

## 13.4: Pressure, Temperature, and RMS Speed

### Learning Objectives

By the end of this section, you will be able to:

- Explain the relations between microscopic and macroscopic quantities in a gas
- Solve problems involving mixtures of gases
- Solve problems involving the distance and time between a gas molecule's collisions

We have examined pressure and temperature based on their macroscopic definitions. Pressure is the force divided by the area on which the force is exerted, and temperature is measured with a thermometer. We can gain a better understanding of pressure and temperature from the **kinetic theory of gases**, the theory that relates the macroscopic properties of gases to the motion of the molecules they consist of. First, we make two assumptions about molecules in an ideal gas.

1. There is a very large number  $N$  of molecules, all identical and each having mass  $m$ .
2. The molecules obey Newton's laws and are in continuous motion, which is random and isotropic, that is, the same in all directions.

To derive the ideal gas law and the connection between microscopic quantities such as the energy of a typical molecule and macroscopic quantities such as temperature, we analyze a sample of an ideal gas in a rigid container, about which we make two further assumptions:

3. The molecules are much smaller than the average distance between them, so their total volume is much less than that of their container (which has volume  $V$ ). In other words, we take the Van der Waals constant  $b$ , the volume of a mole of gas molecules, to be negligible compared to the volume of a mole of gas in the container.
4. The molecules make perfectly elastic collisions with the walls of the container and with each other. Other forces on them, including gravity and the attractions represented by the Van der Waals constant  $a$ , are negligible (as is necessary for the assumption of isotropy).

The collisions between molecules do not appear in the derivation of the ideal gas law. They do not disturb the derivation either, since collisions between molecules moving with random velocities give new random velocities. Furthermore, if the velocities of gas molecules in a container are initially not random and isotropic, molecular collisions are what make them random and isotropic.

We make still further assumptions that simplify the calculations but do not affect the result. First, we let the container be a rectangular box. Second, we begin by considering **monatomic** gases, those whose molecules consist of single atoms, such as helium. Then, we can assume that the atoms have no energy except their translational kinetic energy; for instance, they have neither rotational nor vibrational energy. (Later, we discuss the validity of this assumption for real monatomic gases and dispense with it to consider diatomic and polyatomic gases.)

Figure 13.4.1 shows a collision of a gas molecule with the wall of a container, so that it exerts a force on the wall (by Newton's third law). These collisions are the source of pressure in a gas. As the number of molecules increases, the number of collisions, and thus the pressure, increases. Similarly, if the average velocity of the molecules is higher, the gas pressure is higher.

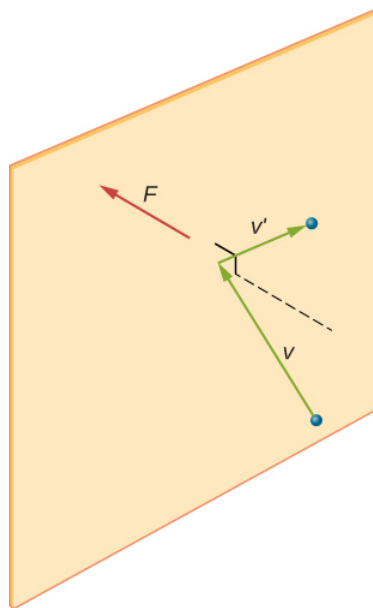


Figure 13.4.1: When a molecule collides with a rigid wall, the component of its momentum perpendicular to the wall is reversed. A force is thus exerted on the wall, creating pressure.

In a sample of gas in a container, the randomness of the molecular motion causes the number of collisions of molecules with any part of the wall in a given time to fluctuate. However, because a huge number of molecules collide with the wall in a short time, the number of collisions on the scales of time and space we measure fluctuates by only a tiny, usually unobservable fraction from the average. We can compare this situation to that of a casino, where the outcomes of the bets are random and the casino's takings fluctuate by the minute and the hour. However, over long times such as a year, the casino's takings are very close to the averages expected from the odds. A tank of gas has enormously more molecules than a casino has bettors in a year, and the molecules make enormously more collisions in a second than a casino has bets.

