

## 19.1: Mass Suspended by Two Ropes

As a typical example of a problem in statics, consider the situation shown in Fig. 19.1.1(a). A block of mass  $m$  is suspended by a wire, and the upper end of the wire is attached to two more ropes or wires that connect to the ceiling. Each of the three ropes is under tension; the tensions are labeled  $T_1$ ,  $T_2$ , and  $T_3$ .

To begin the analysis of this situation, it is often helpful to draw a free-body diagram for each body in the problem. A free-body diagram shows all the forces acting on the body, and helps clarify your thinking when doing the analysis. For this problem, there are two bodies present: the block and the knot. Fig. 19.1.1(b) is a free-body diagram for the block, and Fig. 19.1.1(c) is a free-body diagram for the knot.

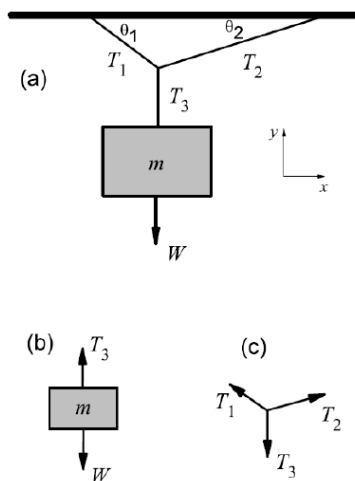


Figure 19.1.1: A block suspended from the ceiling by ropes. (a) Diagram of the situation. The block of mass  $m$  is suspended by a rope; the upper end of the rope is attached to two other ropes that are attached to the ceiling. (b) Free-body diagram for the block. (c) Free-body diagram for the knot.

Now let's begin the analysis; our goal will be to determine the three tensions  $T_1$ ,  $T_2$ , and  $T_3$ , given the mass  $m$  and two angles  $\theta_1$  and  $\theta_2$ . First, let's look at the free-body diagram for the block (Fig. 19.1.1(b)). For the block, the tension and weight vectors are given by

$$\mathbf{T}_3 = T_3 \mathbf{j} \quad (19.1.1)$$

$$\mathbf{W} = -mg \mathbf{j} \quad (19.1.2)$$

(Note the  $x$  and  $y$  directions indicated in Fig. 19.1.1(a).) Now let's apply Newton's second law in both the  $x$  and  $y$  directions, noting that  $F = ma = 0$  in this case:

$$x : \sum F_x = ma_x \Rightarrow 0 = 0 \quad (19.1.3)$$

$$y : \sum F_y = ma_y \Rightarrow T_3 - mg = 0 \quad (19.1.4)$$

Both right-hand sides are zero because the acceleration of the block is zero. The  $x$  equation (Eq. 19.1.3) yields a tautology  $0 = 0$ , which gives us no information. The  $y$  equation (Eq. 19.1.4) tells us  $T_3 = mg$ , so we've just found tension  $T_3$ .

We can find the other two tensions ( $T_1$  and  $T_2$ ) by analyzing the other body: the knot (Fig. 19.1.1(c)). For the knot, the three tension vectors are given by

$$\mathbf{T}_1 = -T_1 \cos \theta_1 \mathbf{i} + T_1 \sin \theta_1 \mathbf{j} \quad (19.1.5)$$

$$\mathbf{T}_2 = T_2 \cos \theta_2 \mathbf{i} + T_2 \sin \theta_2 \mathbf{j} \quad (19.1.6)$$

$$\mathbf{T}_3 = -T_3 \mathbf{j} \quad (19.1.7)$$

Now let's apply Newton's second law ( $F = ma = 0$ ) individually to the  $x$  and  $y$  components:

$$x : \sum F_x = ma_x \Rightarrow -T_1 \cos \theta_1 + T_2 \cos \theta_2 = 0 \quad (19.1.8)$$

$$y : \sum F_y = ma_y \Rightarrow T_1 \sin \theta_1 + T_2 \sin \theta_2 - T_3 = 0 \quad (19.1.9)$$

