $\pm 0.183\ 434\ 642\ 496$	0.101 228 536 29 0.222 381 034 45 0.313 706 645 88 0.362 683 783 38
9	$\pm 0.968\ 160\ 239\ 508$ $\pm 0.836\ 031\ 107\ 327$ $\pm 0.613\ 371\ 432\ 701$ $\pm 0.324\ 253\ 423\ 404$ $0.000\ 000\ 000\ 000$	0.081 274 388 36 0.180 648 160 69 0.260 610 696 40 0.312 347 077 04 0.330 239 355 00
10	$\pm 0.973\ 906\ 528\ 517$ $\pm 0.865\ 063\ 366\ 689$ $\pm 0.679\ 409\ 568\ 299$ $\pm 0.433\ 395\ 394\ 129$ $\pm 0.148\ 874\ 338\ 982$	0.066 671 343 99 0.149 451 349 64 0.219 086 362 24 0.269 266 719 47 0.295 524 224 66
11	$\pm 0.978\ 228\ 658\ 146$ $\pm 0.887\ 062\ 599\ 768$ $\pm 0.730\ 152\ 005\ 574$ $\pm 0.519\ 096\ 129\ 207$ $+0.269\ 543\ 155\ 952$	0.055 668 567 12 0.125 580 369 46 0.186 290 210 93 0.233 193 764 59 0.262 804 544 51

	0.000 000 000 000	0.272 925 086 78	
12	±0.981 560 634 247 ±0.904 117 256 370 ±0.769 902 674 194 ±0.587 317 954 287 ±0.367 831 498 998 ±0.125 233 408 511	0.047 175 336 39 0.106 939 325 99 0.160 078 328 54 0.203 167 426 72 0.233 492 536 54 0.249 147 045 81	(1.14.1)
13	±0.984 183 054 719 ±0.917 598 399 223 ±0.801 578 090 733 ±0.642 349 339 440 ±0.448 492 751 036 ±0.230 458 315 955 0.000 000 000 000	0.040 484 004 77 0.092 121 499 84 0.138 873 510 22 0.178 145 980 76 0.207 816 047 54 0.226 283 180 26 0.232 551 553 23	
14	±0.986 283 808 697 ±0.928 434 883 664 ±0.827 201 315 070 ±0.687 292 904 812 ±0.515 248 636 358 ±0.319 112 368 928 ±0.108 054 948 707	0.035 119 460 33 0.080 158 087 16 0.121 518 570 69 0.157 203 167 16 0.185 538 397 48 0.205 198 463 72 0.215 263 853 46	
15	±0.987 992 518 020 ±0.937 273 392 401 ±0.848 206 583 410 ±0.724 417 731 360 ±0.570 972 172 609 ±0.394 151 347 078 ±0.201 194 093 997 0.000 000 000 000	0.030 753 242 00 0.070 366 047 49 0.107 159 220 47 0.139 570 677 93 0.166 269 205 82 0.186 161 000 02 0.198 431 485 33 0.202 578 241 92	
16	±0.989 400 934 992 ±0.944 575 023 073 ±0.865 631 202 388 ±0.755 404 408 355 ±0.617 876 244 403 ±0.458 016 777 657 ±0.281 603 550 779 ±0.095 012 509 838	0.027 152 459 41 0.062 253 523 94 0.095 158 511 68 0.124 628 971 26 0.149 595 988 82 0.169 156 519 39 0.182 603 415 04 0.189 450 610 46	
17	±0.990 575 475 315 ±0.950 675 521 769 ±0.880 239 153 727 ±0.781 514 003 897 ±0.657 671 159 217	0.024 148 302 87 0.055 459 529 38 0.085 036 148 32 0.111 883 847 19 0.135 136 368 47	

$\pm 0.512\ 690\ 537\ 086\ 0.154\ 045\ 761\ 08$   
 $\pm 0.351\ 231\ 763\ 454\ 0.168\ 004\ 102\ 16$   
 $\pm 0.178\ 484\ 181\ 496\ 0.176\ 562\ 705\ 37$   
 $0.000\ 000\ 000\ 000\ 0.179\ 446\ 470\ 35$

For interest, I draw graphs of the Legendre polynomials in figures I.7 and I.8.

Figure I.7. Legendre polynomials for even I

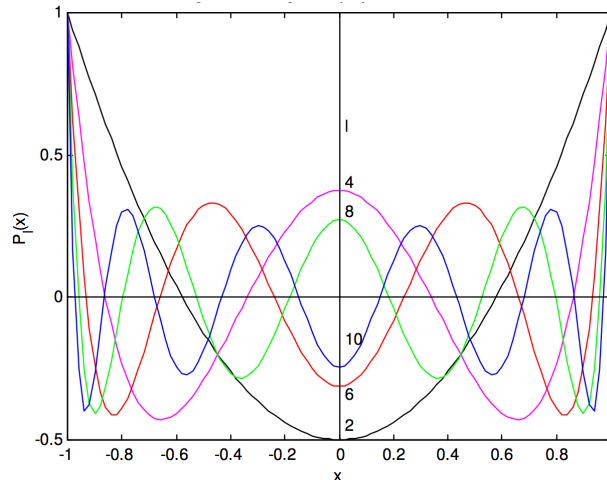
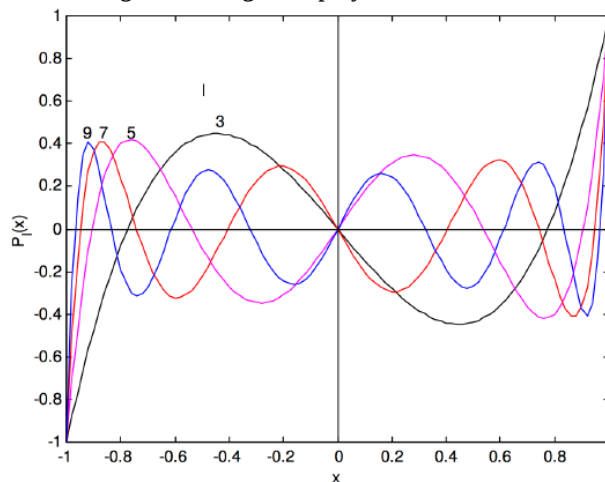


Figure I.8. Legendre polynomials for odd I



For further interest, it should be easy to verify, by substitution, that the Legendre polynomials are solutions of the differential Equation

$$(1 - x^2)y'' - 2xy' + l(l+1)y = 0. \quad (1.14.7)$$

The Legendre polynomials are solutions of this and related Equations that appear in the study of the vibrations of a solid sphere (spherical harmonics) and in the solution of the Schrödinger Equation for hydrogen-like atoms, and they play a large role in quantum mechanics.

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