A calculation of the average force exerted by molecules on the walls of the box leads us to the ideal gas law and to the connection between temperature and molecular kinetic energy. (In fact, we will take two averages: one over time to get the average force exerted by one molecule with a given velocity, and then another average over molecules with different velocities.) This approach was developed by Daniel **Bernoulli** (1700–1782), who is best known in physics for his work on fluid flow (hydrodynamics). Remarkably, Bernoulli did this work before Dalton established the view of matter as consisting of atoms.

Figure 13.4.2 shows a container full of gas and an expanded view of an elastic collision of a gas molecule with a wall of the container, broken down into components. We have assumed that a molecule is small compared with the separation of molecules in the gas, and that its interaction with other molecules can be ignored. Under these conditions, the ideal gas law is experimentally valid. Because we have also assumed the wall is rigid and the particles are points, the collision is elastic (by conservation of energy—there's nowhere for a particle's kinetic energy to go). Therefore, the molecule's kinetic energy remains constant, and hence, its

speed and the magnitude of its momentum remain constant as well. This assumption is not always valid, but the results in the rest of this module are also obtained in models that let the molecules exchange energy and momentum with the wall.

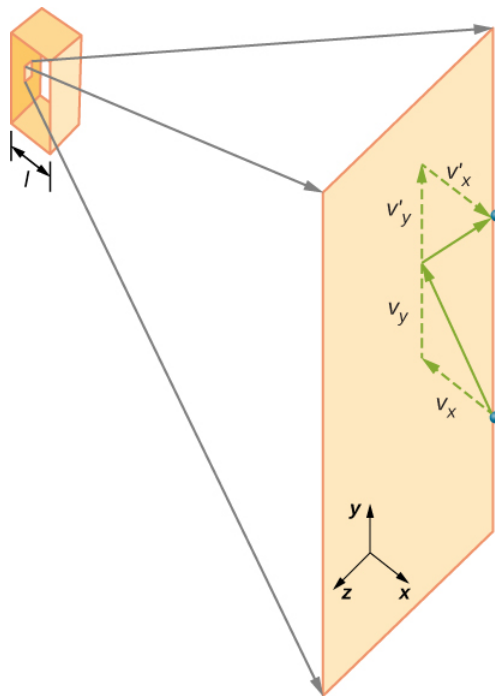


Figure 13.4.2: Gas in a box exerts an outward pressure on its walls. A molecule colliding with a rigid wall has its velocity and momentum in the  $x$ -direction reversed. This direction is perpendicular to the wall. The components of its velocity momentum in the  $y$ - and  $z$ -directions are not changed, which means there is no force parallel to the wall.

If the molecule's velocity changes in the  $x$ -direction, its momentum changes from  $-mv_x$  to  $+mv_x$ . Thus, its change in momentum is  $\Delta mv = +mv_x - (-mv_x) = 2mv_x$ . According to the impulse-momentum theorem given in the chapter on linear momentum and collisions, the force exerted on the  $i$ th molecule, where  $i$  labels the molecules from 1 to  $N$ , is given by

$$F_i = \frac{\Delta P_i}{\Delta t} = \frac{2mv_{ix}}{\Delta t}.$$

(In this equation alone,  $\mathbf{p}$  represents momentum, not pressure.) There is no force between the wall and the molecule except while the molecule is touching the wall. During the short time of the collision, the force between the molecule and wall is relatively large, but that is not the force we are looking for. We are looking for the average force, so we take  $\Delta t$  to be the average time between collisions of the given molecule with this wall, which is the time in which we expect to find one collision. Let  $l$  represent the length of the box in the  $x$ -direction. Then  $\Delta t$  is the time the molecule would take to go across the box and back, a distance  $2l$ , at a speed of  $v_x$ . Thus  $\delta t = 2l/v_x$ , and the expression for the force becomes

$$F_i = \frac{2mv_{ix}}{2l/v_{ix}} = \frac{mv_{ix}^2}{l}.$$

This force is due to **one** molecule. To find the total force on the wall,  $\mathbf{F}$ , we need to add the contributions of all  $N$  molecules:

$$F = \sum_{i=1}^N F_i = \sum_{i=1}^N \frac{mv_{ix}^2}{l} = \frac{m}{l} \sum_{i=1}^N v_{ix}^2.$$

