

42.1: Introduction to Simple Harmonic Motion

The small-angle approximation of the simple plane pendulum is an example of what is called simple harmonic motion. Simple harmonic motion is the motion that a particle exhibits when under the influence of a force of the form given by Hooke's law (named for the 17th century English scientist Robert Hooke):

$$F = -kx. \quad (42.1.1)$$

A force of this form describes, for example, the force on a mass attached to a spring with spring constant k , where k is a measure of the stiffness of the spring. In this case F is the force exerted by the spring, and x is the distance of the mass from its equilibrium position - that is, the "resting" position at which the mass can be left where it will not oscillate.

Substituting Hooke's law as the force in Newton's second law $F = ma$ (and recalling the acceleration $a = d^2x/dt^2$) gives the equation

$$-kx = m \frac{d^2x}{dt^2} \quad (42.1.2)$$

This is a second-order linear differential equation with constant coefficients, and can be solved for $x(t)$ using standard methods from the theory of differential equations. We won't go into the theory of differential equations here, but just present the result. The solution is

$$x(t) = A \cos(\omega t + \delta). \quad (42.1.3)$$

Here ω is called the angular frequency of the motion, and measures how fast the particle oscillates back and forth. The constant A is called the amplitude of the motion, and is the maximum distance the particle travels from its equilibrium position, $x = 0$. The constant δ called the phase constant, and determines where in its cycle the particle is at time $t = 0$. A plot of $x(t)$ is shown in Fig. 42.1.1.

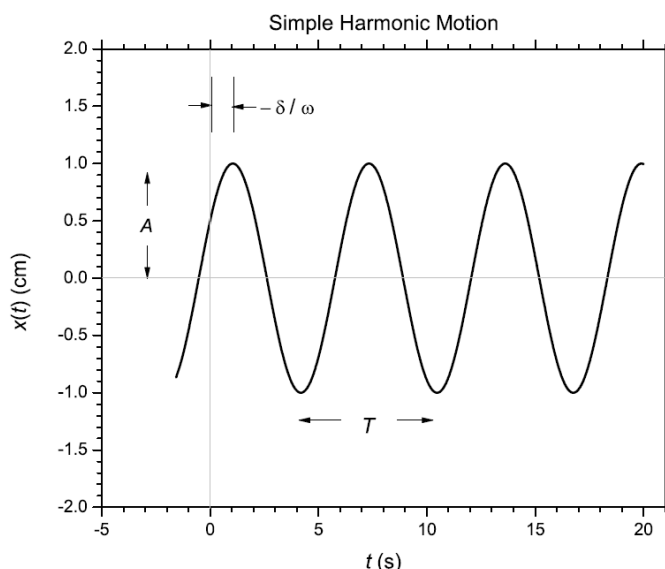


Figure 42.1.1: Simple harmonic motion. Shown are the amplitude A , period T , and phase constant δ . The horizontal line $x(t) = 0$ is the equilibrium position.

Since the sine and cosine function differ only by a phase shift ($\sin \theta \equiv \cos(\theta - \pi/2)$), we could replace the cosine function in Eq. 42.1.3 with a sine by simply adding an extra $\pi/2$ to the phase constant δ . So either the sine or the cosine can be used equally well to describe simple harmonic motion; here we will choose to use the cosine function.

The calculus may also be used to find the velocity of the particle at any time t ; the result is

$$v(t) = -A\omega \sin(\omega t + \delta) \quad (42.1.4)$$

so that the maximum speed of the simple harmonic oscillator is

$$|v_{\max}| = A\omega \quad (42.1.5)$$

Further, it can be shown that the acceleration at any time t is

$$a(t) = -A\omega^2 \cos(\omega t + \delta) \quad (42.1.6)$$

$$= -\omega^2 x(t). \quad (42.1.7)$$

Multiplying Eq. 42.1.7 by the particle mass m , we find

$$ma(t) = F(t) = -m\omega^2 x(t). \quad (42.1.8)$$

Comparing this with Eq. 42.1.1 we see that

$$k = m\omega^2, \quad (42.1.9)$$

or

$$\omega = \sqrt{\frac{k}{m}}. \quad (42.1.10)$$

In Eq. 42.1.3, the amplitude A depends on how far the particle was displaced from equilibrium before being released; the phase constant δ just depends on when we choose time $t = 0$; but the angular frequency ω depends on the physical parameters of the system: the stiffness of the spring k and the mass of the particle m .

42.1: Introduction to Simple Harmonic Motion is shared under a [CC BY-NC-SA 4.0](https://creativecommons.org/licenses/by-nc-sa/4.0/) license and was authored, remixed, and/or curated by LibreTexts.