Again both right-hand sides are zero because the knot is not accelerating. Since  $T_3$  is already known, this gives two simultaneous equations in the two unknown tensions  $T_1$  and  $T_2$ . One method for solving this system of equations is to write the equations in matrix form:

$$\begin{pmatrix} -\cos \theta_1 \cos \theta_2 \\ \sin \theta_1 \sin \theta_2 \end{pmatrix} \begin{pmatrix} T_1 \\ T_2 \end{pmatrix} = \begin{pmatrix} 0 \\ T_3 \end{pmatrix} \quad (19.1.10)$$

Now multiplying both sides on the left by the inverse of the  $2 \times 2$  matrix, we have

$$\begin{pmatrix} T_1 \\ T_2 \end{pmatrix} = \begin{pmatrix} -\cos \theta_1 \cos \theta_2 \\ \sin \theta_1 \sin \theta_2 \end{pmatrix}^{-1} \begin{pmatrix} 0 \\ T_3 \end{pmatrix} \quad (19.1.11)$$

Since the tension  $T_3$  and the angles  $\theta_1$  and  $\theta_2$  are all known, this gives the two unknown tensions  $T_1$  and  $T_2$ .

We can further simplify this by computing the matrix inverse explicitly. The determinant of the  $2 \times 2$  matrix is ([Appendix 63.17](#))

$$\det \begin{pmatrix} -\cos \theta_1 \cos \theta_2 \\ \sin \theta_1 \sin \theta_2 \end{pmatrix} = -\cos \theta_1 \sin \theta_2 - \sin \theta_1 \cos \theta_2 \quad (19.1.12)$$

$$= -\sin(\theta_1 + \theta_2) \quad (19.1.13)$$

and the matrix of cofactors is

$$\text{cof} \begin{pmatrix} -\cos \theta_1 \cos \theta_2 \\ \sin \theta_1 \sin \theta_2 \end{pmatrix} = \begin{pmatrix} \sin \theta_2 - \sin \theta_1 \\ -\cos \theta_2 - \cos \theta_1 \end{pmatrix} \quad (19.1.14)$$

Hence the matrix inverse, which is the transposed matrix of cofactors divided by the determinant, is

$$\begin{pmatrix} -\cos \theta_1 \cos \theta_2 \\ \sin \theta_1 \sin \theta_2 \end{pmatrix}^{-1} = -\frac{1}{\sin(\theta_1 + \theta_2)} \begin{pmatrix} \sin \theta_2 - \cos \theta_2 \\ -\sin \theta_1 - \cos \theta_1 \end{pmatrix} \quad (19.1.15)$$

$$= \frac{1}{\sin(\theta_1 + \theta_2)} \begin{pmatrix} -\sin \theta_2 & \cos \theta_2 \\ \sin \theta_1 & \cos \theta_1 \end{pmatrix}. \quad (19.1.16)$$

The tensions  $T_1$  and  $T_2$  are therefore

$$\begin{pmatrix} T_1 \\ T_2 \end{pmatrix} = \frac{1}{\sin(\theta_1 + \theta_2)} \begin{pmatrix} -\sin \theta_2 & \cos \theta_2 \\ \sin \theta_1 & \cos \theta_1 \end{pmatrix} \begin{pmatrix} 0 \\ T_3 \end{pmatrix} \quad (19.1.17)$$

$$= \frac{T_3}{\sin(\theta_1 + \theta_2)} \begin{pmatrix} \cos \theta_2 \\ \cos \theta_1 \end{pmatrix}. \quad (19.1.18)$$

Recall that we've already found  $T_3 = mg$ ; then the final results are

$$T_1 = \frac{mg \cos \theta_2}{\sin(\theta_1 + \theta_2)}, \quad (19.1.19)$$

$$T_2 = \frac{mg \cos \theta_1}{\sin(\theta_1 + \theta_2)}, \quad (19.1.20)$$

$$T_3 = mg. \quad (19.1.21)$$

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