We now use the definition of the average, which we denote with a bar, to find the force:

$$F = N \frac{m}{l} \left( \frac{1}{N} \sum_{i=1}^N v_{ix}^2 \right) = N \frac{mv_{ix}^2}{l}.$$

We want the force in terms of the speed  $\mathbf{v}$ , rather than the  $x$ -component of the velocity. Note that the total velocity squared is the sum of the squares of its components, so that

$$\bar{v}^2 = \bar{v}_x^2 + \bar{v}_y^2 + \bar{v}_z^2.$$

With the assumption of isotropy, the three averages on the right side are equal, so

$$\bar{v}^2 = 3\bar{v}_{ix}^2.$$

Substituting this into the expression for  $\mathbf{F}$  gives

$$F = N \frac{m\bar{v}^2}{3l}.$$

The pressure is  $\mathbf{F}/\mathbf{A}$ , so we obtain

$$p = \frac{F}{A} = N \frac{m\bar{v}^2}{3Al} = \frac{Nm\bar{v}^2}{3V},$$

where we used  $V = Al$  for the volume. This gives the important result

$$pV = \frac{1}{3}Nm\bar{v}^2.$$

Combining this equation with  $pV = Nk_B T$  gives

$$\frac{1}{3}Nm\bar{v}^2 = Nk_B T.$$

We can get the average kinetic energy of a molecule,  $\frac{1}{2}m\bar{v}^2$ , from the left-hand side of the equation by dividing out  $\mathbf{N}$  and multiplying by  $3/2$ .

#### Average Kinetic Energy per Molecule

The average kinetic energy of a molecule is directly proportional to its absolute temperature:

$$\bar{K} = \frac{1}{2}m\bar{v}^2 = \frac{3}{2}k_B T.$$

The equation  $\bar{K} = \frac{3}{2}k_B T$  is the average kinetic energy per molecule. Note in particular that nothing in this equation depends on the molecular mass (or any other property) of the gas, the pressure, or anything but the temperature. If samples of helium and xenon gas, with very different molecular masses, are at the same temperature, the molecules have the same average kinetic energy.

The **internal energy** of a thermodynamic system is the sum of the mechanical energies of all of the molecules in it. We can now give an equation for the internal energy of a monatomic ideal gas. In such a gas, the molecules' only energy is their translational kinetic energy. Therefore, denoting the internal energy by  $E_{int}$  we simply have  $E_{int} = N\bar{K}$ , or

$$E_{int} = \frac{3}{2}Nk_B T.$$

Often we would like to use this equation in terms of moles:

$$E_{int} = \frac{3}{2}nRT.$$

We can solve  $\bar{K} = \frac{1}{2}m\bar{v}^2 = \frac{3}{2}k_B T$  for a typical speed of a molecule in an ideal gas in terms of temperature to determine what is known as the **root-mean-square (rms) speed** of a molecule.

#### RMS Speed of a Molecule

The root-mean-square (rms) speed of a molecule, or the square root of the average of the square of the speed  $\bar{v}^2$ , is

$$v_{rms} = \sqrt{\bar{v}^2} = \sqrt{\frac{3k_B T}{m}}.$$

The rms speed is not the average or the most likely speed of molecules, as we will see in [Distribution of Molecular Speeds](#), but it provides an easily calculated estimate of the molecules' speed that is related to their kinetic energy. Again we can write this equation in terms of the gas constant **R** and the molar mass **M** in kg/mol:

$$v_{rms} = \sqrt{\frac{3RT}{M}}.$$

We digress for a moment to answer a question that may have occurred to you: When we apply the model to atoms instead of theoretical point particles, does rotational kinetic energy change our results? To answer this question, we have to appeal to quantum mechanics. In quantum mechanics, rotational kinetic energy cannot take on just any value; it's limited to a discrete set of values, and the smallest value is inversely proportional to the rotational inertia. The rotational inertia of an atom is tiny because almost all of its mass is in the nucleus, which typically has a radius less than  $10^{-14}m$ . Thus the minimum rotational energy of an atom is much more than  $\frac{1}{2}k_B T$  for any attainable temperature, and the energy available is not enough to make an atom rotate. We will return to this point when discussing diatomic and polyatomic gases in the next section.

### ✓ Example 13.4.1: Calculating Kinetic Energy and Speed of a Gas Molecule

- What is the average kinetic energy of a gas molecule at  $20.0^\circ C$  (room temperature)?
- Find the rms speed of a nitrogen molecule ( $N_2$ ) at this temperature.

#### Strategy

(a) The known in the equation for the average kinetic energy is the temperature:

$$\bar{K} = \frac{1}{2}m\bar{v}^2 = \frac{3}{2}k_B T.$$

Before substituting values into this equation, we must convert the given temperature into kelvin:  $T = (20.0 + 273) K = 293 K$ . We can find the rms speed of a nitrogen molecule by using the equation

$$v_{rms} = \sqrt{\bar{v}^2} = \sqrt{\frac{3k_B T}{m}},$$

but we must first find the mass of a nitrogen molecule. Obtaining the molar mass of nitrogen  $N_2$  from the periodic table, we find

$$m = \frac{M}{N_A} = \frac{2(14.0067 \times 10^{-3} kg/mol)}{6.02 \times 10^{23} mol^{-1}} = 4.65 \times 10^{-26} kg.$$

#### Solution

- The temperature alone is sufficient for us to find the average translational kinetic energy. Substituting the temperature into the translational kinetic energy equation gives

$$\bar{K} = \frac{3}{2}k_B T = \frac{3}{2}(1.38 \times 10^{-23} J/K)(293 K) = 6.07 \times 10^{-21} J.$$

- Substituting this mass and the value for  $k_B$  into the equation for  $v_{rms}$  yields

$$v_{rms} = \sqrt{\frac{3k_B T}{m}} = \sqrt{\frac{3(1.38 \times 10^{-23} J/K)(293 K)}{4.65 \times 10^{-26} kg}} = 511 m/s.$$

#### Significance

Note that the average kinetic energy of the molecule is independent of the type of molecule. The average translational kinetic energy depends only on absolute temperature. The kinetic energy is very small compared to macroscopic energies, so that we do not feel when an air molecule is hitting our skin. On the other hand, it is much greater than the typical difference in gravitational potential energy when a molecule moves from, say, the top to the bottom of a room, so our neglect of gravitation is justified in typical real-world situations. The rms speed of the nitrogen molecule is surprisingly large. These large molecular velocities do not yield macroscopic movement of air, since the molecules move in all directions with equal likelihood. The **mean free path** (the distance a molecule moves on average between collisions, discussed a bit later in this section) of

molecules in air is very small, so the molecules move rapidly but do not get very far in a second. The high value for rms speed is reflected in the speed of sound, which is about 340 m/s at room temperature. The higher the rms speed of air molecules, the faster sound vibrations can be transferred through the air. The speed of sound increases with temperature and is greater in gases with small molecular masses, such as helium (see Figure 13.4.3).

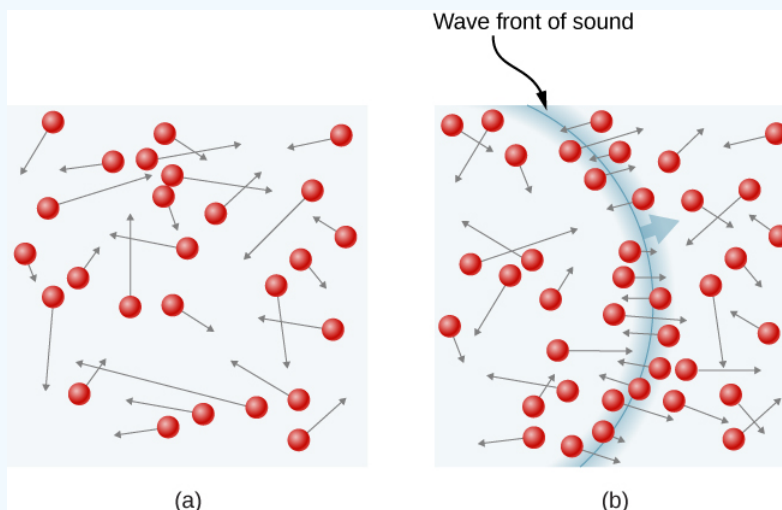


Figure 13.4.3: (a) In an ordinary gas, so many molecules move so fast that they collide billions of times every second. Individual molecules do not move very far in a small amount of time, but disturbances like sound waves are transmitted at speeds related to the molecular speeds.

#### ✓ Example 13.4.2: Calculating Temperature: Escape Velocity of Helium Atoms

To escape Earth's gravity, an object near the top of the atmosphere (at an altitude of 100 km) must travel away from Earth at 11.1 km/s. This speed is called the **escape velocity**. At what temperature would helium atoms have an rms speed equal to the escape velocity?

##### Strategy

Identify the knowns and unknowns and determine which equations to use to solve the problem.

##### Solution

1. Identify the knowns:  $v$  is the escape velocity, 11.1 km/s.
2. Identify the unknowns: We need to solve for temperature,  $T$ . We also need to solve for the mass  $m$  of the helium atom.
3. Determine which equations are needed.
  - To get the mass  $m$  of the helium atom, we can use information from the periodic table:

$$m = \frac{M}{N_A},$$

- To solve for temperature  $T$ , we can rearrange

$$\frac{1}{2} m \bar{v}^2 = \frac{3}{2} k_B T$$

to yield

$$T = \frac{m \bar{v}^2}{3 k_B}.$$

4. Substitute the known values into the equations and solve for the unknowns,

$$m = \frac{M}{N_A} = \frac{4.0026 \times 10^{-3} \text{ kg/mol}}{6.02 \times 10^{23} \text{ mol}} = 6.65 \times 10^{-27} \text{ kg}$$

and

$$T = \frac{(6.65 \times 10^{-27} \text{ kg})(11.1 \times 10^3 \text{ m/s})^2}{3(1.38 \times 10^{-23} \text{ J/K})} = 1.98 \times 10^4 \text{ K}.$$

### Significance

This temperature is much higher than atmospheric temperature, which is approximately 250 K ( $-25^\circ\text{C}$  or  $-10^\circ\text{F}$ ) at high elevation. Very few helium atoms are left in the atmosphere, but many were present when the atmosphere was formed, and more are always being created by radioactive decay (see the chapter on nuclear physics). The reason for the loss of helium atoms is that a small number of helium atoms have speeds higher than Earth's escape velocity even at normal temperatures. The speed of a helium atom changes from one collision to the next, so that at any instant, there is a small but nonzero chance that the atom's speed is greater than the escape velocity. The chance is high enough that over the lifetime of Earth, almost all the helium atoms that have been in the atmosphere have reached escape velocity at high altitudes and escaped from Earth's gravitational pull. Heavier molecules, such as oxygen, nitrogen, and water, have smaller rms speeds, and so it is much less likely that any of them will have speeds greater than the escape velocity. In fact, the likelihood is so small that billions of years are required to lose significant amounts of heavier molecules from the atmosphere. Figure 13.4.4 shows the effect of a lack of an atmosphere on the Moon. Because the gravitational pull of the Moon is much weaker, it has lost almost its entire atmosphere. The atmospheres of Earth and other bodies are compared in this chapter's exercises.



Figure 13.4.4: This photograph of Apollo 17 Commander Eugene Cernan driving the lunar rover on the Moon in 1972 looks as though it was taken at night with a large spotlight. In fact, the light is coming from the Sun. Because the acceleration due to gravity on the Moon is so low (about 1/6 that of Earth), the Moon's escape velocity is much smaller. As a result, gas molecules escape very easily from the Moon, leaving it with virtually no atmosphere. Even during the daytime, the sky is black because there is no gas to scatter sunlight. (credit: Harrison H. Schmitt/NASA)

### ? Exercise 13.4.2

If you consider a very small object, such as a grain of pollen, in a gas, then the number of molecules striking its surface would also be relatively small. Would you expect the grain of pollen to experience any fluctuations in pressure due to statistical fluctuations in the number of gas molecules striking it in a given amount of time?

### Answer

Yes. Such fluctuations actually occur for a body of any size in a gas, but since the numbers of molecules are immense for macroscopic bodies, the fluctuations are a tiny percentage of the number of collisions, and the averages spoken of in this section vary imperceptibly. Roughly speaking, the fluctuations are inversely proportional to the square root of the number of collisions, so for small bodies, they can become significant. This was actually observed in the nineteenth century for pollen grains in water and is known as Brownian motion.

### 13.4.1 Vapor Pressure, Partial Pressure, and Dalton's Law

The pressure a gas would create if it occupied the total volume available is called the gas's **partial pressure**. If two or more gases are mixed, they will come to thermal equilibrium as a result of collisions between molecules; the process is analogous to heat conduction as described in the chapter on temperature and heat. As we have seen from kinetic theory, when the gases have the same temperature, their molecules have the same average kinetic energy. Thus, each gas obeys the ideal gas law separately and exerts the same pressure on the walls of a container that it would if it were alone. Therefore, in a mixture of gases, **the total pressure is the sum of partial pressures of the component gases**, assuming ideal gas behavior and no chemical reactions between the components. This law is known as **Dalton's law of partial pressures**, after the English scientist John **Dalton** (1766–1844) who proposed it. Dalton's law is consistent with the fact that pressures add according to Pascal's principle.

In a mixture of ideal gases in thermal equilibrium, the number of molecules of each gas is proportional to its partial pressure. This result follows from applying the ideal gas law to each in the form  $p/n = RT/V$ . Because the right-hand side is the same for any gas at a given temperature in a container of a given volume, the left-hand side is the same as well.

- Partial pressure is the pressure a gas would create if it existed alone.
- Dalton's law states that the total pressure is the sum of the partial pressures of all of the gases present.
- For any two gases (labeled 1 and 2) in equilibrium in a container,  $\frac{p_1}{n_1} = \frac{p_2}{n_2}$ .

An important application of partial pressure is that, in chemistry, it functions as the concentration of a gas in determining the rate of a reaction. Here, we mention only that the partial pressure of oxygen in a person's lungs is crucial to life and health. Breathing air that has a partial pressure of oxygen below 0.16 atm can impair coordination and judgment, particularly in people not acclimated to a high elevation. Lower partial pressures of  $O_2$  have more serious effects; partial pressures below 0.06 atm can be quickly fatal, and permanent damage is likely even if the person is rescued. However, the sensation of needing to breathe, as when holding one's breath, is caused much more by high concentrations of carbon dioxide in the blood than by low concentrations of oxygen. Thus, if a small room or closet is filled with air having a low concentration of oxygen, perhaps because a leaking cylinder of some compressed gas is stored there, a person will not feel any "choking" sensation and may go into convulsions or lose consciousness without noticing anything wrong. Safety engineers give considerable attention to this danger.

Another important application of partial pressure is **vapor pressure**, which is the partial pressure of a vapor at which it is in equilibrium with the liquid (or solid, in the case of sublimation) phase of the same substance. At any temperature, the partial pressure of the water in the air cannot exceed the vapor pressure of the water at that temperature, because whenever the partial pressure reaches the vapor pressure, water condenses out of the air. Dew is an example of this condensation. The temperature at which condensation occurs for a sample of air is called the **dew point**. It is easily measured by slowly cooling a metal ball; the dew point is the temperature at which condensation first appears on the ball.

The vapor pressures of water at some temperatures of interest for meteorology are given in Table 13.4.1.

Table 13.4.1: Vapor Pressure of Water at Various Temperatures

$T(^{\circ}C)$	Vapor Pressure (Pa)
0	610.5
3	757.9
5	872.3
8	1073
10	1228
13	1497
15	1705
18	2063
20	2338
23	2809
25	3167



$T(^{\circ}C)$	Vapor Pressure (Pa)
30	4243
35	5623
40	7376

The **relative humidity** (R.H.) at a temperature  $T$  is defined by

$$R. H. = \frac{\text{Partial pressure of water vapor at } T}{\text{Vapor pressure of water at } T} \times 100\%.$$

A relative humidity of 100% means that the partial pressure of water is equal to the vapor pressure; in other words, the air is saturated with water.

#### ✓ Example 13.4.3: Calculating Relative Humidity

What is the relative humidity when the air temperature is  $25^{\circ}C$  and the dew point is  $15^{\circ}C$ ?

##### Strategy

We simply look up the vapor pressure at the given temperature and that at the dew point and find the ratio.

##### Solution

$$R. H. = \frac{\text{Partial pressure of water vapor at } 15^{\circ}C}{\text{Vapor pressure of water at } 25^{\circ}C} \times 100\% = \frac{1705 \text{ Pa}}{3167 \text{ Pa}} \times 100\% = 53.8\%.$$

##### Significance

R.H. is important to our comfort. The value of 53.8% is within the range of 40% to 60% recommended for comfort indoors.

As noted in the chapter on temperature and heat, the temperature seldom falls below the dew point, because when it reaches the dew point or frost point, water condenses and releases a relatively large amount of latent heat of vaporization.

### 13.4.2 Mean Free Path and Mean Free Time

We now consider collisions explicitly. The usual first step (which is all we'll take) is to calculate the mean free path,  $\lambda$ , the average distance a molecule travels between collisions with other molecules, and the **mean free time**  $\tau$ , the average time between the collisions of a molecule. If we assume all the molecules are spheres with a radius  $r$  then a molecule will collide with another if their centers are within a distance  $2r$  of each other. For a given particle, we say that the area of a circle with that radius,  $4\pi r^2$ , is the "cross-section" for collisions. As the particle moves, it traces a cylinder with that cross-sectional area. The mean free path is the length  $\lambda$  such that the expected number of other molecules in a cylinder of length  $\lambda$  and cross-section  $4\pi r^2$  is 1. If we temporarily ignore the motion of the molecules other than the one we're looking at, the expected number is the number density of molecules,  $N/V$ , times the volume, and the volume is  $4\pi r^2 \lambda$ , so we have  $(N/V)4\pi r^2 \lambda = 1$ , or

$$\lambda = \frac{V}{4\pi r^2 N}.$$

Taking the motion of all the molecules into account makes the calculation much harder, but the only change is a factor of  $\sqrt{2}$ . The result is

$$\lambda = \frac{V}{4\sqrt{2}\pi r^2 N}.$$

In an ideal gas, we can substitute  $V/N = k_B T/p$  to obtain

$$\lambda = \frac{k_B T}{4\sqrt{2}\pi r^2 p}.$$

The **mean free time**  $\tau$  is simply the mean free path divided by a typical speed, and the usual choice is the rms speed. Then

$$\tau = \frac{k_B T}{4\sqrt{2}\pi r^2 p v_{rms}}$$

### ✓ Calculating Mean Free Time

Find the mean free time for argon atoms ( $M = 39.9 \text{ g/mol}$ ) at a temperature of  $0^\circ\text{C}$  and a pressure of  $1.00 \text{ atm}$ . Take the radius of an argon atom to be  $1.70 \times 10^{-10} \text{ m}$ .

**Solution**

1. Identify the knowns and convert into SI units. We know the molar mass is  $0.0399 \text{ kg/mol}$ , the temperature is  $273 \text{ K}$ , the pressure is  $1.01 \times 10^5 \text{ Pa}$ , and the radius is  $1.70 \times 10^{-10} \text{ m}$ .
2. Find the rms speed:  $v_{rms} = \sqrt{\frac{3RT}{M}} = 413 \text{ m/s}$ .
3. Substitute into the equation for the mean free time:

$$\tau = \frac{k_B T}{4\sqrt{2}\pi r^2 p v_{rms}} = \frac{(1.38 \times 10^{-23} \text{ J/K})(273 \text{ K})}{4\sqrt{2}\pi (1.70 \times 10^{-10} \text{ m})^2 (1.01 \times 10^5 \text{ Pa})(413 \text{ m/s})} = 1.76 \times 10^{-10} \text{ s}.$$

### Significance

We can hardly compare this result with our intuition about gas molecules, but it gives us a picture of molecules colliding with extremely high frequency.

### ? Exercise 13.4.4

Which has a longer mean free path, liquid water or water vapor in the air?

### Answer

In a liquid, the molecules are very close together, constantly colliding with one another. For a gas to be nearly ideal, as air is under ordinary conditions, the molecules must be very far apart. Therefore the mean free path is much longer in the air.

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